

Classical Mechanics Homework
January 30, 2008
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1 The 2-Body Problem

The goal of this problem is to understand a pair of particles interacting via a central force such as gravity. We'll reduce it to problem you've already studied — the case of a *single* particle in a central force.

Suppose we have a system of two particles interacting by a central force. Their positions are functions of time, say $q_1, q_2: \mathbb{R} \rightarrow \mathbb{R}^3$, satisfying Newton's law:

$$m_1 \ddot{q}_1 = f(|q_1 - q_2|) \frac{q_1 - q_2}{|q_1 - q_2|}$$
$$m_2 \ddot{q}_2 = f(|q_2 - q_1|) \frac{q_2 - q_1}{|q_2 - q_1|}.$$

Here m_1, m_2 are their masses, and the force is described by some smooth function $f: (0, \infty) \rightarrow \mathbb{R}$. Let's write the force in terms of a potential as follows:

$$f(r) = -\frac{dV}{dr}.$$

Using conservation of momentum and symmetry under translations and Galilei boosts we can work in coordinates where

$$(1) \quad m_1 q_1(t) + m_2 q_2(t) = 0$$

for all times t . This coordinate system is called the **center-of-mass frame**.

We could use equation (1) to solve for q_2 in terms of q_1 , or vice versa, but we can also use it to express both q_1 and q_2 in terms of the **relative position**

$$(2) \quad q(t) = q_1(t) - q_2(t).$$

This is more symmetrical, so this is what we will do. Henceforth we only need to talk about q . Thus we have reduced the problem to a 1-body problem!

Now here's where you come in:

1. Show that $q(t)$ satisfies the equation

$$m \ddot{q} = f(|q|) \frac{q}{|q|}$$

where m is the so-called **reduced mass**

$$m = \frac{m_1 m_2}{m_1 + m_2}.$$

Note that this looks exactly like Newton's second law for a single particle!

Solution to 1:

From (2) we can immediately see that $\ddot{q} = \ddot{q}_1 - \ddot{q}_2$. We can re-write \ddot{q}_1 and \ddot{q}_2 using their respective equations of motion. This leaves us with:

$$\ddot{q}_1 = f(|q_1 - q_2|) \frac{q_1 - q_2}{m_1 |q_1 - q_2|} \quad \text{and} \quad \ddot{q}_2 = f(|q_2 - q_1|) \frac{q_2 - q_1}{m_1 |q_2 - q_1|}.$$

Subtracting these two equations we get an expression for \ddot{q} as follows:

$$\ddot{q} = f(|q_1 - q_2|) \frac{q_1 - q_2}{m_1 |q_1 - q_2|} - f(|q_2 - q_1|) \frac{q_2 - q_1}{m_2 |q_2 - q_1|}.$$

Now since $|q_1 - q_2| = |q_2 - q_1|$ the above expression can be written as:

$$\ddot{q} = f(|q|) \frac{q_1 - q_2}{m_1 |q|} - f(|q|) \frac{q_2 - q_1}{m_2 |q|} = \frac{f(|q|)}{|q|} \left(\frac{q_1 - q_2}{m_1} - \frac{q_2 - q_1}{m_2} \right).$$

Combining q_i terms of equal i we get:

$$\ddot{q} = \frac{f(|q|)}{|q|} \left[\left(\frac{m_1 + m_2}{m_1 m_2} \right) q_1 - \left(\frac{m_1 + m_2}{m_1 m_2} \right) q_2 \right].$$

The coefficients in front of q_1 and q_2 are just the inverse of the reduced mass.

$$\frac{1}{m} = \frac{m_1 + m_2}{m_1 m_2}$$

Thus we have shown that:

$$\begin{aligned} \ddot{q} &= \frac{1}{m} \frac{f(|q|)}{|q|} (q_1 - q_2) \\ \Rightarrow m\ddot{q} &= f(|q|) \frac{q}{|q|}. \end{aligned}$$

2. Recall that the total energy E of the 2-particle system is the sum of the kinetic energies of the particles plus the potential energy. Express E in terms of q and the reduced mass. Show that

$$E = \frac{1}{2} m |\dot{q}|^2 + V(|q|)$$

Note that this looks exactly like the energy of a single particle!

Solution to 2:

The total energy E for the two particle system is given by:

$$(3) \quad E = \frac{1}{2} m_1 \dot{q}_1^2 + \frac{1}{2} m_2 \dot{q}_2^2 + V(|q_1 - q_2|)$$

From (1) and (2) we express \dot{q}_1 and \dot{q}_2 in terms of \dot{q} , the mass m_1 and the mass m_2 by noting the following:

$$m_1 \dot{q}_1 + m_2 \dot{q}_2 = 0 \quad \Rightarrow \quad \dot{q}_1 = -\frac{m_2}{m_1} \dot{q}_2 \quad \text{and} \quad \dot{q}_2 = -\frac{m_1}{m_2} \dot{q}_1$$

$$(4a) \quad \dot{q}_1 = -\frac{m_2}{m_1} \dot{q}_2 \quad \text{and} \quad \dot{q}_1 - \dot{q}_2 = \dot{q} \quad \Rightarrow \quad \dot{q}_2 + \frac{m_2}{m_1} \dot{q}_2 = -\dot{q} \quad \Rightarrow \quad \dot{q}_2 = \frac{-m_1}{m_1 + m_2} \dot{q}$$

$$(4b) \quad \dot{q}_2 = -\frac{m_1}{m_2} \dot{q}_1 \quad \text{and} \quad \dot{q}_1 - \dot{q}_2 = \dot{q} \quad \Rightarrow \quad \dot{q}_1 + \frac{m_1}{m_2} \dot{q}_1 = \dot{q} \quad \Rightarrow \quad \dot{q}_1 = \frac{m_2}{m_1 + m_2} \dot{q}$$

Substituting the right most equalities of (4a) and (4b) for their respective values in (3) and noting that $|q_1 - q_2| = |q|$ we get:

$$(5) \quad E = \frac{1}{2} \frac{m_1 m_2^2}{(m_1 + m_2)^2} \dot{q}^2 + \frac{1}{2} \frac{m_2 m_1^2}{(m_1 + m_2)^2} \dot{q}^2 + V(|q|).$$

Finally, summing the two kinetic terms explicitly gives the desired result.

$$E = \frac{1}{2} m |\dot{q}|^2 + V(|q|)$$

Where as before m is the reduced mass of the system.

3. Let J be the total angular momentum of the 2-particle system. Show that

$$J = m q \times \dot{q}$$

Note that this looks exactly like the angular momentum of a single particle!

Solution to 3:

Solving for the positions q_1 and q_2 , gives identical equations to (4a) and (4b) only without time derivatives. We summarize the results below:

$$(6a) \quad q_1 = \frac{m_2}{m_1 + m_2} q \quad \dot{q}_1 = \frac{m_2}{m_1 + m_2} \dot{q}$$

$$(6b) \quad q_2 = \frac{-m_1}{m_1 + m_2} q \quad \dot{q}_2 = \frac{-m_1}{m_1 + m_2} \dot{q}$$

The total angular momentum, J , is the sum of the angular momentum of each particle. We give the full expression below.

$$(7) \quad J = J_1 + J_2 = m_1 q_1 \times \dot{q}_1 + m_2 q_2 \times \dot{q}_2$$

Again, inserting (6a) and (6b) into (7) in favor of q we get:

$$\begin{aligned} J &= m_1 q_1 \times \dot{q}_1 + m_2 q_2 \times \dot{q}_2 \\ &= m_1 \frac{m_2}{m_1 + m_2} q \times \frac{m_2}{m_1 + m_2} \dot{q} + m_2 \frac{-m_1}{m_1 + m_2} \times \frac{-m_1}{m_1 + m_2} \dot{q} \\ &= \frac{m_1 m_2^2}{(m_1 + m_2)^2} q \times \dot{q} + \frac{m_2 m_1^2}{(m_1 + m_2)^2} \times \dot{q} \\ &= \frac{m_1 m_2}{(m_1 + m_2)} \frac{m_1 + m_2}{(m_1 + m_2)} q \times \dot{q} \end{aligned}$$

$$J = m q \times \dot{q}$$

This proves the claim in question 3.

At this point we're back to a problem you've already solved: a *single* particle in a central force. The only difference is that now q stands for the *relative* position and m stands for the *reduced* mass!

So, we instantly conclude that two bodies orbiting each other due to the force of gravity will *both* have an orbit that's either an ellipse, or a parabola, or a hyperbola... when viewed in the center-of-mass frame.

2 Poisson brackets

Let \mathbb{R}^{2n} be the **phase space** of a particle in \mathbb{R}^n , with coordinates q_i, p_i ($1 \leq i \leq n$). Let $C^\infty(\mathbb{R}^{2n})$ be the set of smooth real-valued functions on \mathbb{R}^{2n} , which becomes an commutative algebra using pointwise addition and multiplication of functions.

We define the **Poisson bracket** of functions $F, G \in C^\infty(\mathbb{R}^{2n})$ by:

$$\{F, G\} = \sum_{i=1}^n \frac{\partial F}{\partial p_i} \frac{\partial G}{\partial q_i} - \frac{\partial G}{\partial p_i} \frac{\partial F}{\partial q_i}.$$

4. Show that Poisson brackets make the vector space $C^\infty(\mathbb{R}^{2n})$ into a **Lie algebra**. In other words, check the **antisymmetry** of the bracket:

$$\{F, G\} = -\{G, F\}$$

the **bilinearity** of the bracket:

$$\{F, \alpha G + \beta H\} = \alpha \{F, G\} + \beta \{F, H\}$$

$$\{\alpha F + \beta G, H\} = \alpha \{F, H\} + \beta \{G, H\}$$

and **Jacobi identity**:

$$\{F, \{G, H\}\} = \{\{F, G\}, H\} + \{G, \{F, H\}\}$$

for all $F, G, H \in C^\infty(\mathbb{R}^{2n})$ and $\alpha, \beta \in \mathbb{R}$.

(Note the Jacobi identity resembles the product rule $d(GH) = (dG)H + GdH$, with bracketing by F playing the role of d . This is no accident!)

Solution to 4:

Since,

$$\{F, G\} = \sum_{i=1}^n \frac{\partial F}{\partial p_i} \frac{\partial G}{\partial q_i} - \frac{\partial G}{\partial p_i} \frac{\partial F}{\partial q_i},$$

we can calculate the bracket when F and G are interchanged. We get:

$$\{G, F\} = \sum_{i=1}^n \frac{\partial G}{\partial p_i} \frac{\partial F}{\partial q_i} - \frac{\partial F}{\partial p_i} \frac{\partial G}{\partial q_i}.$$

But then it becomes clear that:

$$-\{G, F\} = -\left[\sum_{i=1}^n \frac{\partial G}{\partial p_i} \frac{\partial F}{\partial q_i} - \frac{\partial F}{\partial p_i} \frac{\partial G}{\partial q_i} \right] = \sum_{i=1}^n -\frac{\partial G}{\partial p_i} \frac{\partial F}{\partial q_i} + \frac{\partial F}{\partial p_i} \frac{\partial G}{\partial q_i} = \sum_{i=1}^n \frac{\partial F}{\partial p_i} \frac{\partial G}{\partial q_i} - \frac{\partial G}{\partial p_i} \frac{\partial F}{\partial q_i} = \{F, G\}$$

Proving the **antisymmetry** of the bracket, namely that:

$$\{F, G\} = -\{G, F\}.$$

Next, we investigate the **bilinearity**, writing $\{F, \alpha G + \beta H\}$ explicitly gives:

$$\begin{aligned}
\{F, \alpha G + \beta H\} &= \sum_{i=1}^n \frac{\partial F}{\partial p_i} \frac{\partial(\alpha G + \beta H)}{\partial q_i} - \frac{\partial(\alpha G + \beta H)}{\partial p_i} \frac{\partial F}{\partial q_i} \\
&= \sum_{i=1}^n \frac{\partial F}{\partial p_i} \left[\alpha \frac{\partial G}{\partial q_i} + \beta \frac{\partial H}{\partial q_i} \right] - \left[\alpha \frac{\partial G}{\partial p_i} + \beta \frac{\partial H}{\partial p_i} \right] \frac{\partial F}{\partial q_i} \\
&= \sum_{i=1}^n \alpha \frac{\partial F}{\partial p_i} \frac{\partial G}{\partial q_i} + \beta \frac{\partial F}{\partial p_i} \frac{\partial H}{\partial q_i} - \alpha \frac{\partial G}{\partial p_i} \frac{\partial F}{\partial q_i} - \beta \frac{\partial H}{\partial p_i} \frac{\partial F}{\partial q_i} \\
&= \sum_{i=1}^n \alpha \frac{\partial F}{\partial p_i} \frac{\partial G}{\partial q_i} - \alpha \frac{\partial G}{\partial p_i} \frac{\partial F}{\partial q_i} + \beta \frac{\partial F}{\partial p_i} \frac{\partial H}{\partial q_i} - \beta \frac{\partial H}{\partial p_i} \frac{\partial F}{\partial q_i} \\
&= \alpha \sum_{i=1}^n \frac{\partial F}{\partial p_i} \frac{\partial G}{\partial q_i} - \frac{\partial G}{\partial p_i} \frac{\partial F}{\partial q_i} + \beta \sum_{i=1}^n \frac{\partial F}{\partial p_i} \frac{\partial H}{\partial q_i} - \frac{\partial H}{\partial p_i} \frac{\partial F}{\partial q_i} \\
&= \alpha \{F, H\} + \beta \{G, H\}
\end{aligned}$$

For completeness we must now check, $\{\alpha F + \beta G, H\}$, as well.

$$\begin{aligned}
\{\alpha F + \beta G, H\} &= \sum_{i=1}^n \frac{\partial(\alpha F + \beta G)}{\partial p_i} \frac{\partial H}{\partial q_i} - \frac{\partial H}{\partial p_i} \frac{\partial(\alpha F + \beta G)}{\partial q_i} \\
&= \sum_{i=1}^n \left[\alpha \frac{\partial F}{\partial p_i} + \beta \frac{\partial G}{\partial p_i} \right] \frac{\partial H}{\partial q_i} - \frac{\partial H}{\partial p_i} \left[\alpha \frac{\partial F}{\partial q_i} + \beta \frac{\partial G}{\partial q_i} \right] \\
&= \sum_{i=1}^n \alpha \frac{\partial F}{\partial p_i} \frac{\partial H}{\partial q_i} + \beta \frac{\partial G}{\partial p_i} \frac{\partial H}{\partial q_i} - \alpha \frac{\partial H}{\partial p_i} \frac{\partial F}{\partial q_i} - \beta \frac{\partial H}{\partial p_i} \frac{\partial G}{\partial q_i} \\
&= \sum_{i=1}^n \alpha \frac{\partial F}{\partial p_i} \frac{\partial H}{\partial q_i} - \alpha \frac{\partial H}{\partial q_i} \frac{\partial F}{\partial p_i} + \beta \frac{\partial G}{\partial p_i} \frac{\partial H}{\partial q_i} - \beta \frac{\partial H}{\partial q_i} \frac{\partial G}{\partial p_i} \\
&= \alpha \sum_{i=1}^n \frac{\partial F}{\partial p_i} \frac{\partial H}{\partial q_i} - \frac{\partial H}{\partial q_i} \frac{\partial F}{\partial p_i} + \beta \sum_{i=1}^n \frac{\partial G}{\partial p_i} \frac{\partial H}{\partial q_i} - \frac{\partial H}{\partial q_i} \frac{\partial G}{\partial p_i} \\
&= \alpha \{F, H\} + \beta \{G, H\}
\end{aligned}$$

Thus the bracket fulfills the property of **bilinearity**. Lastly, we show that the **Jacobi Identity** holds.

$$\begin{aligned}
\{F, \{G, H\}\} &= \sum_{i=1}^n \frac{\partial F}{\partial p_i} \frac{\partial\{G, H\}}{\partial q_i} - \frac{\partial\{G, H\}}{\partial p_i} \frac{\partial F}{\partial q_i} \\
&= \sum_{i=1}^n \frac{\partial F}{\partial p_i} \frac{\partial}{\partial q_i} \left[\sum_{j=1}^n \frac{\partial G}{\partial p_j} \frac{\partial H}{\partial q_j} - \frac{\partial H}{\partial p_j} \frac{\partial G}{\partial q_j} \right] - \frac{\partial}{\partial p_i} \left[\sum_{j=1}^n \frac{\partial G}{\partial p_j} \frac{\partial H}{\partial q_j} - \frac{\partial H}{\partial p_j} \frac{\partial G}{\partial q_j} \right] \frac{\partial F}{\partial q_i} \\
&= \sum_{i=1}^n \sum_{j=1}^n \left[\frac{\partial F}{\partial p_i} \frac{\partial}{\partial q_i} \left(\frac{\partial G}{\partial p_j} \frac{\partial H}{\partial q_j} - \frac{\partial H}{\partial p_j} \frac{\partial G}{\partial q_j} \right) - \frac{\partial}{\partial p_i} \left(\frac{\partial G}{\partial p_j} \frac{\partial H}{\partial q_j} - \frac{\partial H}{\partial p_j} \frac{\partial G}{\partial q_j} \right) \frac{\partial F}{\partial q_i} \right] \\
&= \sum_{i=1}^n \sum_{j=1}^n \frac{\partial F}{\partial p_i} \left(\frac{\partial^2 G}{\partial p_j \partial q_i} \frac{\partial H}{\partial q_j} + \frac{\partial G}{\partial p_j} \frac{\partial^2 H}{\partial q_j \partial q_i} - \frac{\partial^2 H}{\partial p_j \partial q_i} \frac{\partial G}{\partial q_j} - \frac{\partial H}{\partial p_j} \frac{\partial^2 G}{\partial q_j \partial q_i} \right) \\
&\quad - \left(\frac{\partial^2 G}{\partial p_j \partial p_i} \frac{\partial H}{\partial q_j} + \frac{\partial G}{\partial p_j} \frac{\partial^2 H}{\partial q_j \partial p_i} - \frac{\partial^2 H}{\partial p_j \partial p_i} \frac{\partial G}{\partial q_j} - \frac{\partial H}{\partial p_j} \frac{\partial^2 G}{\partial q_j \partial p_i} \right) \frac{\partial F}{\partial q_i}
\end{aligned}$$

(Again this identity resembles the product rule!)

Solution to 5:

Expanding the bracket of a product below leads to:

$$\begin{aligned}\{F, GH\} &= \sum_{i=1}^n \frac{\partial F}{\partial p_i} \frac{\partial(GH)}{\partial q_i} - \frac{\partial(GH)}{\partial p_i} \frac{\partial F}{\partial q_i} \\ &= \sum_{i=1}^n \frac{\partial F}{\partial p_i} \left[\frac{\partial G}{\partial q_i} H + G \frac{\partial H}{\partial q_i} \right] - \left[\frac{\partial G}{\partial p_i} H + G \frac{\partial H}{\partial p_i} \right] \frac{\partial F}{\partial q_i} \\ &= \sum_{i=1}^n \frac{\partial F}{\partial p_i} \frac{\partial G}{\partial q_i} H + \frac{\partial F}{\partial p_i} G \frac{\partial H}{\partial q_i} - \frac{\partial G}{\partial p_i} H \frac{\partial F}{\partial q_i} - G \frac{\partial H}{\partial p_i} \frac{\partial F}{\partial q_i} \\ &= \sum_{i=1}^n \frac{\partial F}{\partial p_i} \frac{\partial G}{\partial q_i} H - \frac{\partial G}{\partial p_i} \frac{\partial F}{\partial q_i} H + G \frac{\partial F}{\partial p_i} \frac{\partial H}{\partial q_i} - G \frac{\partial H}{\partial p_i} \frac{\partial F}{\partial q_i} \\ &= \sum_{i=1}^n \left[\frac{\partial F}{\partial p_i} \frac{\partial G}{\partial q_i} - \frac{\partial G}{\partial p_i} \frac{\partial F}{\partial q_i} \right] H + \sum_{i=1}^n G \left[\frac{\partial F}{\partial p_i} \frac{\partial H}{\partial q_i} - \frac{\partial H}{\partial p_i} \frac{\partial F}{\partial q_i} \right]\end{aligned}$$

Arriving at:

$$\{F, GH\} = \{F, G\}H + G\{F, H\},$$

quod erat demonstrandum.