Classical Mechanics, Lecture 18 March 11, 2008 lecture by John Baez notes by Alex Hoffnung

1 The Category of Classical Systems

A category is a bunch of things (objects) and processes (morphisms):

 $\bullet \longrightarrow \bullet$

For us, the 'things' are physical systems in classical mechanics: Poisson manifolds. The 'processes' are certain maps between physical systems: Poisson maps.

Definition 1 Given two Poisson manifolds X and Y, a map $\phi: X \to Y$ is **Poisson** if it is smooth and

$$\phi^*: C^{\infty}(Y) \to C^{\infty}(X)$$
$$f \mapsto f\phi$$

preserves the Poisson brackets: given $f, g \in C^{\infty}(Y)$

$$\{\phi^*(f), \phi^*(g)\} = \phi^*\{f, g\}$$

In fact, we have seen lots of Poisson maps. Any physical process in classical mechanics gives a Poisson map.

Example: $X = T^*(\mathbb{R}^n)$ is the phase space for a particle in \mathbb{R}^n ; the Galilei group G(n+1) acts on X:

$$A: G(n+1) \times X \to X$$

and in fact for each $g \in G(n+1)$ the map $A(g): X \to X$ is Poisson!

2

Example: X is any Poisson manifold, $H \in C^{\infty}(X)$. Then we get a vector field on X and if v_H is integrable we get a flow

$$\phi \colon \mathbb{R} \times X \to X$$
$$(t, x) \mapsto \phi_t(x)$$

with

$$\frac{d}{dt}\phi_t(x) = v_H(\phi_t(x))$$

Given this, each map $\phi_t \colon X \to X$ is Poisson!

Check: given $F, G \in C^{\infty}(X)$ we want:

$$\{\phi_t^* F, \phi_t^* G\} = \phi_t^* \{F, G\}$$

This is true for t = 0 since $\phi_0: X \to X$ is the identity, hence $\phi_0^* F = F$. So, it is enough to show:

$$\frac{d}{dt}\left\{\phi_t^*F,\phi_t^*G\right\} = \frac{d}{dt}\phi_t^*\left\{F,G\right\}$$

i.e.

$$\frac{d}{dt} \left\{ \phi_t^* F, \phi_t^* G \right\} (x) = \frac{d}{dt} \phi_t^* \left\{ F, G \right\} (x)$$

$$\frac{d}{dt}\phi_t^* \{F,G\}(x) \quad \stackrel{\text{def. of pullback}}{=} \quad \frac{d}{dt} \{F,G\}(\phi_t(x)) \\
\stackrel{\text{chain rule}}{=} \quad d\{F,G\}\left(\frac{d}{dt}\phi_t(x)\right) \\
\stackrel{\text{def. of }\phi_t}{=} \quad d\{F,G\}v_H(\phi_t(x)) \\
\stackrel{\text{def. of }v_H}{=} \quad v_H \{F,G\}(\phi_t(x)) \\
\stackrel{\text{Jacobi id.}}{=} \quad \{H,\{F,G\}\}(\phi_t(x)) \\
= \quad (\{\{H,F\},G\}+\{F,\{H,G\}\})\phi_t(x)$$

On the other hand:

$$\frac{d}{dt} \left\{ \phi_t^* F, \phi_t^* G \right\} (x) \qquad \begin{array}{l} \text{bilin. of } \left\{ \cdot, \cdot \right\} \text{ gives prod. rule} \\ = & \left\{ \left\{ \frac{d}{dt} \phi_t^* F, \phi_t^* G \right\} + \left\{ \phi_t^* F, \frac{d}{dt} \phi_t^* G \right\} \right\} (x) \\ = & \left\{ \left\{ H, \phi_t^* F \right\}, \phi_t^* G \right\} + \left\{ \phi_t^* F, \left\{ H, \phi_t^* G \right\} \right\} (x) \\ \end{array}$$

This agrees with the other side at t = 0. (Note: $H = \phi_t^* H$ since energy is conserved.) So we have:

$$\frac{d}{dt}\phi_{t}^{*}\left\{F,G\right\}\left(x\right)\Big|_{t=0} = \frac{d}{dt}\left\{\phi_{t}^{*}F,\phi_{t}^{*}G\right\}\left(x\right)\Big|_{t=0}$$

for all x, so use $\phi_s(x) \in X$:

$$\frac{d}{dt}\phi_t^* \{F, G\}\phi_s(x)\Big|_{t=0} = \frac{d}{dt} \{\phi_t^*F, \phi_t^*G\}(\phi_s(x))\Big|_{t=0}$$
$$\frac{d}{dt}\{F, G\}\phi_{t+s}(x)\Big|_{t=0} = \frac{d}{ds}\{F, G\}\phi_s(x) =$$

We would like:

$$\frac{d}{ds}\left\{F,G\right\}\phi_s(x) = \frac{d}{ds}\left\{\phi_s^*F,\phi_s^*G\right\}$$

Ugh!

So: we have a category of Poisson manifolds and Poisson maps and time evolution for any Hamiltonian is a Poisson map.

Definition 2 A category consists of a collection of objects and for any pair of objects X and Y a set of morphisms $f: X \to Y$ such that given $f: X \to Y$ and $g: Y \to Z$ we have a morphism $gf: X \to Z$, such that:

- 1. (hg)f = h(gf)
- 2. each X has an identity morphism $1_X: X \to X$ such that:

$$f1_X = f$$
$$1_X g = g$$