

Fractional Index Theory

$$\text{Index}_a(\not\partial^+) = \int_Z \widehat{A}(Z) \in \mathbb{Q}$$

Workshop on Geometry and Lie Groups

The University of Hong Kong Institute of Mathematical Research

26 March 2011

Mathai Varghese

School of Mathematical Sciences



[MMS4]

V. Mathai, R.B. Melrose and I.M. Singer,

Equivariant and fractional index of projective elliptic operators,

J. Differential Geometry, **78** (2008), no.3, 465-473.

[\[math.DG/0611819\]](#).

[MMS4]

V. Mathai, R.B. Melrose and I.M. Singer,

Equivariant and fractional index of projective elliptic operators,

J. Differential Geometry, **78** (2008), no.3, 465-473.

[[math.DG/0611819](#)].

[MMS3]

V. Mathai, R.B. Melrose and I.M. Singer,

Fractional Analytic Index,

J. Differential Geometry, **74**, (2006), no. 2, 265-292.

[[math.DG/0402329](#)]

The Atiyah-Singer Index Theorem

In the 1960s, **Atiyah** and **Singer** proved what was to become one of the most important and widely applied theorems in 20th century mathematics, viz. the **Atiyah-Singer Index Theorem**.

The Atiyah-Singer Index Theorem

In the 1960s, **Atiyah** and **Singer** proved what was to become one of the most important and widely applied theorems in 20th century mathematics, viz. the **Atiyah-Singer Index Theorem**. Roughly speaking, the laws of nature are often expressed in terms of differential equations, which if **elliptic**, have an **index** being the number of solutions minus the number of constraints imposed.

The Atiyah-Singer Index Theorem

In the 1960s, **Atiyah** and **Singer** proved what was to become one of the most important and widely applied theorems in 20th century mathematics, viz. the **Atiyah-Singer Index Theorem**. Roughly speaking, the laws of nature are often expressed in terms of differential equations, which if **elliptic**, have an **index** being the number of solutions minus the number of constraints imposed. The Atiyah-Singer Index Theorem gives a striking calculation of this **index**.

The Atiyah-Singer Index Theorem

In the 1960s, **Atiyah** and **Singer** proved what was to become one of the most important and widely applied theorems in 20th century mathematics, viz. the **Atiyah-Singer Index Theorem**. Roughly speaking, the laws of nature are often expressed in terms of differential equations, which if **elliptic**, have an **index** being the number of solutions minus the number of constraints imposed. The Atiyah-Singer Index Theorem gives a striking calculation of this **index**. Atiyah and Singer were jointly awarded the prestigious **Abel Prize** in 2004.

The Atiyah-Singer Index Theorem

In the 1960s, **Atiyah** and **Singer** proved what was to become one of the most important and widely applied theorems in 20th century mathematics, viz. the **Atiyah-Singer Index Theorem**. Roughly speaking, the laws of nature are often expressed in terms of differential equations, which if **elliptic**, have an **index** being the number of solutions minus the number of constraints imposed. The Atiyah-Singer Index Theorem gives a striking calculation of this **index**. Atiyah and Singer were jointly awarded the prestigious **Abel Prize** in 2004.



The Atiyah-Singer Index Theorem

In the 1960s, **Atiyah** and **Singer** proved what was to become one of the most important and widely applied theorems in 20th century mathematics, viz. the **Atiyah-Singer Index Theorem**. Roughly speaking, the laws of nature are often expressed in terms of differential equations, which if **elliptic**, have an **index** being the number of solutions minus the number of constraints imposed. The Atiyah-Singer Index Theorem gives a striking calculation of this **index**. Atiyah and Singer were jointly awarded the prestigious **Abel Prize** in 2004.



The Atiyah-Singer Index Theorem

In the 1960s, **Atiyah** and **Singer** proved what was to become one of the most important and widely applied theorems in 20th century mathematics, viz. the **Atiyah-Singer Index Theorem**. Roughly speaking, the laws of nature are often expressed in terms of differential equations, which if **elliptic**, have an **index** being the number of solutions minus the number of constraints imposed. The Atiyah-Singer Index Theorem gives a striking calculation of this **index**. Atiyah and Singer were jointly awarded the prestigious **Abel Prize** in 2004.



Dirac operators

Paul Dirac was one of the founders of quantum mechanics, and was awarded the **Nobel Prize in Physics** in 1933.

Dirac operators

Paul Dirac was one of the founders of quantum mechanics, and was awarded the **Nobel Prize in Physics** in 1933.



Dirac operators

Paul Dirac was one of the founders of quantum mechanics, and was awarded the **Nobel Prize in Physics** in 1933.



Dirac defined an operator $\not{\partial}$ on \mathbb{R}^n that solved the **square root problem** for the Laplacian on \mathbb{R}^n , that is, $\not{\partial}^2 = \Delta$.

Dirac operators

Paul Dirac was one of the founders of quantum mechanics, and was awarded the **Nobel Prize in Physics** in 1933.



Dirac defined an operator $\not{\partial}$ on \mathbb{R}^n that solved the **square root problem** for the Laplacian on \mathbb{R}^n , that is, $\not{\partial}^2 = \Delta$. The construction was novel as it used **Clifford algebras and spinors** in an essential way.

Dirac operators

More precisely, if $\{\gamma_j\}_{j=1}^n$ denote Clifford multiplication by an orthonormal basis of \mathbb{R}^n , then the Clifford algebra relations are

$$\gamma_j \gamma_k + \gamma_k \gamma_j = 2\delta_{ij}.$$

Dirac operators

More precisely, if $\{\gamma_j\}_{j=1}^n$ denote Clifford multiplication by an orthonormal basis of \mathbb{R}^n , then the Clifford algebra relations are $\gamma_j\gamma_k + \gamma_k\gamma_j = 2\delta_{ij}$. When $n = 2$, these are **Pauli matrices**;

$$\gamma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \gamma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \epsilon = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Dirac operators

More precisely, if $\{\gamma_j\}_{j=1}^n$ denote Clifford multiplication by an orthonormal basis of \mathbb{R}^n , then the Clifford algebra relations are $\gamma_j\gamma_k + \gamma_k\gamma_j = 2\delta_{ij}$. When $n = 2$, these are **Pauli matrices**;

$$\gamma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \gamma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \epsilon = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\not\partial = \sum_{j=1}^n \gamma_j \frac{\partial}{\partial x_j}.$$

It turns out that this operator plays a fundamental role in quantum mechanics, and is known as the **Dirac operator**.

Dirac operators

More precisely, if $\{\gamma_j\}_{j=1}^n$ denote Clifford multiplication by an orthonormal basis of \mathbb{R}^n , then the Clifford algebra relations are $\gamma_j\gamma_k + \gamma_k\gamma_j = 2\delta_{ij}$. When $n = 2$, these are **Pauli matrices**;

$$\gamma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \gamma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \epsilon = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\not{\partial} = \sum_{j=1}^n \gamma_j \frac{\partial}{\partial x_j}.$$

It turns out that this operator plays a fundamental role in quantum mechanics, and is known as the **Dirac operator**. By construction,

$$\not{\partial}^2 = \Delta \cdot \text{Id}.$$

The Index Theorem for Dirac operators

Atiyah and Singer extended the definition of the **Dirac operator**, \not{D}^+ on any compact **spin** manifold Z of even dimension, and computed the analytic index,

$$\begin{aligned}\text{Index}_a(\not{D}^+) &= \dim(\text{nullspace } \not{D}^+) - \dim(\text{nullspace } \not{D}^-) \\ &= \int_Z \widehat{A}(Z) \in \mathbb{Z}\end{aligned}$$

The Index Theorem for Dirac operators

Atiyah and Singer extended the definition of the **Dirac operator**, \not{D}^+ on any compact **spin** manifold Z of even dimension, and computed the analytic index,

$$\begin{aligned}\text{Index}_a(\not{D}^+) &= \dim(\text{nullspace } \not{D}^+) - \dim(\text{nullspace } \not{D}^-) \\ &= \int_Z \widehat{A}(Z) \in \mathbb{Z}\end{aligned}$$

where RHS is the A-hat genus of the manifold Z . In terms of the Riemannian curvature Ω_Z of Z , $\widehat{A}(Z) = \sqrt{\det \left(\frac{\frac{1}{4\pi} \Omega_Z}{\sinh(\frac{1}{4\pi} \Omega_Z)} \right)}$.

The Index Theorem for Dirac operators

Atiyah and Singer extended the definition of the **Dirac operator**, \not{D}^+ on any compact **spin** manifold Z of even dimension, and computed the analytic index,

$$\begin{aligned}\text{Index}_a(\not{D}^+) &= \dim(\text{nullspace } \not{D}^+) - \dim(\text{nullspace } \not{D}^-) \\ &= \int_Z \hat{A}(Z) \in \mathbb{Z}\end{aligned}$$

where RHS is the A-hat genus of the manifold Z . In terms of the Riemannian curvature Ω_Z of Z , $\hat{A}(Z) = \sqrt{\det\left(\frac{\frac{1}{4\pi}\Omega_Z}{\sinh(\frac{1}{4\pi}\Omega_Z)}\right)}$.

?Question?: Since $\int_Z \hat{A}(Z) \notin \mathbb{Z}$ continues to make sense for **non-spin** manifolds Z , **what corresponds to the analytic index** in this situation, since the usual Dirac operator does not exist?

Outline of talk

We propose **2** solutions to the question, and relate them.

Outline of talk

We propose **2** solutions to the question, and relate them.

- 1 In **[MMS3]**, we generalize the notion of " ψ do", to "projective ψ do". In particular, on an oriented even dimensional Riemannian manifold, we define the notion of **projective spin Dirac operator**. We define its fractional analytic index, and prove an **index theorem** showing that it equals the \widehat{A} -genus (proof sketched in the talk).

Outline of talk

We propose **2** solutions to the question, and relate them.

- 1 In **[MMS3]**, we generalize the notion of " ψ do", to "projective ψ do". In particular, on an oriented even dimensional Riemannian manifold, we define the notion of **projective spin Dirac operator**. We define its fractional analytic index, and prove an **index theorem** showing that it equals the \widehat{A} -genus (proof sketched in the talk).
- 2 On the oriented orthonormal frame bundle of such a manifold, we show in **[MMS4]** that there also always exists a Spin-equivariant transversally elliptic Dirac operator. The relation between the **fractional analytic index** of the projective Dirac operator and the **equivariant index** of the associated Spin-equivariant transversally elliptic Dirac operator is explained there and sketched in the talk.

Projective vector bundles

A projective vector bundle on a manifold Z is **not** a global bundle on Z , but rather it is a vector bundle $E \rightarrow Y$, where

$$PU(n) \longrightarrow Y \xrightarrow{\phi} Z$$

is a principal $PU(n)$ -bundle,

Projective vector bundles

A projective vector bundle on a manifold Z is **not** a global bundle on Z , but rather it is a vector bundle $E \rightarrow Y$, where

$$PU(n) \longrightarrow Y \xrightarrow{\phi} Z$$

is a principal $PU(n)$ -bundle, where E also satisfies

$$\mathcal{L}_g \otimes E_y \cong E_{g.y}, \quad g \in PU(n), y \in Y \quad (1)$$

Projective vector bundles

A projective vector bundle on a manifold Z is **not** a global bundle on Z , but rather it is a vector bundle $E \rightarrow Y$, where

$$PU(n) \longrightarrow Y \xrightarrow{\phi} Z$$

is a principal $PU(n)$ -bundle, where E also satisfies

$$\mathcal{L}_g \otimes E_y \cong E_{g.y}, \quad g \in PU(n), y \in Y \quad (1)$$

where $\mathcal{L} = U(n) \times_{U(1)} \mathbb{C} \rightarrow PU(n)$ is the **primitive line bundle**,

$$\mathcal{L}_{g_1} \otimes \mathcal{L}_{g_2} \cong \mathcal{L}_{g_1.g_2}, \quad g_i \in PU(n).$$

Projective vector bundles

A projective vector bundle on a manifold Z is **not** a global bundle on Z , but rather it is a vector bundle $E \rightarrow Y$, where

$$PU(n) \longrightarrow Y \xrightarrow{\phi} Z$$

is a principal $PU(n)$ -bundle, where E also satisfies

$$\mathcal{L}_g \otimes E_y \cong E_{g.y}, \quad g \in PU(n), y \in Y \quad (1)$$

where $\mathcal{L} = U(n) \times_{U(1)} \mathbb{C} \rightarrow PU(n)$ is the **primitive line bundle**,

$$\mathcal{L}_{g_1} \otimes \mathcal{L}_{g_2} \cong \mathcal{L}_{g_1.g_2}, \quad g_i \in PU(n).$$

The identification (1) gives a **projective action** of $PU(n)$ on E , i.e. an action of $U(n)$ on E s.t. the center $U(1)$ acts as scalars.

The **Dixmier-Douady invariant** of Y ,

$$DD(Y) = \delta(Y) \in \text{Torsion}(H^3(Z, \mathbb{Z}))$$

is the obstruction to lifting the principal $PU(n)$ -bundle Y to a principal $U(n)$ -bundle. (The construction also works for any principal G bundle P over Z , together with a central extension \widehat{G} of G .)

The **Dixmier-Douady invariant** of Y ,

$$DD(Y) = \delta(Y) \in \text{Torsion}(H^3(Z, \mathbb{Z}))$$

is the obstruction to lifting the principal $PU(n)$ -bundle Y to a principal $U(n)$ -bundle. (The construction also works for any principal G bundle P over Z , together with a central extension \widehat{G} of G .) The **associated algebra bundle**

$$\mathcal{A} = Y \times_{PU(n)} M_n(\mathbb{C})$$

is called the associated **Azumaya bundle**.

Projective vector bundle of spinors

Let $E \rightarrow Z$ be a real oriented Riemannian vector bundle,

$$SO(n) \longrightarrow SO(E) \xrightarrow{\psi} Z$$

the principal bundle of oriented orthonormal frames on E .

Projective vector bundle of spinors

Let $E \rightarrow Z$ be a real oriented Riemannian vector bundle,

$$SO(n) \longrightarrow SO(E) \xrightarrow{\psi} Z$$

the principal bundle of oriented orthonormal frames on E .

Let N denote the (co)normal bundle to the fibres. Then it is easy to see that $w_2(N) = 0$, so that N always has a bundle of spinors \mathcal{S} , which is a projective vector bundle over Z .

Projective vector bundle of spinors

Let $E \rightarrow Z$ be a real oriented Riemannian vector bundle,

$$SO(n) \longrightarrow SO(E) \xrightarrow{\psi} Z$$

the principal bundle of oriented orthonormal frames on E .

Let N denote the (co)normal bundle to the fibres. Then it is easy to see that $w_2(N) = 0$, so that N always has a bundle of spinors \mathcal{S} , which is a projective vector bundle over Z .

Also $End(\mathcal{S}) \cong \psi^* Cl(E)$.

Schwartz kernel theorem

For a compact manifold, Z , and vector bundles E and F over Z , the **Schwartz kernel theorem** gives a 1-1 correspondence,

continuous linear operators, $C^\infty(Z, E) \longrightarrow C^{-\infty}(Z, F)$



distributional sections, $C^{-\infty}(Z^2, \text{Hom}(E, F) \otimes \Omega_R)$

Schwartz kernel theorem

For a compact manifold, Z , and vector bundles E and F over Z , the **Schwartz kernel theorem** gives a 1-1 correspondence,

continuous linear operators, $C^\infty(Z, E) \longrightarrow C^{-\infty}(Z, F)$



distributional sections, $C^{-\infty}(Z^2, \text{Hom}(E, F) \otimes \Omega_R)$

where $\text{Hom}(E, F)_{(z, z')} = F_z \boxtimes E_{z'}^*$ is the 'big' homomorphism bundle over Z^2 and Ω_R the density bundle from the right factor.

When restricted to **pseudodifferential operators**, $\Psi^m(Z, E, F)$, get an isomorphism with the space of **conormal distributions** with respect to the diagonal, $I^m(Z^2, \Delta; \text{Hom}(E, F))$. i.e.

$$\Psi^m(Z, E, F) \iff I^m(Z^2, \Delta; \text{Hom}(E, F))$$

When restricted to **pseudodifferential operators**, $\Psi^m(Z, E, F)$, get an isomorphism with the space of **conormal distributions** with respect to the diagonal, $I^m(Z^2, \Delta; \text{Hom}(E, F))$. i.e.

$$\Psi^m(Z, E, F) \iff I^m(Z^2, \Delta; \text{Hom}(E, F))$$

When further restricted to **differential operators** $\text{Diff}^m(Z, E, F)$ (which by definition have the property of being local operators) this becomes an isomorphism with the space of conormal distributions, $I^m_{\Delta}(Z^2, \Delta; \text{Hom}(E, F))$, with respect to the diagonal, **supported within the diagonal**, Δ . i.e.

$$\text{Diff}^m(Z, E, F) \iff I^m_{\Delta}(Z^2, \Delta; \text{Hom}(E, F))$$

Projective differential operators/ Ψ DOs

The previous facts motivates our definition of projective differential and pseudodifferential operators when E and F are only projective vector bundles associated to a fixed finite-dimensional Azumaya bundle \mathcal{A} .

Projective differential operators/ Ψ DOs

The previous facts motivates our definition of projective differential and pseudodifferential operators when E and F are only projective vector bundles associated to a fixed finite-dimensional Azumaya bundle \mathcal{A} .

Since a projective vector bundle E is not global on Z , one **cannot** make sense of sections of E , let alone operators acting between sections! However, it still makes sense to talk about Schwartz kernels even in this case, as we explain.

Notice that $\text{Hom}(E, F) = F \boxtimes E^*$ is a projective bundle on Z^2 associated to the Azumaya bundle, $\mathcal{A}_L \boxtimes \mathcal{A}'_R$.

Notice that $\text{Hom}(E, F) = F \boxtimes E^*$ is a projective bundle on Z^2 associated to the Azumaya bundle, $\mathcal{A}_L \boxtimes \mathcal{A}'_R$.

The restriction $\Delta^* \text{Hom}(E, F) = \text{hom}(E, F)$ to the diagonal is an ordinary vector bundle, it is therefore reasonable to expect that $\text{Hom}(E, F)$ also restricts to an ordinary vector bundle in a tubular nbd N_ϵ of the diagonal.

Notice that $\text{Hom}(E, F) = F \boxtimes E^*$ is a projective bundle on Z^2 associated to the Azumaya bundle, $\mathcal{A}_L \boxtimes \mathcal{A}'_R$.

The restriction $\Delta^* \text{Hom}(E, F) = \text{hom}(E, F)$ to the diagonal is an ordinary vector bundle, it is therefore reasonable to expect that $\text{Hom}(E, F)$ also restricts to an ordinary vector bundle in a tubular nbd N_ϵ of the diagonal.

In [MMS3], it is shown that there is a **canonical such choice**, $\text{Hom}^{\mathcal{A}}(E, F)$ of such that the **composition properties hold**.

This allows us to **define** the space of **projective pseudo-differential operators** $\Psi_\epsilon^\bullet(Z; E, F)$ with Schwartz kernels supported in an ϵ -neighborhood N_ϵ of the diagonal Δ in Z^2 , with the space of **conormal distributions**, $I_\epsilon^\bullet(N_\epsilon, \Delta; \text{Hom}^A(E, F))$.

$$\Psi_\epsilon^\bullet(Z; E, F) \quad := \quad I_\epsilon^\bullet(N_\epsilon, \Delta; \text{Hom}^A(E, F)).$$

This allows us to **define** the space of **projective pseudo-differential operators** $\Psi_\epsilon^\bullet(Z; E, F)$ with Schwartz kernels supported in an ϵ -neighborhood N_ϵ of the diagonal Δ in Z^2 , with the space of **conormal distributions**, $I_\epsilon^\bullet(N_\epsilon, \Delta; \text{Hom}^A(E, F))$.

$$\Psi_\epsilon^\bullet(Z; E, F) \quad := \quad I_\epsilon^\bullet(N_\epsilon, \Delta; \text{Hom}^A(E, F)).$$

Despite **not** being a space of operators, this has precisely the **same local structure** as in the standard case and has similar composition properties provided supports are restricted to appropriate neighbourhoods of the diagonal.

The space of **projective smoothing operators**, $\Psi_\epsilon^{-\infty}(Z; E, F)$ is defined as the smooth sections, $C_c^\infty(N_\epsilon; \text{Hom}^A(E, F) \otimes \pi_R^* \Omega)$.

The space of **projective smoothing operators**, $\Psi_\epsilon^{-\infty}(Z; E, F)$ is defined as the smooth sections, $C_c^\infty(N_\epsilon; \text{Hom}^A(E, F) \otimes \pi_R^* \Omega)$.

The space of all **projective differential operators**, $\text{Diff}^\bullet(Z; E, F)$ is defined as those conormal distributions that are **supported within the diagonal** Δ in Z^2 ,

$$\text{Diff}^\bullet(Z; E, F) := I_\Delta^\bullet(N_\epsilon, \Delta; \text{Hom}^A(E, F)).$$

In fact, $\text{Diff}^\bullet(Z; E, F)$ is even a **ring** when $E = F$.

Projective spin Dirac operator

Recall that there is a projective bundle of spinors $\mathcal{S} = \mathcal{S}^+ \oplus \mathcal{S}^-$ on any even dimensional oriented manifold Z .

Projective spin Dirac operator

Recall that there is a projective bundle of spinors $\mathcal{S} = \mathcal{S}^+ \oplus \mathcal{S}^-$ on any even dimensional oriented manifold Z .

There are natural spin connections on the Clifford algebra bundle $Cl(Z)$ and \mathcal{S}^\pm induced from the Levi-Civita connection on T^*Z .

Projective spin Dirac operator

Recall that there is a projective bundle of spinors $\mathcal{S} = \mathcal{S}^+ \oplus \mathcal{S}^-$ on any even dimensional oriented manifold Z .

There are natural spin connections on the Clifford algebra bundle $Cl(Z)$ and \mathcal{S}^\pm induced from the Levi-Civita connection on T^*Z .

Recall also that $\text{hom}(\mathcal{S}, \mathcal{S}) \cong Cl(Z)$, has an extension to $\tilde{Cl}(Z)$ in a tubular neighbourhood of the diagonal Δ , with an induced connection ∇ .

The **projective spin Dirac operator** is defined as the distributional section

$$\not{D} = c/ \cdot \nabla_L(\kappa_{Id}), \quad \kappa_{Id} = \delta(z - z') Id_S$$

Here ∇_L is the connection ∇ restricted to the left variables with $c/$ the contraction given by the Clifford action of T^*Z on the left. As in the usual case, the projective spin Dirac operator \not{D} is **elliptic** and odd wrt \mathbb{Z}_2 grading of S .

The **principal symbol map** is well defined for conormal distributions, leading to the globally defined symbol map,

$$\sigma : \Psi_{\epsilon}^m(Z; E, F) \longrightarrow C^{\infty}(T^*Z, \pi^* \text{hom}(E, F)),$$

homogeneous of degree m ; here $\text{hom}(E, F)$, is a globally defined, **ordinary vector bundle** with fibre, $\text{hom}(E, F)_z = F_z \otimes E_z^*$. Thus **ellipticity** is well defined, as the **invertibility of this symbol**.

The **principal symbol map** is well defined for conormal distributions, leading to the globally defined symbol map,

$$\sigma : \Psi_{\epsilon}^m(Z; E, F) \longrightarrow C^{\infty}(T^*Z, \pi^* \text{hom}(E, F)),$$

homogeneous of degree m ; here $\text{hom}(E, F)$, is a globally defined, **ordinary vector bundle** with fibre, $\text{hom}(E, F)_z = F_z \otimes E_z^*$. Thus **ellipticity** is well defined, as the **invertibility of this symbol**.

Equivalently, $A \in \Psi_{\epsilon/2}^m(Z; E, F)$ is **elliptic** if there exists a parametrix $B \in \Psi_{\epsilon/2}^{-m}(Z; F, E)$ and smoothing operators $Q_R \in \Psi_{\epsilon}^{-\infty}(Z; E, E)$, $Q_L \in \Psi_{\epsilon}^{-\infty}(Z; F, F)$ such that

$$BA = I_E - Q_R, \quad AB = I_F - Q_L$$

The **trace functional** is defined on projective smoothing operators $\text{Tr} : \Psi_{\epsilon}^{-\infty}(Z; E) \rightarrow \mathbb{C}$ as

$$\text{Tr}(Q) = \int_Z \text{tr } Q(z, z).$$

The **trace functional** is defined on projective smoothing operators $\text{Tr} : \Psi_{\epsilon}^{-\infty}(Z; E) \rightarrow \mathbb{C}$ as

$$\text{Tr}(Q) = \int_Z \text{tr } Q(z, z).$$

It vanishes on commutators, i.e. $\text{Tr}(QR - RQ) = 0$, if $Q \in \Psi_{\epsilon/2}^{-\infty}(Z; F, E)$, $R \in \Psi_{\epsilon/2}^{-\infty}(Z; E, F)$ which follows from Fubini's theorem.

The **trace functional** is defined on projective smoothing operators $\text{Tr} : \Psi_{\epsilon}^{-\infty}(Z; E) \rightarrow \mathbb{C}$ as

$$\text{Tr}(Q) = \int_Z \text{tr} Q(z, z).$$

It vanishes on commutators, i.e. $\text{Tr}(QR - RQ) = 0$, if $Q \in \Psi_{\epsilon/2}^{-\infty}(Z; F, E)$, $R \in \Psi_{\epsilon/2}^{-\infty}(Z; E, F)$ which follows from Fubini's theorem.

The **fractional analytic index** of the projective elliptic operator $A \in \Psi_{\epsilon}^{\bullet}(Z; E, F)$ is defined in the essentially analytic way as,

$$\text{Index}_a(A) = \text{Tr}([A, B]) \in \mathbb{R}$$

where B is a parametrix for A , and the RHS is the notation for $\text{Tr}_F(AB - I_F) - \text{Tr}_E(BA - I_E)$.

Techniques to prove basic properties of the fractional analytic index

For $A \in \Psi_{\epsilon/4}^m(Z; E, F)$, the Guillemin-Wodzicki **residue trace** is,

$$\mathrm{Tr}_R(A) = \lim_{z \rightarrow 0} z \mathrm{Tr}(AD(z))$$

where $D(z) \in \Psi_{\epsilon/4}^z(Z; E)$ is an **entire family** of Ψ DOs of complex order z which is **elliptic** and such that $D(0) = I$. The residue trace is **independent of the choice** of such a family.

Techniques to prove basic properties of the fractional analytic index

For $A \in \Psi_{\epsilon/4}^m(Z; E, F)$, the Guillemin-Wodzicki **residue trace** is,

$$\mathrm{Tr}_R(A) = \lim_{z \rightarrow 0} z \mathrm{Tr}(AD(z))$$

where $D(z) \in \Psi_{\epsilon/4}^z(Z; E)$ is an **entire family** of Ψ DOs of complex order z which is **elliptic** and such that $D(0) = I$. The residue trace is **independent of the choice** of such a family.

- 1 The residue trace Tr_R **vanishes** on all Ψ DOs of sufficiently negative order.

For $A \in \Psi_{\epsilon/4}^m(Z; E, F)$, the Guillemin-Wodzicki **residue trace** is,

$$\mathrm{Tr}_R(A) = \lim_{z \rightarrow 0} z \mathrm{Tr}(AD(z))$$

where $D(z) \in \Psi_{\epsilon/4}^z(Z; E)$ is an **entire family** of Ψ DOs of complex order z which is **elliptic** and such that $D(0) = I$. The residue trace is **independent of the choice** of such a family.

- 1 The residue trace Tr_R **vanishes** on all Ψ DOs of sufficiently negative order.
- 2 The residue trace Tr_R is also a **trace functional**, that is,

$$\mathrm{Tr}_R([A, B]) = 0,$$

for $A \in \Psi_{\epsilon/4}^m(Z; E, F)$, $B \in \Psi_{\epsilon/4}^{m'}(Z; F, E)$.

The **regularized trace**, is defined to be the residue,

$$\mathrm{Tr}_D(A) = \lim_{z \rightarrow 0} \frac{1}{z} (z \mathrm{Tr}(AD(z)) - \mathrm{Tr}_R(A)).$$

The **regularized trace**, is defined to be the residue,

$$\mathrm{Tr}_D(A) = \lim_{z \rightarrow 0} \frac{1}{z} (z \mathrm{Tr}(AD(z)) - \mathrm{Tr}_R(A)).$$

For general A , Tr_D **does** depend on the regularizing family $D(z)$.

The **regularized trace**, is defined to be the residue,

$$\mathrm{Tr}_D(A) = \lim_{z \rightarrow 0} \frac{1}{z} (z \mathrm{Tr}(AD(z)) - \mathrm{Tr}_R(A)).$$

For general A , Tr_D **does** depend on the regularizing family $D(z)$.
But for smoothing operators it **coincides** with the standard operator trace,

$$\mathrm{Tr}_D(S) = \mathrm{Tr}(S), \quad \forall S \in \Psi_\epsilon^{-\infty}(Z, E)$$

The **regularized trace**, is defined to be the residue,

$$\mathrm{Tr}_D(A) = \lim_{z \rightarrow 0} \frac{1}{z} (z \mathrm{Tr}(AD(z)) - \mathrm{Tr}_R(A)).$$

For general A , Tr_D **does** depend on the regularizing family $D(z)$.
But for smoothing operators it **coincides** with the standard operator trace,

$$\mathrm{Tr}_D(S) = \mathrm{Tr}(S), \quad \forall S \in \Psi_\epsilon^{-\infty}(Z, E)$$

Therefore the fractional analytic index is also given by,

$$\mathrm{Index}_a(A) = \mathrm{Tr}_D([A, B])$$

for a projective elliptic operator A , and B a parametrix for A .

The regularized trace Tr_D is **not** a trace function, but however it satisfies the **'trace defect formula'**,

$$\text{Tr}_D([A, B]) = \text{Tr}_R(B\delta_D A)$$

where δ_D is a **derivation** acting on the full symbol algebra.

The regularized trace Tr_D is **not** a trace function, but however it satisfies the '**trace defect formula**',

$$\text{Tr}_D([A, B]) = \text{Tr}_R(B\delta_D A)$$

where δ_D is a **derivation** acting on the full symbol algebra. It also satisfies the condition of being closed,

$$\text{Tr}_R(\delta_D a) = 0 \quad \forall a.$$

Using the derivation δ_D and the trace defect formula, we prove:

Using the derivation δ_D and the trace defect formula, we prove:

- 1 the **homotopy invariance** of the index,

$$\frac{d}{dt} \text{Index}_a(A_t) = 0,$$

where $t \mapsto A_t$ is a smooth 1-parameter family of projective elliptic Ψ DOs;

Using the derivation δ_D and the trace defect formula, we prove:

- ① the **homotopy invariance** of the index,

$$\frac{d}{dt} \text{Index}_a(A_t) = 0,$$

where $t \mapsto A_t$ is a smooth 1-parameter family of projective elliptic Ψ DOs;

- ② the **multiplicativity property** of of the index,

$$\text{Index}_a(A_2 A_1) = \text{Index}_a(A_1) + \text{Index}_a(A_2),$$

where A_i for $i = 1, 2$ are projective elliptic Ψ DOs.

An analogue of the **McKean-Singer formula** holds,

$$\text{Index}_a(\not\partial_E^+) = \lim_{t \downarrow 0} \text{Tr}_s(H_\chi(t))$$

where $H_\chi(t) = \chi(H_t)$ is a globally defined, **truncated heat kernel**, both in space (in a nbd of the diagonal) and in time.

An analogue of the **McKean-Singer formula** holds,

$$\text{Index}_a(\not{D}_E^+) = \lim_{t \downarrow 0} \text{Tr}_s(H_\chi(t))$$

where $H_\chi(t) = \chi(H_t)$ is a globally defined, **truncated heat kernel**, both in space (in a nbd of the diagonal) and in time.

The local index theorem can then be applied, thanks to the McKean-Singer formula, to obtain the **index theorem** for projective spin Dirac operators.

Index of projective spin Dirac operators

Theorem ([MMS3])

The projective spin Dirac operator on an even-dimensional compact oriented manifold Z , has fractional analytic index,

$$\text{Index}_a(\not{D}^+) = \int_Z \widehat{A}(Z) \in \mathbb{Q}.$$

Index of projective spin Dirac operators

Theorem ([MMS3])

The projective spin Dirac operator on an even-dimensional compact oriented manifold Z , has fractional analytic index,

$$\text{Index}_a(\not{D}^+) = \int_Z \widehat{A}(Z) \in \mathbb{Q}.$$

Recall that $Z = \mathbb{C}P^{2n}$ is an oriented but **non-spin** manifold such that $\int_Z \widehat{A}(Z) \notin \mathbb{Z}$, justifying the title of the talk. e.g.

$$Z = \mathbb{C}P^2 \quad \implies \quad \text{Index}_a(\not{D}^+) = -1/8.$$

$$Z = \mathbb{C}P^4 \quad \implies \quad \text{Index}_a(\not{D}^+) = 3/128.$$

Equivariant transversally elliptic Dirac operator

Recall that projective half spinor bundles \mathcal{S}^\pm on Z can be realized as $Spin(n)$ -equivariant honest vector bundles, $(\tilde{\mathcal{S}}^+, \tilde{\mathcal{S}}^-)$, over the total space of the oriented frame bundle \mathcal{P} and in which the center, \mathbb{Z}_2 , acts as ± 1 , as follows:

Equivariant transversally elliptic Dirac operator

Recall that projective half spinor bundles \mathcal{S}^\pm on Z can be realized as $Spin(n)$ -equivariant honest vector bundles, $(\tilde{\mathcal{S}}^+, \tilde{\mathcal{S}}^-)$, over the total space of the oriented frame bundle \mathcal{P} and in which the center, \mathbb{Z}_2 , acts as ± 1 , as follows:
the conormal bundle N to the fibres of \mathcal{P} has vanishing w_2 -obstruction, and $\tilde{\mathcal{S}}^\pm$ are just the 1/2 spin bundles of N .

Equivariant transversally elliptic Dirac operator

Recall that projective half spinor bundles S^\pm on Z can be realized as $Spin(n)$ -equivariant honest vector bundles, $(\tilde{S}^+, \tilde{S}^-)$, over the total space of the oriented frame bundle \mathcal{P} and in which the center, \mathbb{Z}_2 , acts as ± 1 , as follows: the conormal bundle N to the fibres of \mathcal{P} has vanishing w_2 -obstruction, and \tilde{S}^\pm are just the 1/2 spin bundles of N .

One can define the $Spin(n)$ -equivariant transversally elliptic Dirac operator $\tilde{\not{D}}^\pm$ using the Levi-Civita connection on Z together with the Clifford contraction, where **transverse ellipticity** means that the principal symbol is invertible when restricted to directions that are conormal to the fibres.

The nullspaces of $\tilde{\mathcal{D}}^\pm$ are infinite dimensional unitary representations of $Spin(n)$. The transverse ellipticity implies that the characters of these representations are **distributions** on the group $Spin(n)$. In particular, the **multiplicity** of each irreducible unitary representation in these nullspaces is **finite**, and grows at most polynomially.

The nullspaces of $\tilde{\not{D}}^\pm$ are infinite dimensional unitary representations of $Spin(n)$. The transverse ellipticity implies that the characters of these representations are **distributions** on the group $Spin(n)$. In particular, the **multiplicity** of each irreducible unitary representation in these nullspaces is **finite**, and grows at most polynomially.

The **$Spin(n)$ -equivariant index** of $\tilde{\not{D}}^+$ is defined to be the following distribution on $Spin(n)$,

$$\text{Index}_{Spin(n)}(\tilde{\not{D}}^+) = \text{Char}(\text{Nullspace}(\tilde{\not{D}}^+)) - \text{Char}(\text{Nullspace}(\tilde{\not{D}}^-))$$

An alternate, **analytic** description of the $Spin(n)$ -equivariant index of $\tilde{\mathcal{D}}^+$ is: for a function of compact support $\chi \in C_c^\infty(G)$, the action of the group induces a graded operator

$$T_\chi : C^\infty(\mathcal{P}; \tilde{\mathcal{S}}) \longrightarrow C^\infty(\mathcal{P}; \tilde{\mathcal{S}}), \quad T_\chi u(x) = \int_G \chi(g) g^* u dg,$$

which is smoothing along the fibres.

An alternate, **analytic** description of the $Spin(n)$ -equivariant index of $\tilde{\partial}^+$ is: for a function of compact support $\chi \in C_c^\infty(G)$, the action of the group induces a graded operator

$$T_\chi : C^\infty(\mathcal{P}; \tilde{\mathcal{S}}) \longrightarrow C^\infty(\mathcal{P}; \tilde{\mathcal{S}}), \quad T_\chi u(x) = \int_G \chi(g) g^* u dg,$$

which is smoothing along the fibres. $\tilde{\partial}^+$ has a microlocal parametrix Q , in the directions that are conormal to the fibres (ie along N). Then for any $\chi \in C_c^\infty(G)$,

$$T_\chi \circ (\tilde{\partial}^+ \circ Q - I_-) \in \Psi^{-\infty}(\mathcal{P}; \tilde{\mathcal{S}}^-); \quad T_\chi \circ (Q \circ \tilde{\partial}^+ - I_+) \in \Psi^{-\infty}(\mathcal{P}; \tilde{\mathcal{S}}^+)$$

are smoothing operators. The **$Spin(n)$ -equivariant index** of $\tilde{\partial}^+$, evaluated at $\chi \in C_c^\infty(G)$, is also given by:

An alternate, **analytic** description of the $Spin(n)$ -equivariant index of $\tilde{\mathcal{D}}^+$ is: for a function of compact support $\chi \in C_c^\infty(G)$, the action of the group induces a graded operator

$$T_\chi : C^\infty(\mathcal{P}; \tilde{\mathcal{S}}) \longrightarrow C^\infty(\mathcal{P}; \tilde{\mathcal{S}}), \quad T_\chi u(x) = \int_G \chi(g) g^* u dg,$$

which is smoothing along the fibres. $\tilde{\mathcal{D}}^+$ has a microlocal parametrix Q , in the directions that are conormal to the fibres (ie along N). Then for any $\chi \in C_c^\infty(G)$,

$$T_\chi \circ (\tilde{\mathcal{D}}^+ \circ Q - I_-) \in \Psi^{-\infty}(\mathcal{P}; \tilde{\mathcal{S}}^-); \quad T_\chi \circ (Q \circ \tilde{\mathcal{D}}^+ - I_+) \in \Psi^{-\infty}(\mathcal{P}; \tilde{\mathcal{S}}^+)$$

are smoothing operators. The **$Spin(n)$ -equivariant index** of $\tilde{\mathcal{D}}^+$, evaluated at $\chi \in C_c^\infty(G)$, is also given by:

$$\text{Index}_{Spin(n)}(\tilde{\mathcal{D}}^+)(\chi) = \text{Tr}(T_\chi \circ (\tilde{\mathcal{D}}^+ \circ Q - I_-)) - \text{Tr}(T_\chi \circ (Q \circ \tilde{\mathcal{D}}^+ - I_+))$$

Relation between the two types of Dirac operators

Theorem ([MMS4])

Let $\pi : \mathcal{P}^2 \rightarrow Z^2$ denote the projection. The pushforward map, π_ , maps the Schwartz kernel of the $\text{Spin}(n)$ -transversally elliptic Dirac operator to the projective Dirac operator: That is,*

$$\pi_*(\tilde{\mathcal{D}}^\pm) = \mathcal{D}^\pm.$$

Relation between the two solutions

We will now relate these two pictures.

An easy argument shows that the support of the equivariant index distribution is contained within the center \mathbb{Z}_2 of $Spin(n)$.

Theorem

Let $\phi \in C^\infty(Spin(n))$ be such that :

- 1 $\phi \equiv 1$ in a neighborhood of e ;
- 2 $-1 \notin \text{supp}(\phi)$. Then

$$\text{Index}_{Spin(n)}(\tilde{\mathcal{D}}^+)(\phi) = \text{Index}_a(\mathcal{D}^+)$$

Relation between the two solutions

We will now relate these two pictures.

An easy argument shows that the support of the equivariant index distribution is contained within the center \mathbb{Z}_2 of $Spin(n)$.

Theorem

Let $\phi \in C^\infty(Spin(n))$ be such that :

- 1 $\phi \equiv 1$ in a neighborhood of e ;
- 2 $-1 \notin \text{supp}(\phi)$. Then

$$\text{Index}_{Spin(n)}(\tilde{\mathcal{D}}^+)(\phi) = \text{Index}_a(\mathcal{D}^+)$$

Informally, the **fractional analytic index**, of the projective Dirac operator \mathcal{D}^+ , is the **coefficient of the delta function (distribution) at the identity** in $Spin(n)$ of the **$Spin(n)$ -equivariant index** for the associated transversally elliptic Dirac operator $\tilde{\mathcal{D}}^+$ on \mathcal{P} .