

have a rep on 0-dim'l v.space — acts as identity

11/12/02

Defn: A rep (ρ, V) of a group G is indecomposable if it's not equivalent to a direct sum of 2 nontrivial reps unless one of those reps is 0-dim'l. (rule out case that we write something as a sum of 0 w/ itself) exactly

Note: $(\rho, V) \cong (\rho, V) \oplus (\rho_{\{0\}}, \{0\})$

Defn: A rep (ρ, V) of G has a rep (σ, U) as a subrepresentation if $U \subseteq V$ and

$\rho(g): U \rightarrow U \quad \forall g$ so that we can define

$$\sigma(g) = \rho(g)|_U : U \rightarrow U.$$

Defn: A rep (ρ, V) is irreducible iff it has no subrepresentations except $(\rho|_{\{0\}}, \{0\})$

and (ρ, V) itself, (which must be different)

Note: We don't want the zero rep to be indecomposable or irreducible.

(Just as 1 isn't prime.) (If it were — wouldn't have unique decomposition).

So — we exclude the zero-dim'l rep from being indecomposable or irreducible.

Note: irreducible \Rightarrow indecomposable

Thm: If (ρ, \mathfrak{H}) is a unitary rep, ρ is indecomposable iff ρ is irreducible.

proof: In general, ρ irreducible $\Rightarrow \rho$ indecomposable.

For converse — we use the fact ρ is unitary:

Suppose ρ has a subrep $(\rho|_U, U)$ (other than $\{0\}$ or V) and show that

U^\perp (all vectors \perp to vectors in U) is preserved by $\rho(g) \forall g \in G$. We'll show $V = U \oplus U^\perp$.

$$v \in U^\perp \Rightarrow \langle v, w \rangle = 0 \quad \forall w \in U.$$

Now, consider $\langle \rho(g)v, w \rangle = \langle v, \rho(g)^{-1}w \rangle = \langle v, \underbrace{\rho(g^{-1})w}_{\text{this is in } U \text{ since } U \text{ is a subrep.}} \rangle$

since ρ is unitary,
so $\rho(g)^* = \rho(g)^{-1}$

this is in U
since U is a
subrep.

$$= 0.$$

Thus, $\rho(g)v \in U^\perp$.

Therefore — $(\rho|_{U^\perp}, U^\perp)$ is also a subrep.

of (ρ, V) and so $(\rho, V) \cong (\rho|_{U^\perp}, U^\perp) \oplus (\rho|_U, U)$

Want an example of indecomposable but not irreducible.
(can't be unitary!)

Example: What's a rep of $(\mathbb{R}, +)$ which is indecomposable but not irreducible?

$$\rho(t) = \begin{pmatrix} 1 & t \\ 0 & 1 \end{pmatrix} \text{ is a rep of } \mathbb{R} \text{ on } \mathbb{C}^2.$$

(upper triangular matrix won't be unitary)

Note: This matrix isn't unitary for any inner product.

The subspace $U = \left\{ \begin{pmatrix} a \\ 0 \end{pmatrix} \mid a \in \mathbb{C} \right\} \subseteq \mathbb{C}^2$

is preserved by $\rho(t) \forall t$, so it's a subrep, and thus our rep is reducible.

(trivial rep - $\rho(t)$ is identity)

(neither $\{0\}$ nor \mathbb{C}^2)

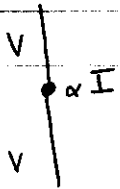
But it's not decomposable. In particular $\nexists U' \subseteq \mathbb{C}^2$ which is also preserved by $\rho(t) \forall t$ and st $\mathbb{C}^2 \cong U \oplus U'$.

Just need to check that the only subreps are $\{0\}$, U , and \mathbb{C}^2 . \square

ρ_k are all 1-dim'l reps on \mathbb{R} .

recall: $\rho_k(t) = e^{kt}$

Schur's Lemma: If a rep (ρ, V) is irreducible and $T: V \rightarrow V$ is an intertwiner then $T = \alpha I_V$ for some $\alpha \in \mathbb{C}$.



proof: Suppose T is an intertwiner. Then any eigenspace

$$U = \{v \in V \mid Tv = \lambda v\} \text{ is a subrep.}$$

Let $v \in U$. Then $Tv = \lambda v$. T an eigenvector

$$\Rightarrow T\rho(g)v = \rho(g)Tv = \lambda\rho(g)v$$

because T an intertwiner

Thus, $\rho(g)v \in U$.

But since U must be $\{0\}$ or V , T has at most one eigenvalue, so $T = \alpha I$. \square

Corollary: If G is abelian, then all its irreps (irreducible reps) are 1-dim'l.

unitary means $|\alpha| = 1$.

pf: Suppose (ρ, V) is an irrep of G .

Since G is abelian:

$$\rho(g)\rho(g') = \rho(gg') = \rho(g'g) = \rho(g')\rho(g) \quad \forall g, g' \in G$$

(a homomorphism maps an abelian grp into an abel. subgroup)

But— we call an intertwiner something that commutes w/ $\rho(g) \forall g$.

So— $\rho(g)$ is an intertwiner $\forall g \in G$.

By Schur's lemma, $\rho(g) = \alpha_g I \quad \forall g \in G$.
(multiple depends on what g is)

Then— $T: V \rightarrow V$ is an intertwiner if T commutes w/ $\rho(g)$ — that is T is an intertwiner for all $T: V \rightarrow V$ because $\rho(g)$ is a mult of identity.

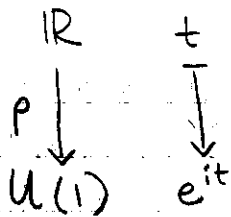
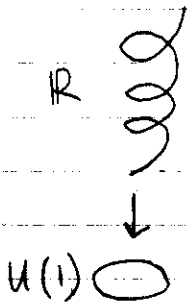
By Schur's lemma, all $T: V \rightarrow V$ must be multiples of I . So— V must be 0-dim'l or 1-dim'l. But, 0-dim'l ones aren't irreducible, so since V is irred. it must be 1-dim'l. \square

Corollary: Every (smooth) irrep of \mathbb{R} is of the form

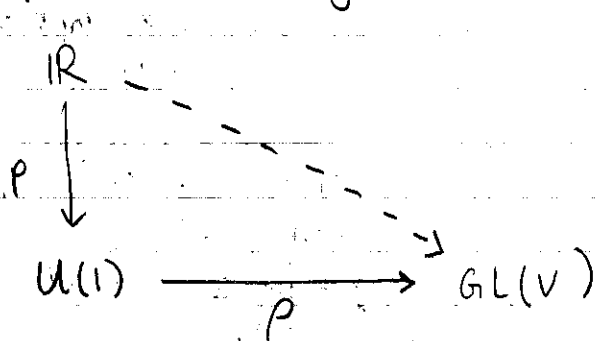
$$\rho_k(t) = e^{kt}$$

pf: We know all the 1-dim'l (smooth) reps are of this form; since they're 1-dim'l they're irreducible and there are no higher-dim'l irreps by previous corollary. \square

How about $U(1)$? We have an onto homo

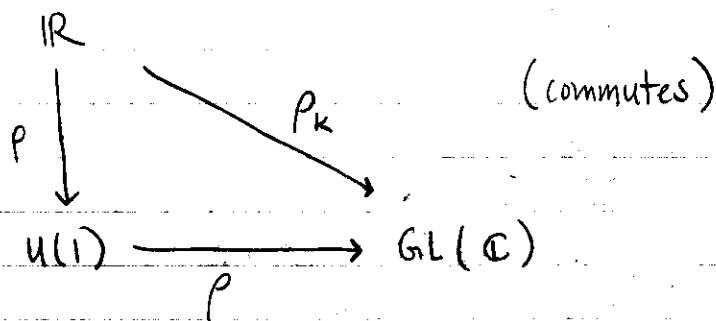


so any rep of $U(1)$ gives a rep of IR :



and any irrep of $U(1)$ must give an irrep of IR .

So - the only options for ρ are:



ie. $\rho(e^{it}) = \rho_k(t) = e^{kt}$ $t \in \mathbb{R}$
 $k \in \mathbb{C}$

Since $\rho(e^{2\pi i}) = 1$ we need

$e^{k \cdot 2\pi i} = 1$ so $ik \in \mathbb{Z}$

reps of \mathbb{Z} form grp $U(1)$ (by Poincaré duality)

Thm: All irreps of $U(1)$ are of the form

$$\rho(e^{it}) = e^{int} \quad \text{for some } n \in \mathbb{Z}.$$

Note: $U(1), \mathbb{R}$ are 'boring' since both are 1-dim'l & abelian.

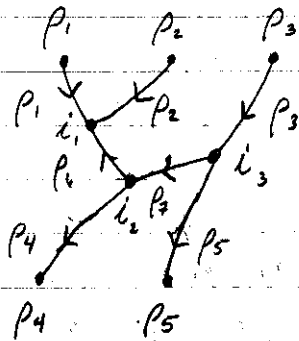
Compact Lie grps	Noncompact Lie grps
countably many smooth irreps, all finite-dim'l	typically uncountably many (smooth) irreps,
all are unitary w/r/t some inner product	typically some infinite-dim'l, typically many can't be made unitary.
e.g.: $\rho(e^{it}) = e^{int} \quad n \in \mathbb{Z}$ $U(1)$	e.g.: $\rho(t) = e^{kt}, k \in \mathbb{C}$ \mathbb{R}

11/14/02

Let G be a (Lie) group. There's a category $\text{Net}(G)$ where the objects are lists of irreducible (smooth) unitary reps of G :

$\rho_1 \quad \rho_2 \quad \rho_3 \quad \rho_4$
 $\bullet \quad \bullet \quad \bullet \quad \bullet$

The morphisms are (certain equiv. classes) of diagrams like:



vertices of the graph are the ones in the middle, not top or bottom.

label vertices w/ i 's: intertwiners

$$i_1: \rho_1 \otimes \rho_2 \otimes \rho_6 \longrightarrow \mathbb{C}$$

↑
no arrows going out!

$$i_2: \rho_7 \longrightarrow \rho_6 \otimes \rho_4$$

$$i_3: \rho_5 \longrightarrow \rho_7 \otimes \rho_5$$

Note: edges are labelled by unitary irreps;
 vertices labelled by intertwiners.

Edges are directed; changing direction and replacing rep by its dual doesn't change diagram.

Incoming & outgoing edges, if directed downwards, must be labelled by reps that match those on corresponding dots at top & bottom.

We get a braided monoidal category w/ duals.

symmetric

In quantum field theory we call a morphism in $\text{Net}(G)$ a "Feynman diagram", edges represent "particles" and vertices represent "interactions."

* We take a diagram (like on prev pg) and compose & tensor the intertwiners i 's and get a big intertwiner from $\rho_1 \otimes \rho_2 \otimes \rho_3 \rightarrow \rho_4 \otimes \rho_5$.

What does it mean to "evaluate" a Feynman diagram?

We take one and compose/tensor all the intertwiners to get a new intertwiner:

$$i: \begin{array}{ccc} (\rho, v) & \longrightarrow & (\rho', v') \\ \rho & \longrightarrow & \rho' \end{array}$$

(linear map bet. the v . spaces)

where ρ is the tensor product of all the irreps labelling the incoming edges; ρ' is the tensor product of reps labelling outgoing edges.

So, we turn objects in $\text{Net}(G)$ into reps
 morphisms in $\text{Net}(G)$ into intertwiners

That is - we get a functor:

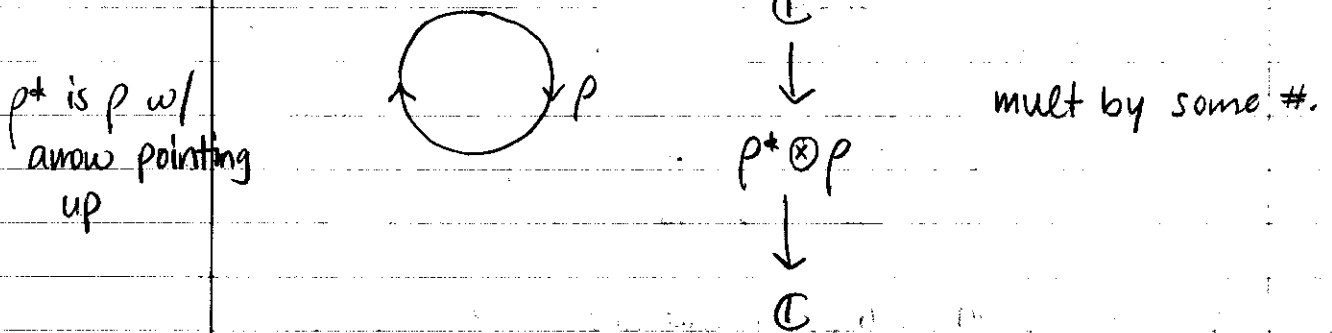
$$\text{eval}: \text{Net}(G) \longrightarrow \text{Rep}(G)$$

where $\text{Rep}(G)$ is the category whose objects are unitary (smooth) reps of G and whose morphisms are intertwiners.

Note: Can't always compose / tensor intertwiners and get a new one - we need finite dim'l.

* Actually this "eval" is well-defined if we restrict attention to finite-dim'l unitary reps when defining $\text{Net}(G)$ and $\text{Rep}(G)$.

Otherwise the evaluation of a diagram may "diverge".



The linear map from \mathbb{C} to \mathbb{C} is ill-defined if dim of our rep is infinite.

↙ We're talking about unitary reps!
(Hilbert spaces)

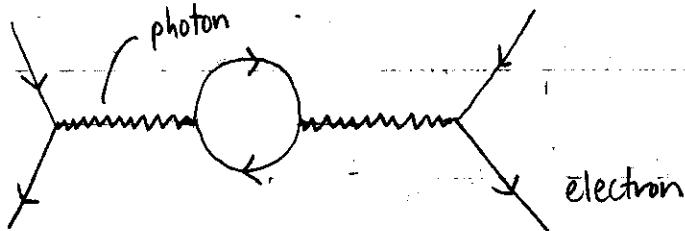
Say our rep is (ρ, H) and Hilbert space H has orthonormal basis e_i and H^* has dual basis e^i . Then:

$$\begin{array}{ccc}
 \mathbb{C} & \xrightarrow{\quad} & \mathbb{C} \\
 \downarrow & & \downarrow \\
 H^* \otimes H & \xrightarrow{\quad} & \sum e^i \otimes e_i \\
 \downarrow & & \downarrow \\
 \mathbb{C} & \xrightarrow{\quad} & \sum_i e^i(e_i) = \sum_i 1 = \dim(H)
 \end{array}$$

since i is running over a basis.

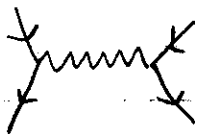
When H is infinite dim'l, this blows up.

In general — whenever you allow infinite dim'l reps to label edges of Feynman diagrams which are not simply-connected, you run the risk of having the evaluation diverge:

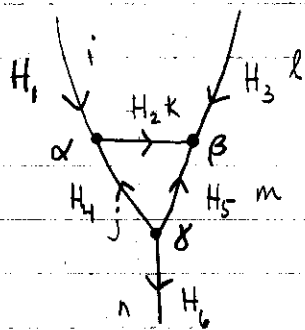


2 electrons exchange a photon

But simply-connected diagrams are okay:



Example:



$$\alpha: H_1 \otimes H_4 \rightarrow H_2$$

$$\beta: H_3 \otimes H_2 \otimes H_5 \rightarrow \mathbb{C}$$

$$\gamma: \mathbb{C} \rightarrow H_4 \otimes H_5 \otimes H_6$$

want to put these together to get something from

$$I: H_1 \otimes H_3$$



$$H_6$$

α_{ij}^k is a matrix elt. for i running over basis of H_1
 j " " " " H_4
 k " " " " H_2

Similarly we get: β^{klm} and γ_{jmn}

Then we multiply matrices: actually tensors

$$\sum_{k,j,m} \alpha_{ij}^k \beta^{klm} \gamma_{jmn} = I_n$$

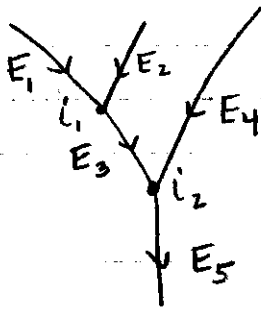
* the edges we're summing over are the "inside" ones.

The answer is basis-independent!

Some examples:

- 1) $G = \mathbb{R}$. Then unitary irreps. are 1-dim'l:
 $\rho_E(t) = e^{iEt} \quad t \in \mathbb{R}, E \in \mathbb{R}$

We call E the "energy" and our diagrams are "energy networks":



For there to be an intertwiner at each vertex, we need "conservation of energy"

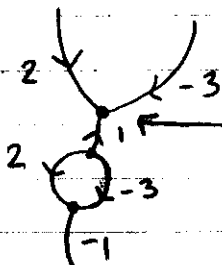
$$\underbrace{E_1 + E_2}_{\text{in}} = \underbrace{E_3}_{\text{out}} \text{ for } i_1$$

* We think of \mathbb{R} as time translation, and the conserved quantity is energy.

- 2) $G = U(1)$ "gauge grp for electromagnetism"
 ("phase changes" in electromagnetism)

Then the unitary irreps $\rho_\gamma(e^{it}) = e^{i\gamma t} \quad e^{it} \in U(1)$
 $\gamma \in \mathbb{Z}$

We call γ the "charge" and our diagrams are "charge networks":



now we need "conservation of charge."

$$3) G = \mathbb{R}^4$$

↑
spacetime translations

Here irreps are labelled by 4 numbers:

$$\left(\underset{E}{\parallel}, p_x, p_y, p_z \right) \in \mathbb{R}^4$$

p_t

energy

momentum in x, y, z direction

need

$$p = p' + p''$$

for there to

be an intertwiner

- conservation of

"energy-momentum"

