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Suppose (ρ, V) is a rep of G on V . Define the "symmetrizer"

$P_S: V^{\otimes n} \rightarrow V^{\otimes n}$ by

$$P_S(v_1 \otimes \dots \otimes v_n) = \frac{1}{n!} \sum_{\sigma \in S_n} v_{\sigma(1)} \otimes \dots \otimes v_{\sigma(n)} \quad v_i \in V$$

with $P_S^2 = P_S$. (check)

Check: P_S is an intertwiner:

$(\rho^{\otimes n}, V^{\otimes n})$ is a rep of G with

$$\rho^{\otimes n}(g)(v_1 \otimes \dots \otimes v_n) = \rho(g)v_1 \otimes \dots \otimes \rho(g)v_n.$$

and

$$P_S \rho^{\otimes n}(g) = \rho^{\otimes n}(g) P_S.$$

Either of these applied to $v_1 \otimes \dots \otimes v_n$ gives:

$$\frac{1}{n!} \sum_{\sigma \in S_n} \rho(g)v_{\sigma(1)} \otimes \dots \otimes \rho(g)v_{\sigma(n)}$$

Since p_s is an intertwiner,

$$S^n V = \text{range } p_s \subseteq V^{\otimes n}$$

is a subrep. of $V^{\otimes n}$.

We need to check:

$$x \in S^n V \Rightarrow \rho^{\otimes n}(g)x \in S^n V.$$

pf: $x \in S^n V \Rightarrow x \in \text{range } p_s$

$$\Rightarrow \exists y \in V^{\otimes n} \text{ st } x = p_s(y)$$

$$\Rightarrow \rho^{\otimes n}(g)x = \rho^{\otimes n}(g)p_s y$$

$$= p_s \underbrace{\rho^{\otimes n}(g)y}_{\text{range } p_s}$$

$$\Rightarrow \rho^{\otimes n}(g)x \in S^n V.$$

Similarly - $\Lambda^n V = \text{range } p_A \subseteq V^{\otimes n}$

is also a subrep, since p_A is also an intertwiner.

Now let (ρ, V) be the defining rep of $G = SL(2, \mathbb{C})$ on $V = \mathbb{C}^2$. We get

$$\text{reps } \left(\begin{array}{c|c} \rho^{\otimes n} & S^n V \\ \hline & S^n V \end{array}, S^n V \right) := (S^n \rho, S^n V)$$

on the symmetrized tensor powers of V . Since $SU(2, \mathbb{C}) \subseteq SL(2, \mathbb{C})$, we get a rep of $SU(2)$, and it's irreducible in both cases.

($SL(2)$ -double cover of Lorentz grp)

Note: If smaller space has no invariant subspace, then neither does larger space.

These are irreducible, even when we restrict them to $SU(2) \subseteq SL(2, \mathbb{C})$. In fact, these are all the irreducible reps of $SU(2)$; but let's just show they are irreducible.

Thm: $(S^n \rho, S^n V)$ is an irrep. of $SU(2)$.

pf: $V = \mathbb{C}^2$ has a basis

$$\uparrow = \text{spin up} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \downarrow = \text{spin down} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$J_z = \begin{pmatrix} 1/2 & 0 \\ 0 & -1/2 \end{pmatrix}$$

$$J_z \uparrow = 1/2 \uparrow \quad \text{spin up state}$$

$$J_z \downarrow = -1/2 \downarrow \quad \text{spin down state}$$

$V^{\otimes n}$ has a basis of vectors like

$$\uparrow \otimes \downarrow \otimes \downarrow \otimes \uparrow \otimes \uparrow \otimes \downarrow \quad (n=6)$$

$S^n V$ has a basis of vectors like

$$P_S(\uparrow \otimes \downarrow \otimes \uparrow \otimes \uparrow \otimes \uparrow \otimes \downarrow) =$$

$$P_S(\uparrow \otimes \uparrow \otimes \uparrow \otimes \uparrow \otimes \downarrow \otimes \downarrow)$$

i.e. $P_S(\underbrace{\uparrow \otimes \uparrow \otimes \dots \otimes \uparrow}_n) = \uparrow^n$

$$P_S(\underbrace{\uparrow \otimes \uparrow \otimes \dots \otimes \uparrow}_{n-1} \otimes \downarrow) = \uparrow^{n-1} \downarrow$$

\vdots

$$P_S(\underbrace{\downarrow \otimes \downarrow \otimes \dots \otimes \downarrow}_n) = \downarrow^n$$

$$\dim S^n V = n+1$$

$S^n V = \{ \text{homogeneous degree-}n \text{ polynomials in 2 variables} \}$

$S^n \mathbb{C}^k = \{ \text{homogeneous degree-}n \text{ polynomials in } k \text{ variables} \}$

Want to show these form an irrep of $SU(2)$.

Suppose $U \subseteq S^n V$ is a subrepresentation:

$$S^n \rho(g): U \longrightarrow U \quad \forall g \in SU(2)$$

Need to show $U = \{0\}$ or $U = S^n V$.

Assume $U \neq \{0\}$. Let $0 \neq \psi \in U$, and we'll show $U = S^n V$.

Recall - if $X \in \mathfrak{su}(2)$, then $\exp(tX) \in SU(2)$,

so -

$$S^n \rho(\exp(tX)) \psi \in U.$$

Differentiating w/r/t t :

(how to get a rep of Lie alg from rep of Lie grp)

$$\frac{d}{dt} S^n \rho(e^{tX}) \psi \in U \quad \forall t.$$

$$\text{Recall - } \left. \frac{d}{dt} S^n \rho(e^{tX}) \right|_{t=0} = dS^n \rho(X)$$

and $dS^n \rho$ is a rep of $\mathfrak{su}(2)$ on $S^n V$.

So - $dS^n \rho(X) \psi \in U \quad \forall X \in \mathfrak{su}(2)$.

Take X to be:

$$J_+ = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$$

$$J_- = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

$$\mathfrak{su}(2) = \{x: \mathbb{C}^2 \rightarrow \mathbb{C}^2 \mid \text{tr}(x) = 0, x^* = -x\}$$

Note - J_+ , J_- aren't in $\mathfrak{su}(2)$, they're in $\mathfrak{sl}(2)$.

$$J_+ = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} - \frac{i}{2} \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}$$

$$J_- = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} = -\frac{1}{2} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} - \frac{i}{2} \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}$$

↑
in $\mathfrak{sl}(2)$

↑ ↑
these are in $\mathfrak{su}(2)$.

J_+ , J_- are linear combs. of complex multiples of matrices in $\mathfrak{su}(2)$.

* In fact, every matrix in $\mathfrak{sl}(2, \mathbb{C})$ is of the form $x + iy$ where $x, y \in \mathfrak{su}(2)$.

Thus, starting from our rep. dS^n of $\mathfrak{su}(2)$ we can define a rep α of $\mathfrak{sl}(2, \mathbb{C})$ by:

$$\alpha(x + iy) = \alpha(x) + i\alpha(y)$$

$$:= dS^n \rho(x) + i dS^n \rho(y)$$

(replacing i by $-i$ we get conjugate rep.)

Since $dS^n \rho(x) \psi \in U \quad \forall x \in su(2)$ we also know

$$\alpha(x+iy) \psi \in U \quad \forall x+iy \in sl(2, \mathbb{C})$$

(by defn of $\alpha(x+iy)$).

Write $\psi \in S^n V$ as:

$$\psi = a_0 \uparrow^n + a_1 \uparrow^{n-1} \downarrow + \dots + a_n \downarrow^n$$

$$U \ni \alpha(J_-) \psi = a_0 n \uparrow^{n-1} \downarrow + a_1 (n-1) \uparrow^{n-2} \downarrow^2 + \dots +$$

takes one
 \uparrow & turns
it to \downarrow .

$$a_{n-1} \downarrow^n + 0$$

what it does to \downarrow^n .

Since U is an invariant subspace, we can do this again & end up in U .

$$U \ni [\alpha(J_-)]^2 \psi = a_0 n(n-1) \uparrow^{n-2} \downarrow^2 + \dots$$

Continuing doing this — we get

⋮

$$U \ni [\alpha(J_-)]^k \psi = (\text{some } \#) \downarrow^n$$

for some k , we get all spin down state.

We know \downarrow^n is in U .

We can do the same thing w/ J_+ :

$$\alpha(J_+) \downarrow^n = n \uparrow \downarrow^{n-1}$$

$$\alpha(J_+)^2 \downarrow^n = n(n-1) \uparrow^2 \downarrow^{n-2}$$

\vdots

so we get a basis of $S^n V$ lying in U , so

$$\downarrow^n, \uparrow^n \in U \Rightarrow U = S^n V. \quad \square$$

Riemann sphere, heavenly sphere. (space of directions we can look), same as states of spin- $1/2$ reps.

We call $(S^n \rho, S^n V)$ the "spin- j rep" of $SU(2)$ or of $SL(2, \mathbb{C})$ where $j = n/2$.

So (ρ, V) is called spin- $1/2$ representation.

V itself is closely related to the heavenly sphere: $\mathbb{C}P^1 = S^2$ is space of all 1-d subspaces of $V = \mathbb{C}^2$.

Let $L_x \subseteq \mathbb{C}^2$ be the 1-dim'l ^{complex} subspace corresp. to the point $x \in \mathbb{C}P^1$.

L_x is a 1-dim'l vector bundle (have a v. space \forall pt)

$\mathbb{C}P^1$ -
Riemann
sphere

There's a 1-dim'l vector bundle L over $\mathbb{C}P^1$
w/ the L_x 's as its fibers.



Given $f \in V^*$, it restricts
to a linear functional on each
 $L_x \cong V$.

i.e. we get $f_x \in L_x^*$. So we get a section¹ of
the bundle L^* whose fibers are L_x^* .

(section - $\forall x$, we get an elt. of fiber L_x)

This is a holomorphic section of L^* (locally look
like holomorphic functs) and in fact:

Thm - All holomorphic sections of L^* are of
the form \tilde{f} for some $f \in V^*$.

(only holomorphic functs on Riemann sphere are
constants by Liouville.)

If $\Gamma(L^*) = \{ \text{holomorphic sections of } L^* \}$
we get

$$\Gamma(L^*) \cong V^* \quad (\text{restatement of above thm}).$$

$$\Gamma(L^*) \cong V^* \cong V$$

(as reps of $SL(2, \mathbb{C})$.)