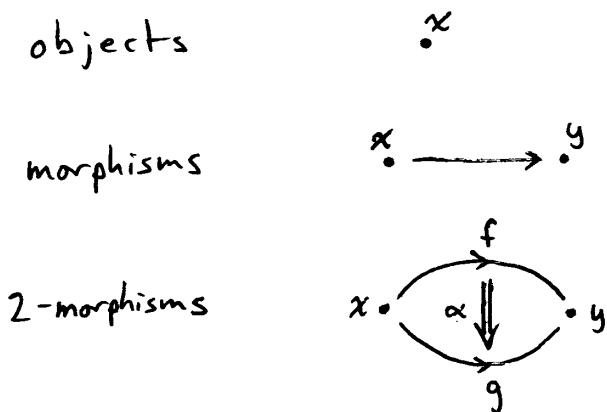


5 October 2004 !

- Bénabou (1967) - invented "bicategories" (aka "weak 2-categories"), which have:



- i.e. like categories but with 2-morphisms $\alpha: f \Rightarrow g$ between parallel 1-morphisms $f, g: x \rightarrow y$. As in category theory, we can compose 1-morphisms:

$$f: x \rightarrow y, g: y \rightarrow z \text{ give } fg: x \rightarrow z$$

& we have identity 1-morphisms

$$1_x: x \rightarrow x$$

but these only satisfy associativity, left & right unit laws up to specified 2-isomorphism:

$$\alpha_{f,g,h}: (fg)h \xrightarrow{\sim} f(gh)$$

$$l_f: 1_x f \xrightarrow{\sim} f$$

$$r_f: f 1_x \xrightarrow{\sim} f$$

What's a 2-isomorphism? Well, you can also compose 2-morphisms:

$$\alpha: f \Rightarrow g \quad \beta: g \Rightarrow h \quad \text{give} \quad \alpha\beta: f \Rightarrow h$$

& there are identity 2-morphisms:

$$1_f: f \Rightarrow f$$

which satisfy associativity, left & right unit laws "on the nose" — i.e. as equations. Thus we can talk about the inverse of a 2-morphism, and an invertible 2-morphism is called a 2-isomorphism.

Also: a (associator), λ (left unit law), — should be called the "uniter"!
 r (right unit law) satisfy some coherence conditions:

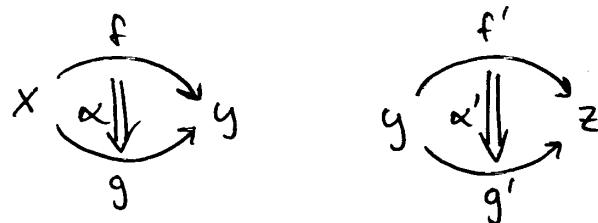
$$\begin{array}{ccc}
 (f(gh))_k & \xrightarrow{\alpha_{f,gh,k}} & f((gh)_k) \\
 \nearrow \alpha_{f,g,h} \cdot 1_k & & \searrow 1_f \cdot \alpha_{g,h,k} \\
 ((fg)h)_k & & f(g(hk)) \\
 \searrow \alpha_{fg,h,k} & & \nearrow \alpha_{f,g,hk} \\
 & (fg)(hk) &
 \end{array}$$

$$\begin{array}{ccc}
 (f1_x)g & \xrightarrow{\alpha_{f,1_x,g}} & f(1_x g) \\
 \searrow r_f \cdot 1_g & & \swarrow 1_f \cdot r_g \\
 & fg &
 \end{array}$$

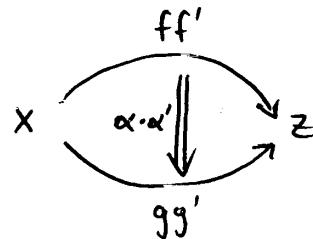
(and these two imply all other imaginable coherence laws)

To make sense of these we also need an operation called "horizontal composition" of 2-morphisms:

Given



we get



(The previous form of composition -

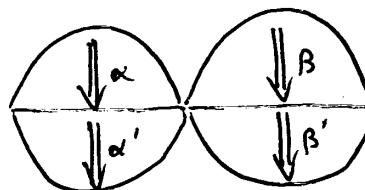


gives



- is called "vertical composition")

Lastly, we require the interchange law for horizontal & vertical composition:



$$(\alpha \cdot \beta)(\alpha' \cdot \beta') = (\alpha \alpha') \cdot (\beta \beta').$$

Bénabou had lots of examples in mind, e.g. :

- Cat is a bicategory with categories as objects functors as morphisms natural transformations as 2-morphisms.

(But this is a rather simplistic example — Cat is a strict 2-category, meaning that the associator and the unit laws are actually identity 2-morphisms)

- Mod is a bicategory with
 - rings as objects
 - (R, S) bimodules as morphisms from R to S
 - (R, S) bimodule homomorphisms as 2-morphisms

Composition of morphisms is tensoring of bimodules.
We have an associator

$$\alpha_{M,N,O} : (M \otimes_S N) \otimes_T O \xrightarrow{\cong} M \otimes_S (N \otimes_T O)$$

where $R \xrightarrow{M} S \xrightarrow{N} T \xrightarrow{O} U$.

- For any space X there's a bicategory $\mathbf{TT}_2(X)$ with
 - points in X as objects
 - paths in X as morphisms
 - homotopy classes of homotopies between paths in X as 2-morphisms.

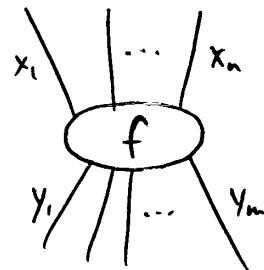
(where we need "homotopy classes" at the top level to get the laws involving 2-morphisms to hold on the nose).

- Just as a monoid is a category with one object, a monoidal category is a bicategory with one object!

monoidal category	bicategory
morphisms	2-morphisms
objects	morphisms
—	one object

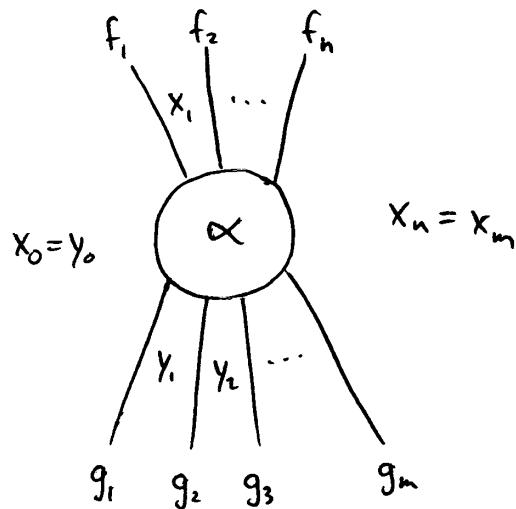
where "tensoring objects" in the monoidal category is secretly composing morphisms in the corresponding bicategory.

Recall that for monoidal categories we could draw a morphism $f: x_1 \otimes \dots \otimes x_n \rightarrow y_1 \otimes \dots \otimes y_m$ as

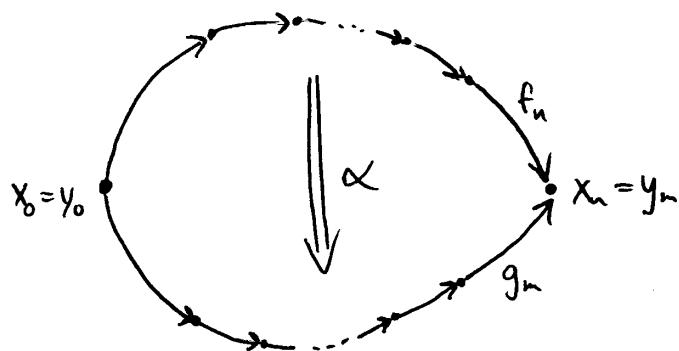


More generally, we can draw 2-morphisms in a bicategory in the same fashion. Given composable morphisms $f_1: x_0 \rightarrow x_1, f_2: x_1 \rightarrow x_2, \dots, f_n: x_{n-1} \rightarrow x_n$ & $g_1: y_0 \rightarrow y_1, \dots, g_m: y_{m-1} \rightarrow y_m$, this is how

we draw $\alpha: f_1 f_2 \dots f_n \Rightarrow g_1 g_2 \dots g_m$ (if $x_0 = y_0$ & $x_n = y_m$)



Note this is Poincaré dual to



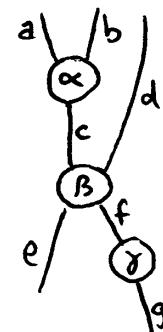
Given 2-morphisms

$$\alpha: a b \Rightarrow c$$

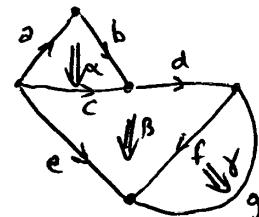
$$\beta: c d \Rightarrow e f$$

$$\gamma: f \Rightarrow g$$

we can combine them using both vertical and horizontal composition



or



7 October 2004

- Penrose (1971) - In GR physicists study "tensors", i.e. linear maps

$$f : \underbrace{V \otimes \cdots \otimes V}_n \longrightarrow \underbrace{V \otimes \cdots \otimes V}_m$$

& they use "index notation" to describe them. I.e. they would pick a basis e_i of V & say

$$f(e_{i_1} \otimes e_{i_2} \otimes \cdots \otimes e_{i_n}) = \sum_{1 \leq j_1, \dots, j_m \leq d} f_{i_1, \dots, i_n}^{j_1, \dots, j_m} e_{j_1} \otimes \cdots \otimes e_{j_m}$$

where $d = \dim V$.

- Mathematicians would mock them (and still do) for indulging in this "debauch of indices," but these indices are needed to describe the most general ways of "composing" tensors. E.g. given

$$\begin{array}{ll} S : V \otimes V \longrightarrow V \otimes V & S_{ij}^{kl} \\ T : V \longrightarrow V \otimes V & T_i^{jk} \end{array}$$

we can form all sorts of other tensors, e.g.: ordinary composition

$$ST : V \longrightarrow V \otimes V$$

$$(ST)_i^{jk} = S_{lm}^{jk} T_i^{lm}$$

(following the Einstein summation convention)

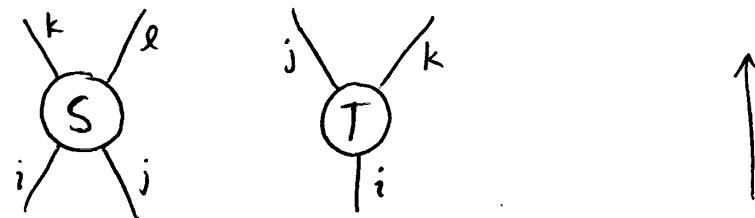
But we could also form

$$(T \otimes 1_V)S : V \otimes V \rightarrow V \otimes V \otimes V$$

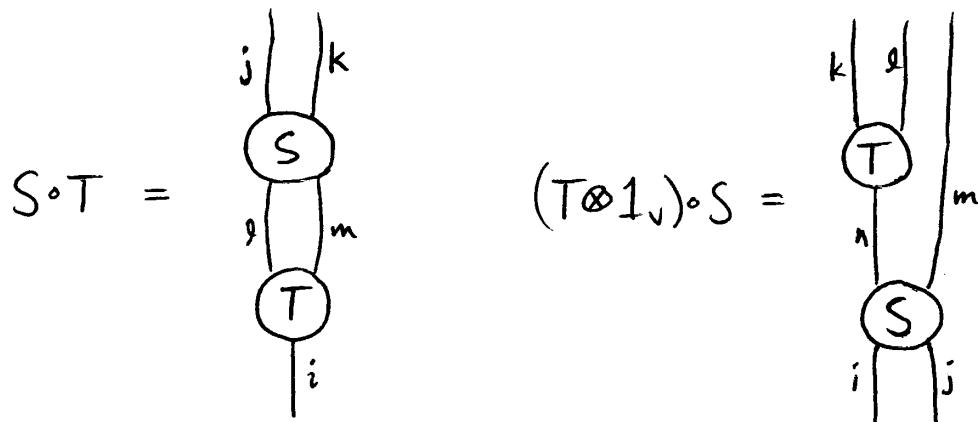
$$((T \otimes 1_V)S)_{ij}^{klm} = T_n^{kl} S_{ij}^{nm}$$

and many many more...

Penrose noticed that in these expressions, the indices don't really need to refer to a specific basis of a vector space. — the results are basis-independent; the indices only say how the tensors are hooked up. This is his "abstract index notation". He also started drawing tensors as "black boxes"

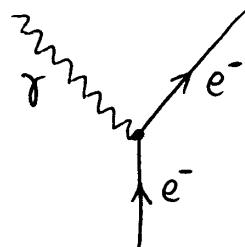


& composing them diagrammatically:



And we can now even leave out the indices — the 'wires' do the work.

Nowadays we know this notation works in any monoidal category, or even a bicategory, where they're called "string diagrams". Penrose didn't know about monoidal categories, but he did know about Feynman diagrams, and he realized what the relationship was.



In Feynman diagrams, each kind of particle has a specific vector space (Hilbert space) of states; these are objects in Vect. The vertices of Feynman diagrams describe interactions; these are morphisms in Vect. In the theory of Feynman diagrams, all the vector spaces are actually representations of some symmetry group G (e.g. Poincaré group). Also, the vertices actually represent morphisms between representations — i.e. intertwiners. In short, our monoidal category is not Vect but $\text{Rep}(G)$ — the category of representations of G .

Penrose played the same game but taking $G = SU(2)$.

This group has one irrep of each dimension, called
(by physicists) the "spin-0" ($= 1d$), "spin- $\frac{1}{2}$ " ($= 2d$),
"spin-1" ($= 3d$), ... representations.