

Smooth Functors And Beyond

We've seen that if M is a smooth space (e.g. a manifold) then there is a smooth category PM where:

- objects are points $x \in M$
- morphisms $y: x \rightarrow y$ are thin homotopy classes of smooth maps $f: [0, 1] \rightarrow M$ with $f(0) = x$, $f(1) = y$
i.e. f constant near 0 & 1.

We think of PM as a category of "configurations" i.e. "processes" for some physical system. So to formulate the

Lagrangian approach to the physics of this system, we need a smooth functor

$$S: PM \rightarrow \mathbb{R}$$

describing the "action" of any process.

Last time we saw:

Theorem - There's a 1-1 correspondence
between smooth functors

$$S: PM \rightarrow \mathbb{R}$$

of 1-forms A on M , given by:

$$S(\gamma) = \int_{\gamma} A$$

where we pick a representative (path)
for γ to define the integral.

Alas, this isn't general enough... as we'll
soon see.

To do quantum physics, what matters is not

$$S: PM \rightarrow \mathbb{R}$$

but the phase

$$e^{is} : PM \rightarrow U(1)$$

which has less information since

$$\exp : \mathbb{R} \rightarrow U(1)$$

is many-to-one. In fact, e^{is}
is also sufficient to do classical physics!



$$\frac{d}{ds} S(\gamma_s) = 0$$

If we seek critical points of the action (instead of minima), we can work with e^{is} instead of S :

$$\frac{d}{ds} e^{is(\gamma_s)} = 0$$

(for all smooth homotopies γ_s of γ holding endpoints fixed).

(4)

The critical points of e^{iS} are the same as those of S , so this doesn't seem like a big deal.

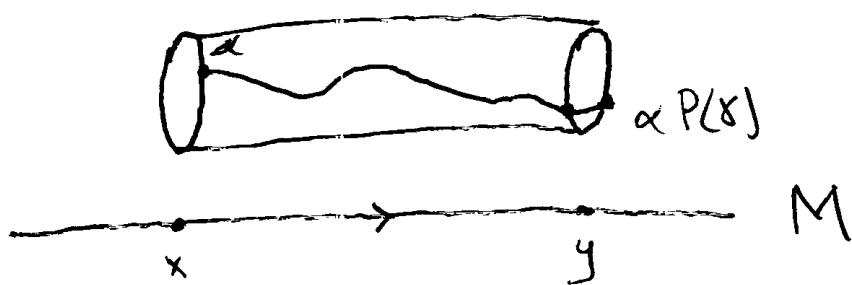
Theorem — There's a 1-1 correspondence between smooth functions

$$P: PM \rightarrow U(1)$$

is 1-forms A on M , given by :

$$P(\gamma) = e^{i\int_A \gamma}$$

Here P stands for "phase". So far the picture looks like:



(5)

For each point $x \in M$ we have a circle of possible phases for the system in configuration x , so we have a "trivial principal $U(1)$ bundle":

$$\begin{array}{ccc} M \times U(1) & \rightarrow & (x, \alpha) \\ \downarrow \pi_1 & & \downarrow \pi_2 \\ M & & x \end{array}$$

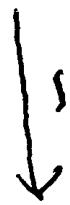
Sitting over $x \in M$ we have a fiber

$$\pi_1^{-1}(x) \subseteq M \times U(1)$$

which is a circle - the set of possible phases our system could have at x .

This example is called "trivial" because each fiber is $U(1)$ - or is canonically isomorphic to $U(1)$:

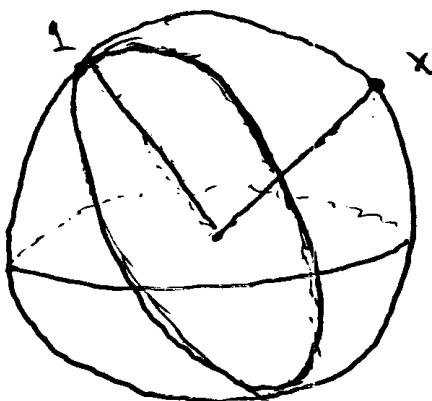
$$\pi_i^{-1}(x) = \{ (x, \alpha) : \alpha \in U_i \}$$



$$U_i$$

where the isomorphism sends (x, α) to α .

More interesting are the nontrivial principal $U(1)$ bundles:

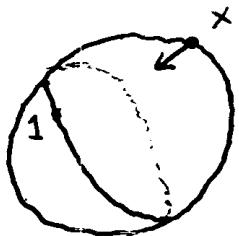


$$M = S^2 = \mathbb{C}P^1$$

For example, let the fiber over x be the set of points in S^2 that are \perp to x .

(7)

We can't smoothly identify all the fibers with $U(1)$ since that would produce a nowhere vanishing smooth vector field on S^2 :



Now let's get a bit more formal.

What's the difference between a circle and the circle? The circle is $U(1) \subseteq \mathbb{C}$.

A circle is a "U(1)-torsor". - a copy of $U(1)$ that's forgotten what the element 1 is.

Def. - For any group G , a G -torsor is a set X equipped with an action (\sim right action) of G :

$$\begin{aligned}\alpha: X \times G &\longrightarrow X \\ (x, g) &\longmapsto xg\end{aligned}$$

s.t.

$$x1 = x$$

$$(xg)h = x(gh)$$

such that X is isomorphic to G as a space with right G -action: there's a bijection

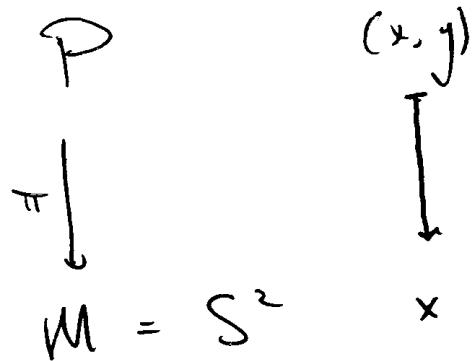
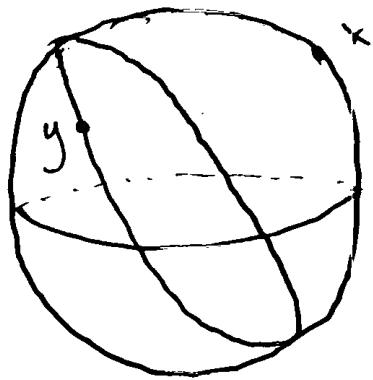
$$\beta: X \rightarrow G$$

s.t.

$$\beta(xg) = \beta(x)g$$

If $G = U(1)$, the difference between right & left actions is inessential since $U(1)$ is abelian. More importantly, any circle equipped with the ability to rotate it by any phase $g \in U(1)$ is a $U(1)$ -torsor.

E.g.



if a point in P is a point $x \in S^1$ together with a point in the circle \perp to x , if $\pi: P \rightarrow M$ is the obvious map, then $\pi^{-1}(x)$ is a $U(1)$ -torsor.

More precisely, $\pi^{-1}(x)$ becomes a $U(1)$ -torsor after we pick a "right" or left-hand "rule" for rotating $y \in \pi^{-1}(x)$ by a phase $g \in U(1)$.

A principal $U(1)$ bundle is (among other things) a smooth space P with a smooth map

$$\begin{matrix} P \\ \pi \downarrow \\ M \end{matrix}$$

such that each fiber $\pi^{-1}(x)$ ($x \in M$) is equipped with the structure of being a $U(1)$ -torsor.

Next time we'll really define "principal $U(1)$ -bundle" & include a clause saying that the $U(1)$ -torsor structure on $\pi^{-1}(x)$ varies smoothly with x .