

***k*-colourings**  
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**1.** *k*-coloring a finite set.

A *k*-colouring of the set  $n$  is a function  $f: n \rightarrow k$ . For each element of  $n$  there are  $k$  choices of the image of  $f$ , and the choices for each of the elements are all independent, so there are  $k^n$  possible  $f: n \rightarrow k$ . Hence,

$$|C(k)_n| = k^n.$$

**2.** The generating function of *k*-colourings.

$$|C(k)|(z) = \sum_{n \geq 0} \frac{k^n}{n!} z^n = e^{kz}.$$

**3, 4.** Annihilation.

Trivially,

$$a|C(k)|(z) = \frac{d}{dz} e^{kz} = k e^{kz} = k|C(k)|(z).$$

Suppose that  $|C(w)|(z)$  is an eigenvector of  $a$  with eigenvalue  $w$ , that is,

$$\frac{d}{dz} |C(w)|(z) = w |C(w)|(z).$$

Letting

$$|C(w)|(z) = \sum_{n \geq 0} \frac{|C(w)_n|}{n!} z^n,$$

the eigenvalue equation becomes

$$\sum_{n \geq 1} \frac{|C(w)_n|}{(n-1)!} z^{n-1} = \sum_{n \geq 0} \frac{w |C(w)_n|}{n!} z^n$$

so

$$\sum_{n \geq 0} \frac{|C(w)_{n+1}|}{n!} z^n = \sum_{n \geq 0} \frac{w |C(w)_n|}{n!} z^n$$

and  $|C(w)_{n+1}| = w |C(w)_n|$ . It follows that  $|C(w)_n| = w^n |C(w)_0|$ , and that

$$|C(w)|(z) = |C(w)_0| e^{wz},$$

so for each complex number  $w$  there is a one-dimensional space of eigenvectors of the annihilation operator with eigenvalue  $w$ .

**5.** Categorized eigenvalue problem.

Let  $K$  be the structure type such that putting it on a set  $S$  is “picking a colour out of  $k$  and  $S$  is empty”. It follows that  $|K|(z) = k$ .

We seek a structure type  $T_k$  such that

$$AT_k \simeq K \times T_k.$$

Observe that “putting an  $AT_k$  structure on a set  $S$ ” is “putting a  $T_k$  structure on the set  $S+1$ ”, and “putting a  $K \times T_k$  structure on the set  $S$ ” is the same as “picking a colour out of  $k$  and putting a  $T_k$  structure on  $S$ ”. That is,

putting a  $T_k$  structure on  $S + 1$  is the same as picking a colour out of  $k$  and putting a  $T_k$  structure on  $S$ .

It easily follows by induction on  $|S|$  that a  $T_k$  structure on the set  $S$  is a  $T_k$  structure on the empty set, and a  $k$ -colouring of  $S$ . How many ways there are to put a  $T_k$  structure on the empty set is undetermined, but it is basically the structure type  $K'$  for some integer  $k'$ . It follows that

$$T_k \simeq K' \times C(k).$$

### 6. Inner product on Fock space.

Assume, without loss of generality, that  $n \geq m$ . Then,

$$\langle z^n, z^m \rangle = \langle (a^*)^n 1, z^m \rangle = \langle 1, a^n z^m \rangle = \left\langle 1, \frac{d^n z^m}{dz^n} \right\rangle = \delta_{n,m} m! = \delta_{n,m} n!.$$

### 7. Normalizing coherent states.

Let  $\psi_w = e^{wz}$ . Then,

$$\langle \psi_w, \psi_w \rangle = \left\langle \sum_{n \geq 0} \frac{w^n}{n!} z^n, \sum_{n \geq 0} \frac{\bar{w}^n}{n!} z^n \right\rangle = \sum_{n,m \geq 0} \frac{\bar{w}^n w^m}{n! m!} \langle z^n, z^m \rangle = \sum_{n \geq 0} \frac{|w|^{2n}}{n!} = e^{|w|^2}.$$

Hence, the normalized coherent state is

$$\psi_w = e^{-\frac{|w|^2}{2} + wz}$$

While we're at it, we are going to need  $\|z\psi_w\|^2$  later:

$$\langle z\psi_w, z\psi_w \rangle = e^{-|w|^2} \sum_{n \geq 0} \frac{|w|^{2n} (n+1)}{n!} = |w|^2 + 1.$$

### 8,9,10,11. Heisenberg uncertainty for coherent states.

We use the facts that  $q = \frac{1}{\sqrt{2}}(a + a^*)$  and that  $p = \frac{1}{i\sqrt{2}}(a - a^*)$ . Then,

$$\begin{aligned} \langle \psi_w, q\psi_w \rangle &= \frac{1}{\sqrt{2}} (\langle \psi_w, a\psi_w \rangle + \langle a\psi_w, \psi_w \rangle) = \frac{w + \bar{w}}{\sqrt{2}} \\ \langle \psi_w, p\psi_w \rangle &= \frac{w - \bar{w}}{\sqrt{2}i}. \end{aligned}$$

Similarly,

$$\begin{aligned} \langle \psi_w, q^2\psi_w \rangle &= \langle q\psi_w, q\psi_w \rangle = \frac{1}{2} (\langle a\psi_w, a\psi_w \rangle + \langle a^2\psi_w, \psi_w \rangle + \langle \psi_w, a^2\psi_w \rangle + \langle a^*\psi_w, a^*\psi_w \rangle) \\ &= \frac{1}{2} (|w|^2 + \bar{w}^2 + w^2 + |w|^2 + 1) = \langle \psi_w, q\psi_w \rangle^2 + \frac{1}{2} \\ \langle \psi_w, p^2\psi_w \rangle &= \frac{1}{2} (|w|^2 - \bar{w}^2 - w^2 + |w|^2 + 1) = \langle \psi_w, p\psi_w \rangle^2 + \frac{1}{2} \end{aligned}$$

It follows that  $\sigma_{\psi_w}^2 p = \sigma_{\psi_w}^2 q = \frac{1}{2}$ , as required.