

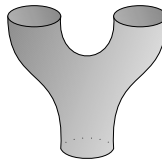
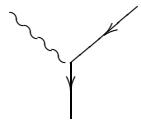
Physics, Topology, Logic and Computation: a Rosetta Stone

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California State University, Fresno

Categories	Physics	Topology	Logic	Computation
object	system	manifold	proposition	data type
morphism	process	cobordism	proof	program

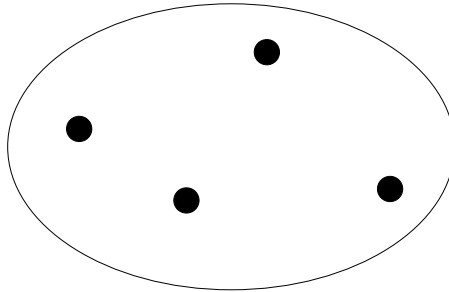


$X \& Y \vdash Z$

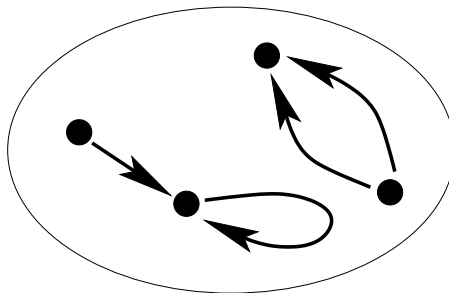
$\lambda x \lambda y . x(y)$

The Big Idea

Once upon a time, mathematics was all about *sets*:



In 1945, Eilenberg and Mac Lane introduced *categories*:



These put *processes* on an equal footing with *things*.

In physics, we often use categories where:

- **objects represent physical *systems*;**
- **morphisms represent physical *processes*.**

In classical physics we often use the category Set, where:

- **an object is a *set***
- **a morphism is a *function***

In quantum physics we often use Hilb, where:

- **an object is a *Hilbert space***
- **a morphism is a *linear operator***

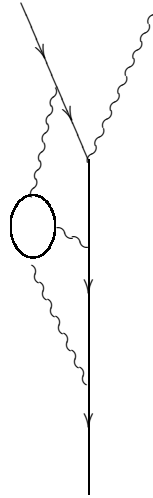
A category C consists of:

- **A collection of *objects*. If X is an object of C we write $X \in C$.**
- **For any $X, Y \in C$, a set of *morphisms* $f: X \rightarrow Y$.**

We require that:

- **Every $X \in C$ has an *identity* morphism $1_X: X \rightarrow X$.**
- **Given $f: X \rightarrow Y$ and $g: Y \rightarrow Z$, there is a *composite* morphism $gf: X \rightarrow Z$.**
- **The *unit laws* hold: if $f: X \rightarrow Y$, then $f1_X = f = 1_Yf$.**
- **Composition is *associative*: $(hg)f = h(gf)$.**

Feynman used diagrams to describe processes in quantum physics:



Now we know that these are pictures of *morphisms* — so we can use these diagrams in other contexts!

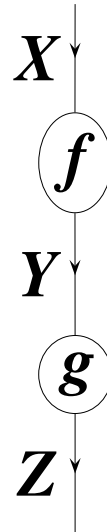
We can draw a morphism

$$f: X \rightarrow Y$$

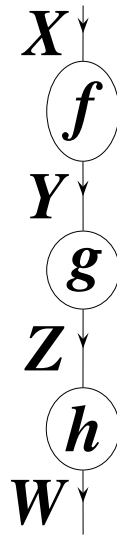
like this:



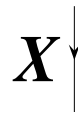
We draw the composite of $f : X \rightarrow Y$ and $g : Y \rightarrow Z$ like this:



Then the associative law is implicit:



If we draw the identity morphism $1_X: X \rightarrow X$ like this:



the unit laws are implicit too!

For theories with at least 1 dimension of space, we need *monoidal* categories.

Here any pair of morphisms $f: X \rightarrow Y, f': X' \rightarrow Y'$ has a tensor product

$$f \otimes f' : X \otimes X' \rightarrow Y \otimes Y'$$

We use this to describe parallel processes:

$$\begin{array}{ccc}
 \begin{array}{c} X \downarrow \\ \textcircled{f} \\ Y \downarrow \end{array} & \begin{array}{c} X' \downarrow \\ \textcircled{f'} \\ Y' \downarrow \end{array} & = & \begin{array}{c} X \otimes X' \downarrow \\ \textcircled{f \otimes f'} \\ Y \otimes Y' \downarrow \end{array}
 \end{array}$$

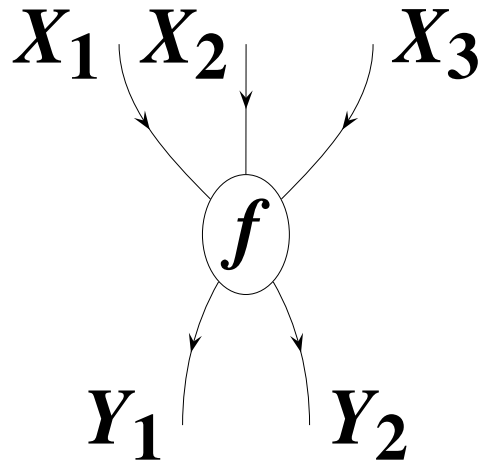
Examples:

- The category **Hilb**, with its usual tensor product \otimes .
- The category **Set**, with the cartesian product \times .

More generally, we can draw any morphism

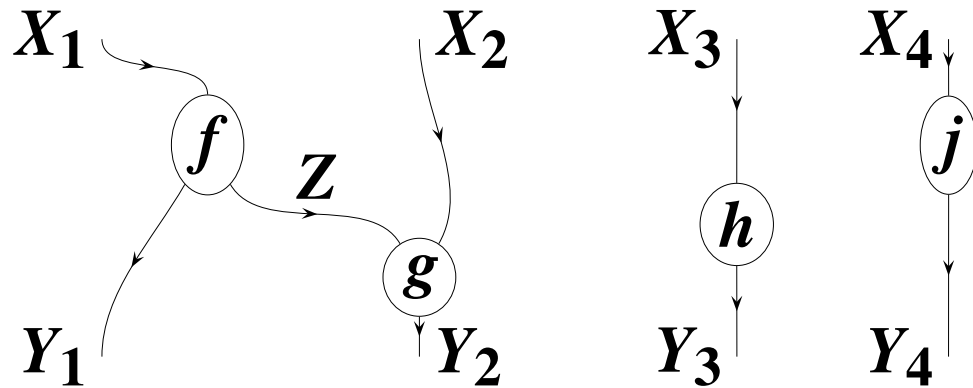
$$f: X_1 \otimes \cdots \otimes X_n \rightarrow Y_1 \otimes \cdots \otimes Y_m$$

like this:

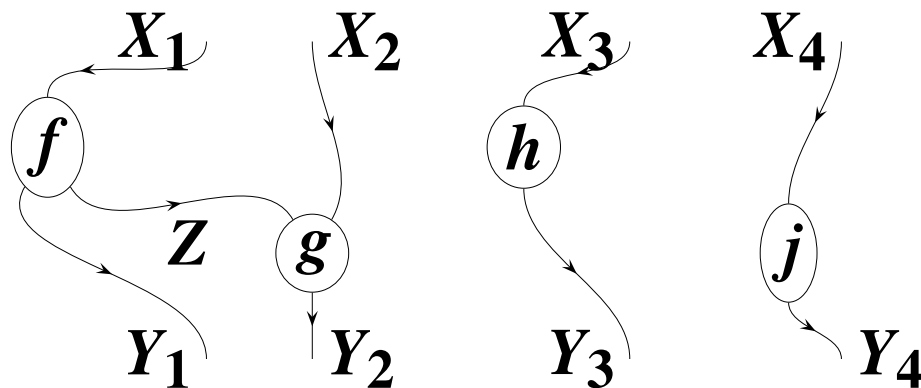


In physics we use this to depict an interaction between particles.

By composing and tensoring, we can build up bigger diagrams:



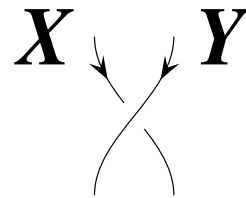
The monoidal category axioms let us deform the picture without changing the morphism:



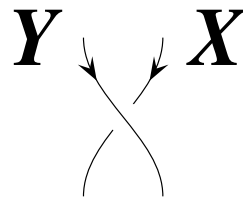
In theories with least 2 dimensions of space, we use *braided* monoidal categories. We can draw the braiding

$$B_{X,Y}: X \otimes Y \rightarrow Y \otimes X$$

like this:



It has an inverse, drawn like this:



Then we have:

$$\begin{array}{c} X & & Y \\ & \searrow & / \\ & & \\ & / & \searrow \\ X & & Y \end{array} = \begin{array}{c} X \\ \downarrow \\ X \\ \downarrow \\ Y \end{array} = \begin{array}{c} X & & Y \\ & \searrow & / \\ & & \\ & / & \searrow \\ X & & Y \end{array}$$

In theories with at least 3 dimensions of space, we use *symmetric* monoidal categories, where:

$$\begin{array}{c} X & & Y \\ & \searrow & / \\ & & \\ & / & \searrow \\ X & & Y \end{array} = \begin{array}{c} X & & Y \\ & / & \searrow \\ & & \\ & \searrow & / \\ X & & Y \end{array}$$

The most familiar braided monoidal categories are symmetric:

- **In Set with its cartesian product, the standard braiding is:**

$$\begin{aligned} B_{X,Y} &: X \times Y \rightarrow Y \times X \\ (x, y) &\mapsto (y, x) \end{aligned}$$

- **In Hilb with its usual tensor product, the standard braiding is:**

$$\begin{aligned} B_{X,Y} &: X \otimes Y \rightarrow Y \otimes X \\ x \otimes y &\mapsto y \otimes x \end{aligned}$$

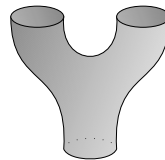
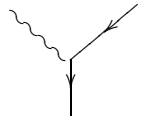
However, in thin films there can be ‘anyons’. These are particle-like excitations described by braided monoidal categories that are *not* symmetric!

- **Superconducting films: the quantum Hall effect.**
- **Graphene (single-layer graphite): fractional-charge anyons are possible, not yet seen.**

But there's a lot more to this story...

THE ROSETTA STONE

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object	system	manifold	proposition	data type
morphism	process	cobordism	proof	program



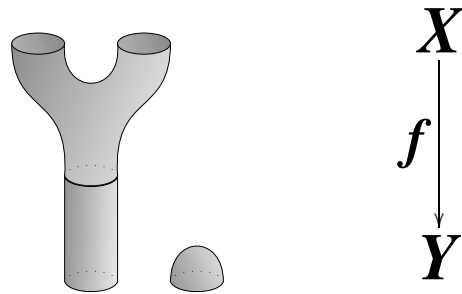
$X \& Y \vdash Z$

$\lambda x \lambda y . x(y)$

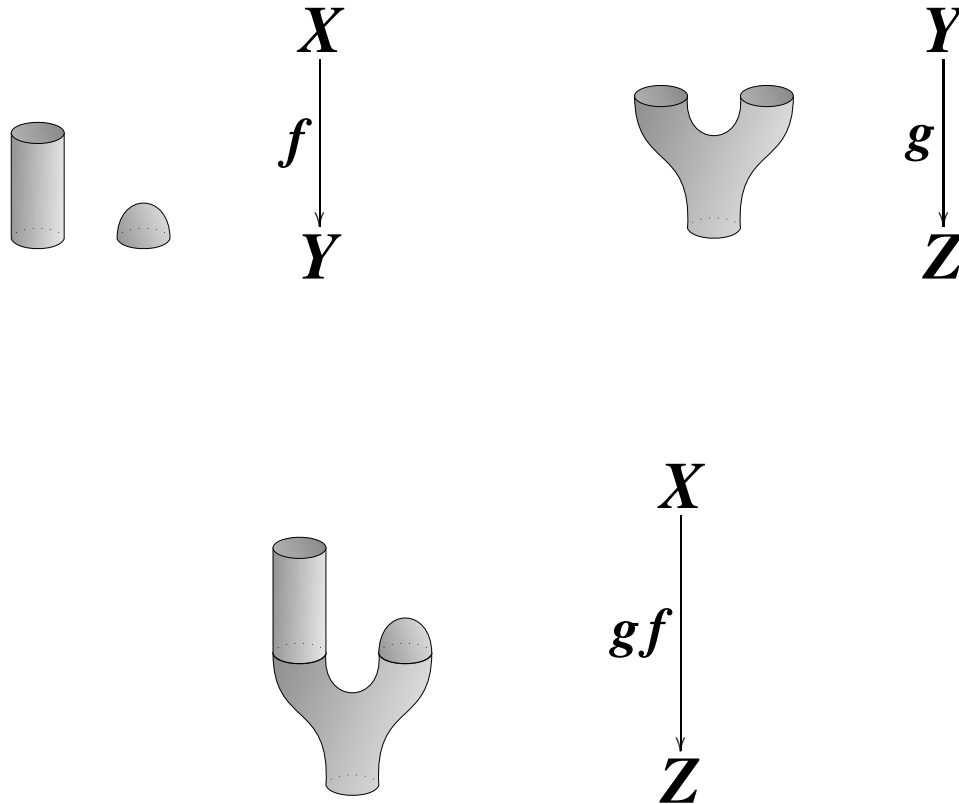
In topology, there is a category $n\text{Cob}$ where:

- objects are $(n - 1)$ -dimensional *manifolds*;
- morphisms are *cobordisms*.

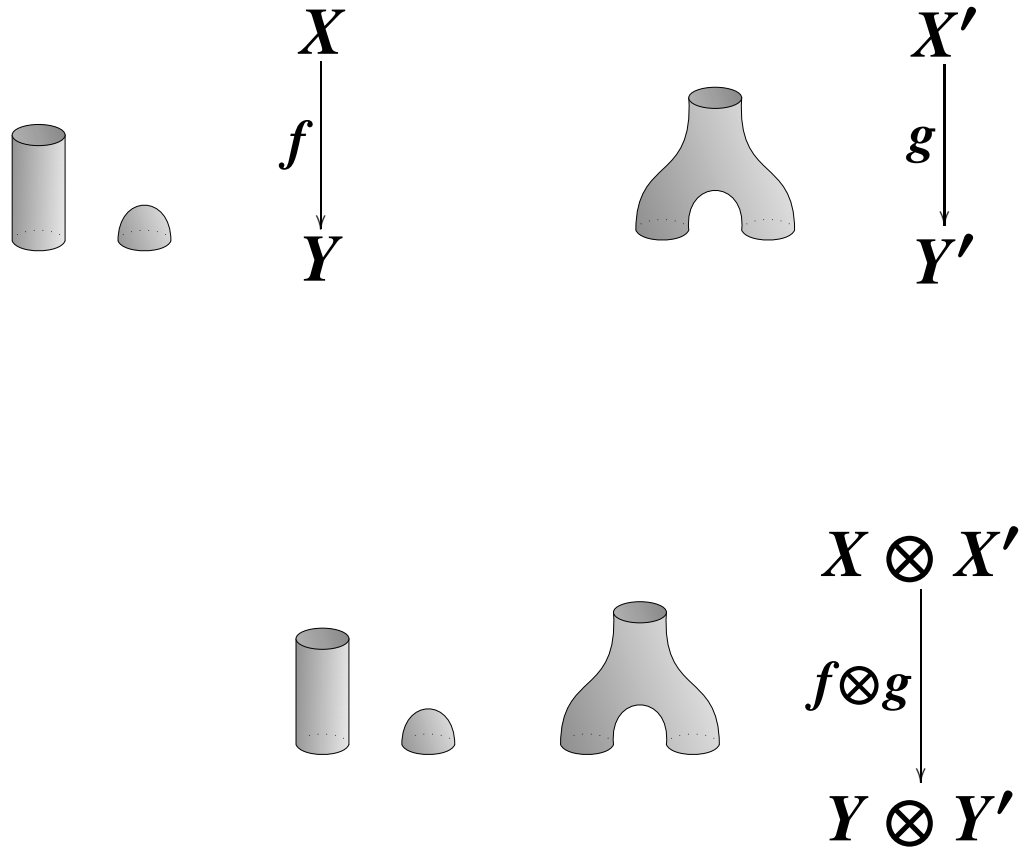
A *cobordism* $f : X \rightarrow Y$ is an n -dimensional manifold whose boundary is the disjoint union of X and Y . For example, when $n = 2$:



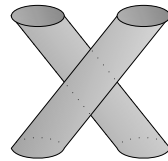
We compose cobordisms by gluing the ‘output’ of one to the ‘input’ of the other:



$n\text{Cob}$ is a monoidal category. We tensor cobordisms by taking their disjoint union:



In fact, $n\text{Cob}$ is a symmetric monoidal category:



$$\begin{array}{c} X \otimes Y \\ \downarrow B_{X,Y} \\ Y \otimes X \end{array}$$

In general relativity, objects in $n\text{Cob}$ describe choices of *space*, while morphisms describe choices of *space-time*. I believe that:

Quantum theory will eventually make more sense, as part of a theory of quantum gravity — but this can only be understood using categories.

Why? The weird features of quantum theory come from the ways that Hilb is less like Set than $n\text{Cob}$. But $n\text{Cob}$ is what we use to describe space and space-time in general relativity!

‘Weird’ properties of quantum theory correspond to unsurprising properties of spacetime.

	object ●	morphism ● → ●
SET THEORY	set	function between sets
QUANTUM THEORY	Hilbert space (state)	operator between Hilbert spaces (process)
GENERAL RELATIVITY	manifold (space)	cobordism between manifolds (spacetime)

For example: Set is ‘cartesian’, while $n\text{Cob}$ and Hilb are not.

If a symmetric monoidal category is cartesian, you can do various things including *duplication*:

$$\Delta_X : X \rightarrow X \otimes X$$

In Set we can duplicate as follows:

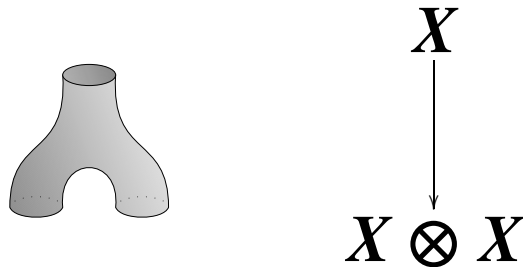
$$\begin{aligned} \Delta_X &: X \rightarrow X \times X \\ x &\mapsto (x, x) \end{aligned}$$

In Hilb we cannot duplicate: the function

$$\begin{aligned} X &\rightarrow X \otimes X \\ x &\mapsto x \otimes x \end{aligned}$$

**is not linear! It's not a morphism in Hilb.
So: we 'cannot clone a quantum state'.**

**Similarly, in $n\text{Cob}$ there is no duplication, despite
this misleading picture for $n = 2$:**



**When $n = 1$ there's typically no cobordism from a
manifold X to $X \otimes X$, and similarly for $n = 4$.**

What about logic and computer science? These too study categories of things and processes:

In proof theory, we use categories where:

- **an object is a *proposition***
- **a morphism is a *proof***

In computer science, we use categories where:

- **an object is a *data type***
- **a morphism is a *program***

In proof theory $X \vdash Y$ means *assuming X, we can prove Y*. But we can also let it mean *the set of proofs leading from assumption X to conclusion Y*.

Since proofs are morphisms, we can compose them:

$$\frac{X \vdash Y \quad Y \vdash Z}{X \vdash Z}$$

The identity morphism:

$$\overline{X \vdash X}$$

Logic uses *monoidal* categories where the tensor product is ‘and’. We can tensor propositions, and tensor proofs:

$$\frac{W \vdash X \quad Y \vdash Z}{W \& Y \vdash X \& Z}$$

In fact, logic uses symmetric monoidal categories:

$$\frac{X \vdash Y \& Z}{X \vdash Z \& Y}$$

Classical logic is cartesian, so it permits duplication:

$$\frac{X \vdash Y}{X \vdash Y \& Y}$$

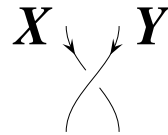
Linear logic does not!

A program that takes data of type X as input and returns data of type Y can be seen as a morphism $f: X \rightarrow Y$.

Categories of data types and programs are monoidal. Given data types X and X' there is a data type $X \otimes X'$. And given programs $f: X \rightarrow Y, f': X' \rightarrow Y'$, we can write a program $f \otimes f'$ that does these two jobs in parallel:

$$\begin{array}{ccc}
 \begin{array}{c} X \downarrow \\ \circlearrowleft f \\ Y \downarrow \end{array} & \begin{array}{c} X' \downarrow \\ \circlearrowleft f' \\ Y' \downarrow \end{array} & = & \begin{array}{c} X \otimes X' \downarrow \\ \circlearrowleft f \otimes f' \\ Y \otimes Y' \downarrow \end{array}
 \end{array}$$

These categories are are typically symmetric monoidal:



They're also cartesian. For example, we can write programs that duplicate data:

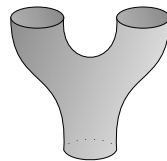
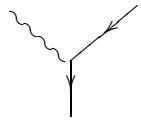
$$\Delta_X: X \rightarrow X \otimes X$$

But for quantum computation, we need programming languages that apply to *noncartesian* categories — because you can't duplicate quantum data!

And in quantum computation using anyons, the relevant categories are *braided*!

For more detail, read our paper in Bob Coecke's forthcoming book *New Structures in Physics*. You can find it now . You can find it now on the arXiv.

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$X \& Y \vdash Z$

$\lambda x \lambda y . x(y)$