We divide the remaining cases into four subcases according to whether  $\alpha + \alpha_i$ and  $\beta + \alpha_i$  are elements of  $\Phi$  or not, and show that  $e_i v_{\alpha,\beta} = 0$ . We first consider the case where one of  $\alpha + \alpha_i$ ,  $\beta + \alpha_i$  belongs to  $\Phi$ . As the argument is the same, we only treat the case where  $\alpha + \alpha_i \in \Phi$ . When  $\alpha + \alpha_i \in \Phi$ , we can write  $[E_{\alpha_i}, E_{\alpha}] =$  $c_{\alpha_1,\alpha}E_{\alpha_1+\alpha}$  with a non-zero number  $c_{\alpha_1,\alpha}$ . Thus,

$$[L_{\alpha_i}, L_{\alpha}] = e_i \iota(E_{\alpha}) = \iota(e_i E_{\alpha}) = \iota([E_{\alpha_i}, E_{\alpha}])$$
$$= c_{\alpha_i,\alpha} \iota(E_{\alpha_i+\alpha}) = c_{\alpha_i,\alpha} L_{\alpha_i+\alpha}.$$

This means that the element

$$[[L_{\alpha_i}, L_{\alpha}], L_{\beta}] - \iota([[E_{\alpha_i}, E_{\alpha}], E_{\beta}])$$

is equal to a scalar multiple of  $v_{\alpha+\alpha_1,\beta}$ . By the maximality of  $\gamma$ , this element is 0. Using this we have

$$e_i v_{\alpha,\beta} = [L_{\alpha}, [L_{\alpha_i}, L_{\beta}]] - \iota ([E_{\alpha}, [E_{\alpha_i}, E_{\beta}]]).$$

If  $\beta + \alpha_i \notin \Phi$ , then  $\beta + \alpha_i \neq 0$  implies that  $[E_{\alpha_i}, E_{\beta}] = 0$ . So

$$[L_{\alpha_i}, L_{\beta}] = e_i \iota(E_{\beta}) = \iota(e_i E_{\beta}) = \iota([E_{\alpha_i}, E_{\beta}]) = 0,$$

which implies that  $e_i v_{\alpha,\beta} = 0$ . If  $\beta + \alpha_i \in \Phi$ , then  $e_i v_{\alpha,\beta}$  is equal to a scalar multiple of  $v_{\alpha,\beta+\alpha_i}$  and the maximality of  $\gamma$  again implies that  $e_i v_{\alpha,\beta} = 0$ .

If both  $\alpha + \alpha_i$ ,  $\beta + \alpha_i$  are not in  $\overline{\Phi}$ , then

$$\begin{split} [L_{\alpha_i}, L_{\alpha}] &= e_i \iota(E_{\alpha}) = \iota(e_i E_{\alpha}) = \iota([E_{\alpha_i}, E_{\alpha}]) = 0, \\ [L_{\alpha_i}, L_{\beta}] &= e_i \iota(E_{\beta}) = \iota(e_i E_{\beta}) = \iota([E_{\alpha_i}, E_{\beta}]) = 0. \end{split}$$

Hence we have  $e_i v_{\alpha,\beta} = 0$  in this case also.

Since i is arbitrary and  $e_i v_{\alpha,\beta} = 0$ , Assertion 3(3) tells us that  $v_{\alpha,\beta}$  is equal to a scalar multiple of  $\iota(E_{1n})$ . In particular, we have

$$\gamma = \alpha_1 + \dots + \alpha_{n-1}$$

and  $\{\alpha,\beta\} = \{\alpha_1 + \cdots + \alpha_k, \alpha_{k+1} + \cdots + \alpha_{n-1}\}$  for some k. However,  $[L_{\alpha}, L_{\beta}]$  is equal to  $L_{1n}$  in this case and so  $v_{\alpha,\beta} = 0$ , which contradicts our choice of  $\gamma$ . Hence, we have  $A = \emptyset$  and Assertion 4 follows.

We are now ready to prove Theorem 2.2. We have constructed a map  $\iota:\mathfrak{g}\to U$ which satisfies  $[\iota(X), \iota(Y)] = \iota([X, Y])$ . Our next task is to prove the universality of the pair  $(U, \iota)$ ; however, this is obvious because the map  $\phi: U \to A$  is uniquely determined by the requirements that  $\phi(e_i) = \rho(E_{i,i+1})$  etc.

## 2.2. The quantum algebra of type $A_{r-1}$

Based on Theorem 2.2 Drinfeld and Jimbo introduced the quantum algebra which is obtained as a "deformation" of the enveloping algebra of  $\mathfrak{sl}_r = \mathfrak{sl}(r,\mathbb{C})$ . The definition is as follows. We choose  $\mathbb{Q}(v)$  as a base field since it is not necessary to assume it to be  $\mathbb{C}(v)$ . The element  $t_i$  is often denoted by  $v^{h_i}$  and  $\alpha_j(h_i) =$  $2\delta_{ij} - \delta_{i,j+1} - \delta_{i+1,j}$  by definition.

DEFINITION 2.5. Let  $K = \mathbb{Q}(v)$  where v is an indeterminate. The quantum algebra of type  $A_{r-1}$  is the unital associative K-algebra  $U_v(\mathfrak{sl}_r)$  defined by the following generators and relations.

Generators:  $t_i^{\pm 1}, e_i, f_i \quad (1 \le i \le r - 1).$ 

Relations:

$$\begin{split} t_i e_j t_i^{-1} &= v^{\alpha_j(h_i)} e_j, \quad t_i f_j t_i^{-1} = v^{-\alpha_j(h_i)} f_j, \\ [e_i, \, f_j] &= \delta_{ij} \frac{t_i - t_i^{-1}}{v - v^{-1}}, \\ [t_i, \, t_j] &= 0, \quad t_i t_i^{-1} = t_i^{-1} t_i = 1, \\ e_i^2 e_j - (v + v^{-1}) e_i e_j e_i + e_j e_i^2 = 0 \; (i - j = \pm 1), \\ e_i e_j &= e_j e_i \; (otherwise), \\ f_i^2 f_j - (v + v^{-1}) f_i f_j f_i + f_j f_i^2 &= 0 \; (i - j = \pm 1), \\ f_i f_j &= f_j f_i \; (otherwise). \end{split}$$

These relations are called the (deformed) Serre relations.

DEFINITION 2.6. Let  $[k] = \frac{v^k - v^{-k}}{v - v^{-1}}$ , for  $k \in \mathbb{N}$ , and let  $[n]! = \prod_{k=1}^n [k]$ . Then  $f_i^{(n)}$  is defined by

$$f_i^{(n)} = \frac{f_i^n}{[n]!}.$$

Roughly speaking, the quantum algebra is the algebra which is obtained by "integrating" the Cartan subalgebra and deforming the other relations "nicely".

We may obtain representations of  $U_v(\mathfrak{sl}_r)$  by deforming the representations of  $\mathfrak{g}$ . We can also define tensor product representations by deforming the coproduct of the enveloping algebra as follows.

$$\Delta(t_i) = t_i \otimes t_i, \ \Delta(e_i) = 1 \otimes e_i + e_i \otimes t_i^{-1},$$
  
$$\Delta(f_i) = f_i \otimes 1 + t_i \otimes f_i.$$

EXERCISE 2.7. Verify that  $\Delta$  defines an algebra homomorphism from  $U_v(\mathfrak{s}l_r)$  to  $U_v(\mathfrak{s}l_r)\otimes U_v(\mathfrak{s}l_r)$ .

EXERCISE 2.8. Let  $V = K^r$  and define  $\rho: U_v(\mathfrak{sl}_r) \to \operatorname{End}(V)$  by

$$\rho(t_i) = I + (v - 1)E_{i,i} + (v^{-1} - 1)E_{i+1,i+1},$$

$$\rho(e_i) = E_{i,i+1}, \quad \rho(f_i) = E_{i+1,i}.$$

Show that  $(\rho, V)$  is a representation of  $U_v(\mathfrak{s}l_r)$ . This is called the natural (or vector) representation of  $U_v(\mathfrak{s}l_r)$ .

EXERCISE 2.9. Let V be the natural representation of  $U_v(\mathfrak{s}l_2)$ . Decompose  $V\otimes V$  into a sum of irreducible  $U_v(\mathfrak{s}l_2)$ -submodules.