The following is extracted from Emmanual Hebey's 1999 book "Nonlinear Analysis on Manifolds: Sobolev Spaces and Inequalities." It spells out what is going with the $\nabla$ operator in more detail.

Let $D$ be an affine connection on a Riemmanian manifold. (In the seminar I was using $\nabla$ for $D$ to agree with Do Carmo, but $D$ is Aubin's and Hebey's and, I believe, Jerry's notation.) Define $\nabla_{i}$ by

$$
\nabla_{i}=D_{\frac{\partial}{\partial x^{i}}},
$$

so $\nabla_{i}$ maps a vector field into a vector field. Define the Christoffel symbols, $\left\{\Gamma_{i j}^{k}\right\}$ by

$$
\nabla_{i}\left(\frac{\partial}{\partial x^{j}}\right)=\Gamma_{i j}^{k} \frac{\partial}{\partial x^{k}} .
$$

If $Y=Y^{j} \frac{\partial}{\partial x^{j}}$ is a vector field, then by the Liebniz rule,

$$
\begin{aligned}
\nabla_{i} Y & =\frac{\partial Y^{j}}{\partial x^{i}} \frac{\partial}{\partial x^{j}}+Y^{j} \nabla_{i}\left(\frac{\partial}{\partial x^{j}}\right) \\
& =\frac{\partial Y^{j}}{\partial x^{i}} \frac{\partial}{\partial x^{j}}+Y^{j} \Gamma_{i j}^{k} \frac{\partial}{\partial x^{k}} \\
& =\left(\frac{\partial Y^{j}}{\partial x^{i}}+\Gamma_{i \alpha}^{j} Y^{\alpha}\right) \frac{\partial}{\partial x^{j}},
\end{aligned}
$$

where we re-indexed in the last step.
We can generalize the operation of $\nabla_{i}$ to a $(p, q)$-tensor $T$ (see my little spiel on tensors below, if you like) by

$$
\begin{aligned}
\left(\nabla_{i} T\right)_{i_{1} \cdots i_{p}}^{j_{1} \cdots j_{q}}=\frac{\partial T_{i_{1} \cdots i_{p}}^{j_{1} \cdots j_{q}}}{\partial x^{i}} & -\sum_{k=1}^{p} \Gamma_{i i_{k}}^{\alpha} T_{i_{1} \cdots i_{k-1}}^{j_{1} \cdots i_{k+1} \cdots i_{p}} \\
& +\sum_{k=1}^{q} \Gamma_{i \alpha}^{j_{k}} T_{i_{1} \cdots i_{p}}^{j_{1} \cdots j_{k-1} \alpha j_{k+1} \cdots j_{q}} .
\end{aligned}
$$

Our calculation for $Y=Y^{j} \frac{\partial}{\partial x^{j}}$ was the special case of a $(0, q)$-tensor-a vector (field).

We now define the $\nabla$ operator in such a way that it maps a $(p, q)$-tensor into a $(p+1, q)$ tensor (so it is really a series of operators, one for each $p$ ). For a smooth function $f$ on $M$ (a ( 0,0 )-tensor), we define $\nabla f=d f$ (a $(1,0)$-tensor). For a ( $p, q$ ) tensor $T, p \geq 1$, we define

$$
(\nabla T)_{i_{1} \cdots i_{p+1}}^{j_{1} \cdots j_{q}}=\left(\nabla_{i_{1}} T\right)_{i_{2} \cdots i_{p+1}}^{j_{1} \cdots j_{q}} .
$$

It makes sense to compose $\nabla$ operators: $\nabla^{k}=\nabla \cdots \nabla$ maps a $(p, q)$-tensor to a $(p+k, q)$-tensor.

For instance, if $f$ is a smooth function on a manifold,

$$
\left(\nabla^{2} f\right)_{i j}=\frac{\partial^{2} f}{\partial x_{i} \partial x_{j}}-\Gamma_{i j}^{k} \frac{\partial f}{\partial x_{k}},
$$

and is called the Hessian of $f$.

Hebey defines the norm of $\nabla^{k} u$ by

$$
\left|\nabla^{k} u\right|=g^{i_{1} j_{1}} \cdots g^{i_{k} j_{k}}\left(\nabla^{k} u\right)_{i_{1} \cdots i_{k}}\left(\nabla^{k} u\right)_{j_{1} \cdots i_{k}}
$$

It appears, though, that this is actually $\left|\nabla^{k} u\right|^{2}$, which I will assume is the case (I don't think it would be a norm, otherwise). Note that this is a definition, not a calculation, since $\nabla^{k} u$ is a vector field only for $k=1$.

Hebey's $\left|\nabla^{k} u\right|^{2}$ (assuming I am right about the square), then, would be the same as Aubin's definition of $\left|\nabla^{k} \varphi\right|^{2}$ on p. 32 (with $\varphi$ in place of $u$ ) if by $\nabla_{\alpha_{1}} \nabla_{\alpha_{2}} \cdots \nabla_{\alpha_{k}} \varphi$ Aubin means what Hebey means by $\left(\nabla^{k} \varphi\right)_{\alpha_{1} \alpha_{2} \cdots \alpha_{k}}$ and if Aubin means to have the indices raised after all the derivatives have been performed (or if raising the indices somehow commutes with the derivatives). Aubin's terse definition 1.14 on p. 4 seems to indicate this may be the case. If so, then Hebey's definition of the Sobolev norm is identical to that of Aubin's.

In any case, when $u$ is a real-valued function on the manifold-a $(0,0)$ -tensor- $\nabla^{k} u$ is a $(k, 0)$-tensor. This looks a lot like a $k$-form, but is not antisymetric; in any case, it is only its coefficients in a coordinate system that enter into the Sobolev norms, and these coefficients contain all the derivatives of order $k$ along with lower-order derivatives. In the special case of a connection with all of $\Gamma_{j k}^{i}=0$, the coefficients would consist exactly of all derivatives of order $k$, and in flat space, the definition of the Sobolev space norm that Aubin gives would coincide with the usual definition.

## Tensors

A $(p, q)$-tensor is a multilinear form that maps, for each $x \in M$ the space

$$
T_{x}(M) \times \cdots \times T_{x}(M) \times T_{x}(M)^{*} \times \cdots \times T_{x}(M)^{*}
$$

into $\mathbb{R}$, where there are $p$ products of $T_{x}(M)$ and $q$ products of $T_{x}(M)^{*}$. A basis element for the space of $(p, q)$-tensors looks like

$$
d x^{i_{1}} \otimes \cdots \otimes d x^{i_{p}} \otimes \frac{\partial}{\partial x^{j_{1}}} \otimes \cdots \otimes \frac{\partial}{\partial x^{j_{q}}}
$$

if that helps clear things up. We can write a $(p, q)$-tensor, $T$, in coordinates as

$$
T_{i_{1} \cdots i_{p}}^{j_{1} \cdots j_{q}} d x^{i_{1}} \otimes \cdots \otimes d x^{i_{p}} \otimes \frac{\partial}{\partial x^{j_{1}}} \otimes \cdots \otimes \frac{\partial}{\partial x^{j_{q}}}
$$

but we usually just write the coefficients $T_{i_{1} \cdots i_{p}}^{j_{1} \cdots j_{q}}$ if the coordinate system is understood.

A $(p, q)$-tensor transforms under changes of coordinates from $x$ to $y$ as follows:

$$
\widetilde{T}_{i_{1} \cdots i_{p}}^{j_{1} \cdots j_{q}}=T_{\alpha_{1} \cdots \alpha_{p}}^{\beta_{1} \cdots \beta_{q}} \frac{\partial x^{\alpha_{1}}}{\partial y^{i_{1}}} \cdots \frac{\partial x^{\alpha_{p}}}{\partial y^{i_{p}}} \frac{\partial y^{j_{1}}}{\partial x^{\beta_{1}}} \cdots \frac{\partial x^{j_{q}}}{\partial y^{\beta_{q}}}
$$

