

2/26/02

## Maxwell's Eqns

$$dF = 0$$

$$*d*F = J$$

2-form  $F$  electromagnetic field

First eqn will hold automatically if  $F = dA$   
- This looks like a trick! Let's understand what  $A$  is.

Quest: Is  $A$  "fundamental" or is  $F$ ?

Note: If spacetime is  $\mathbb{R}^4$ , it doesn't matter which we use ( $A$  or  $F$ ) because any 2-form  $F$  w/  $dF = 0$  is of the form  $dA$  for some 1-form  $A$ .

Now - is  $A$  unique?

Well - while  $A$  is not unique, any 2 1-forms  $A, A'$  with

Actually - we  $dA = dA'$

are related by

$$A' - A = df$$

for some 0-form  $f$ .

(Easy part: If  $A' - A = df$ , then  $dA' - dA = ddf = 0$ .)

Hard part: If  $dA = dA'$  then  $A' - A = df$  for some  $f$ .

This part depends on the topology of our spacetime.

ie) This fails for spacetimes of (other) topologies.

Moreover - on some manifolds,  $dF=0 \not\Rightarrow F=dA$

True for  $\mathbb{R}^n$ : if  $dB=0$  then  $B=d\Omega$  for some  $\Omega$ .

De Rham cohomology: If  $M$  is a smooth manifold we get:

$$\Omega^0(M) \xrightarrow{d} \Omega^1(M) \xrightarrow{d} \Omega^2(M) \xrightarrow{d}$$

with  $d^2=0$ , which implies

$$* \quad \mathbb{Z}^p(M) = \{w \in \Omega^p(M) \mid dw=0\}$$

"cocycles" cycle - boundary is zero.

$d$  is called a 'coboundary' map.

$d$  boundary-map sends from  $p$ -space to  $p-1$  space.

$$* \quad \mathbb{B}^p(M) = \{w \in \Omega^p(M) \mid w = d\alpha \text{ for some } \alpha \in \Omega^{p-1}(M)\}$$

"coboundaries" (are  $d$  of something)

$$d^2=0 \Rightarrow \mathbb{B}^p(M) \subseteq \mathbb{Z}^p(M).$$

but sometimes this is a strict inclusion  $\subset$ .

$$\text{let } H^p(M) = \frac{Z^p(M)}{B^p(M)}$$

"de Rham cohomology"

This tells us something about the holes in space.

If  $M = \mathbb{R}^n$ , then  $H^p(M) = 0 \quad \forall p > 0$ .

Note:  $df = 0 \Rightarrow f$  is constant

But in general  $H^p(M)$  is an interesting topological invariant of  $M$  - finite dim'l if  $M$  compact.

So - for  $\mathbb{R}^4$  it "doesn't matter" whether we use  $A$  or  $F$ , but  $A$  is nice if we're interested in...

sidenote: ( $A = 1$ -form, integrate along world line of particle - action)

**Action**

Suppose we have a particle moving along a path in space:

$$g: \begin{array}{ccc} [0, T] & \longrightarrow & \mathbb{R}^3 \\ \text{time} & & \text{space} \end{array}$$

We want to know how our particle moves.  
We phrase the question as:

Given  $q(0)$ ,  $q(T)$ , what is the particle's path?  
(How does it get from time 0 to time T).

It minimizes some quantity called the action,  $S$ ,  
which is of the form

$$S = \int_0^T \mathcal{L}(q(t), \dot{q}(t)) dt$$

where  $\mathcal{L}$  is some function

$$\mathcal{L}: \mathbb{R}^3 \times \mathbb{R}^3 \longrightarrow \mathbb{R}$$

$\downarrow$                        $\downarrow$   
 $q$                        $\dot{q}$

(or if you replace  $\mathbb{R}^3$  by a manifold  $M$ ,  
 $\mathcal{L}: TM \longrightarrow \mathbb{R}$ .)

tangent bundle of  $M$

Actually the particle will not minimize  $S$ , but  
will merely achieve a critical point:

$$\delta S = 0$$

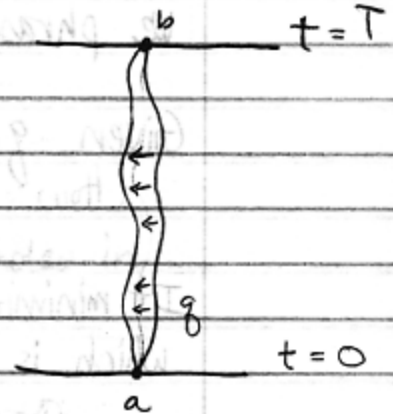
(i.e.  $dS = 0$  where  $S$  is a function  $S: P \longrightarrow \mathbb{R}$ )

$$P = \{q: [0, T] \longrightarrow \mathbb{R}^3 \mid q(0) = a, q(T) = b\}$$

where

a path, not  
particle

$$\lim_{\epsilon \rightarrow 0} \frac{S(q + \epsilon \delta q) - S(q)}{\epsilon} = \delta S$$



$$\delta q: [0, T] \rightarrow \mathbb{R}^3$$

$$\delta q(0) = \delta q(T) = 0$$

So— we want  $\delta S = 0$   
for all  $\delta q$  st  $\delta q(0) = \delta q(T) = 0$ .

Which paths make  $\delta S = 0$  in this sense?

$$0 = \delta S = \delta \int_0^T \mathcal{L}(q(t), \dot{q}(t)) dt$$

This can  
change in 2 ways:  
 $q(t)$  or  $\dot{q}(t)$

$$= \int_0^T \delta \mathcal{L}(q(t), \dot{q}(t)) dt$$

$$= \int_0^T \frac{\partial \mathcal{L}}{\partial q^i} \delta q^i(t) + \frac{\partial \mathcal{L}}{\partial \dot{q}^i} \delta \dot{q}^i(t) dt$$

Einstein  
sum notation  
in denominator,  
upstairs is  
downstairs for  
indices

use  
integration by parts

here we're using Minkowski metric:  $-dt^2 + dx^2 + dy^2 + dz^2$   
 want to minimize distance of this metric

so write  $\delta q^i(t) = \delta \frac{dq^i}{dt}(t) = \frac{d}{dt} \delta q^i(t)$

$$= \int_0^T \left( \frac{\partial \mathcal{L}}{\partial q^i} - \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}^i} \right) \delta q^i dt$$

This can only vanish  $\forall \delta q$  if

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}^i} = \frac{\partial \mathcal{L}}{\partial q^i}$$

Euler-Lagrange eqn.

$$\underbrace{\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}^i}}_{\substack{i^{\text{th}} \text{ component} \\ \text{of momentum}}} = \underbrace{\frac{\partial \mathcal{L}}{\partial q^i}}_{i^{\text{th}} \text{ component of force}}$$

Example:

1) Newtonian particle in a potential:

$$\mathcal{L}(q(t), \dot{q}(t)) = \text{kinetic energy} - \text{potential energy}$$

note - action is what is happening  
 so - kinetic energy - what's happening  
 potential energy - what's not happening / could happen

(+) So, for example, a rock on a cliff has lots of potential energy, so negative action.

$\mathcal{L}(q(t), \dot{q}(t)) = \text{kinetic energy} - \text{potential energy}$

$$= \frac{m}{2} \dot{q}^2 - V(q(t))$$

$m > 0$  mass.  $\dot{q}^i \dot{q}_i$  "potential"  
 $V: \mathbb{R}^3 \rightarrow \mathbb{R}$

$$\text{momentum} = \frac{d\mathcal{L}}{d\dot{q}^i} = m\dot{q}_i$$

$$\frac{d}{dt} \frac{d\mathcal{L}}{d\dot{q}^i} = m\ddot{q}_i = \text{mass} \cdot \text{acceleration}$$

$$\frac{d\mathcal{L}}{dq^i} = -\frac{dV}{dq^i} = \text{force}$$

so we get  $F = ma$ .

2) Special relativistic particle in electromagnetic field.

Now - we won't view our particle's path as space & time parts.

Now - our particle's path will be

$$q: [0, T] \rightarrow \mathbb{R}^4$$

"parameter space"

spacetime

Here we're using Minkowski metric:  $-dt^2 + dx^2 + dy^2 + dz^2$   
 want to minimize distance w/ this metric

and action is:

$$S = \underbrace{-m}_{\text{mass}} \cdot \text{length of path} + e \int_{\underbrace{q}_{\text{our path}}} A$$

or "total proper time" of path

$$S = -m \int_0^T \sqrt{-\frac{dq^i}{ds} \cdot \frac{dq_i}{ds}} ds$$

particle like this  
to be big

inner product

$$+ e \int_0^T A_i(q(s)) \frac{dq^i}{ds} ds$$

$\dot{q}^i = \frac{dq^i}{ds}$

$i = 0, 1, 2, 3$  so,  $\mathcal{L}(q, \dot{q}) = -m \sqrt{-\dot{q}^i \dot{q}_i} + e A_i(q) \dot{q}^i$

momentum =  $\frac{\partial \mathcal{L}}{\partial \dot{q}^i} = m \cdot \frac{\dot{q}^i}{\sqrt{-\dot{q}^i \dot{q}_i}} + e A_i(q)$

S doesn't (strongly) depend on the path!

$$S(q_1) = S(q_2) \text{ if } q_1 \text{ and } q_2 \text{ are paths}$$

$$q_1: [a, T] \rightarrow \mathbb{R}^4 \text{ and } q_2: [0, T] \rightarrow \mathbb{R}^4$$