

"Violin String Theory":

Start out with an infinitely long string. Its motion satisfies the wave equation:

$$\varphi : \mathbb{R}^2 \longrightarrow \mathbb{R}$$

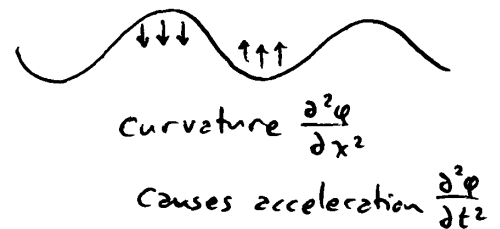
ψ
 (t, x)

time position

height of string

$$\boxed{\frac{\partial \varphi}{\partial t^2} = \frac{\partial^2 \varphi}{\partial x^2}}$$

Acceleration Curvature

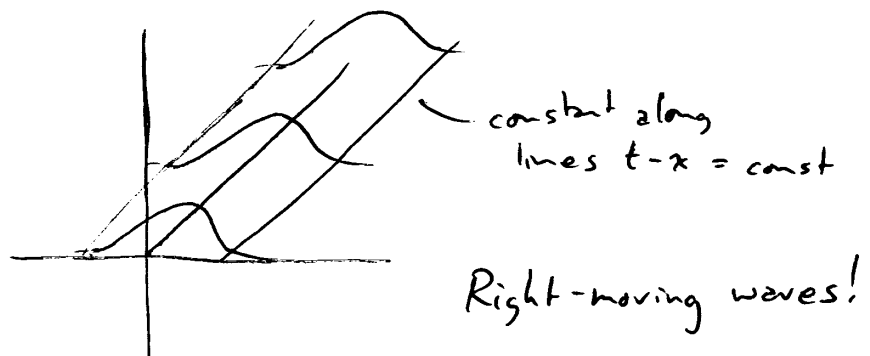


This has some obvious solutions:

$$\varphi(t, x) = f(t-x)$$

$$\begin{array}{ccc}
 & \frac{\partial}{\partial t} & \frac{\partial}{\partial x} \\
 & \swarrow & \searrow \\
 f'(t-x) & & -f'(t-x) \\
 & \swarrow \frac{\partial}{\partial t} & \searrow \frac{\partial}{\partial x} \\
 & f''(t-x) &
 \end{array}$$

Solutions look like:



We also have left-moving waves:

$$\varphi(t, x) = f(t+x)$$

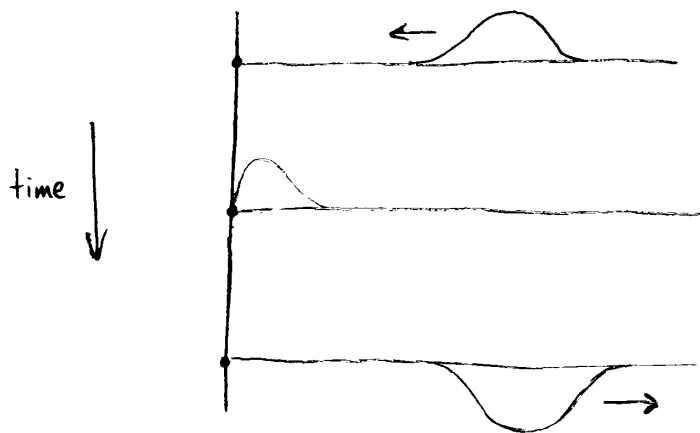
Thm: Every smooth solution of the wave equation on $\mathbb{R} \times \mathbb{R}$ is of the form

$$\varphi(t, x) = f(t+x) + g(t-x)$$

for some $f, g \in C^\infty(\mathbb{R})$.

(Proof: Exercise!)

Suppose we want a solution of the wave equation on $\mathbb{R} \times [0, \infty)$ such that $\varphi(t, 0) = 0$ for all $t \in \mathbb{R}$:



We can find such solutions by taking:

$$\varphi(t, x) = f(t+x) - f(t-x)$$

This satisfies the wave equation on $\mathbb{R} \times \mathbb{R}$ and has $\varphi(t, -x) = -\varphi(t, x)$ so $\varphi(t, 0) = 0$ & φ satisfies wave equation on $\mathbb{R} \times [0, \infty)$.

Previous theorem gives:

Cor: Every smooth soln. of the wave eq. on $\mathbb{R} \times [0, \infty)$ with $\varphi(t, 0) = 0$ is of the form

$$\varphi(t, x) = f(t+x) - f(t-x)$$

for some $f \in C^\infty(\mathbb{R})$.

(Proof: The space of odd functions on \mathbb{R} is isomorphic to the space of functions on $[0, \infty)$ that vanish at 0.)

We can also get some nice complex-valued solutions of this type, taking

$$f(t) = e^{i\omega t} \quad t \in \mathbb{R}$$

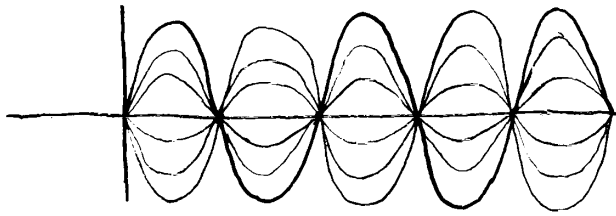
The real & imaginary parts of a complex solution will be real solutions.

$$\begin{aligned} \varphi(t, x) &= e^{i\omega(t+x)} - e^{i\omega(t-x)} \\ &= e^{i\omega t} (e^{i\omega x} - e^{-i\omega x}) \\ &= 2i e^{i\omega t} \sin \omega x \end{aligned}$$

So

$$\begin{aligned} \operatorname{Re} \varphi(t, x) &= -2 \sin \omega t \sin \omega x \\ \operatorname{Im} \varphi(t, x) &= 2 \cos \omega t \sin \omega x \end{aligned}$$

So we get something like



These are standing waves with nodes at points where $\sin \omega x = 0$, i.e.

$$x = k \frac{\pi}{\omega} \quad k = 0, 1, 2, \dots$$

Now suppose we want standing waves that have nodes at $x=0$ & $x=\pi$ i.e.

$$\begin{aligned} \varphi(t, 0) &= 0 \\ \varphi(t, \pi) &= 0 \end{aligned} \quad \forall t \in \mathbb{R}$$

Then we need $\pi = k \frac{\pi}{\omega}$, or $\omega = k = 0, 1, 2, 3, \dots$

(if $\omega < 0$ we just get minus one of these solutions)

This gives solutions:

$$\varphi(t, x) = 2i e^{ikt} \sin kx$$

or real solutions:

$$\cos kt \cdot \sin kx \quad \& \quad \sin kt \cdot \sin kx.$$

These are nonzero for $k \in \mathbb{N}^+$

Thm: The space V of smooth real solutions of the wave equation on $\mathbb{R} \times [0, \pi]$ with $\varphi(t, 0) = \varphi(t, \pi) = 0$ has $\{\underbrace{\cos kt \sin kx}_{= e_k}, \underbrace{\sin kt \sin kx}_{= f_k}\}$ as a topological basis, i.e. every linear comb.

$$\sum_{k \in \mathbb{N}^+} a_k e_k + b_k f_k \quad \text{where } a_k, b_k \rightarrow 0 \text{ faster than}$$

$\frac{1}{k^p}$ for any power of p converges to a point in V (with its C^∞ topology), and every point of V is of this form for a unique choice of $\{a_k\}, \{b_k\}$.

$$\begin{array}{l} f_i \rightarrow f \text{ in } C^\infty \text{ top.} \\ \text{if} \\ \sup_{t,x} |\partial_x^n \partial_x^m (f_i - f)| \rightarrow 0 \\ \forall n, m \end{array}$$

topological direct sum =
closure of algebraic direct sum.

Note that $V = \bigoplus_{k \in \mathbb{N}^+} V_k$ where V_k has basis e_k, f_k .

We have linear operators describing time evolution:

$$U(t): V \longrightarrow V \quad \forall t \in \mathbb{R}$$

given by:

$$(U(t)\varphi)(s, x) = \varphi(t+s, x)$$

&

$$U(t): V_k \longrightarrow V_k$$

So: our physics problem here is a direct sum of little physics problems — harmonic oscillators with different frequencies!

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We saw that the space of smooth solns of the wave equation on $\mathbb{R} \times [0, \pi]$ with Dirichlet boundary conditions:

$$\varphi(t, 0) = \varphi(t, \pi) = 0 \quad \forall t$$

has a topological basis given by

$$e_k = \cos kt \sin kx$$

$$f_k = \sin kt \sin kx$$

Let

$$V_k = \{ae_k + bf_k : a, b \in \mathbb{R}\}$$

and now let's work purely algebraically & redefine V to be the space of finite linear combinations of e_k, f_k (so we're taking a smaller space, but it's dense in the original one) so

$$V = \bigoplus_{k \in \mathbb{N}^+} V_k$$

where now this is an algebraic direct sum.

Time evolution

$$U(t) : V \longrightarrow V$$

preserves each subspace V_k :

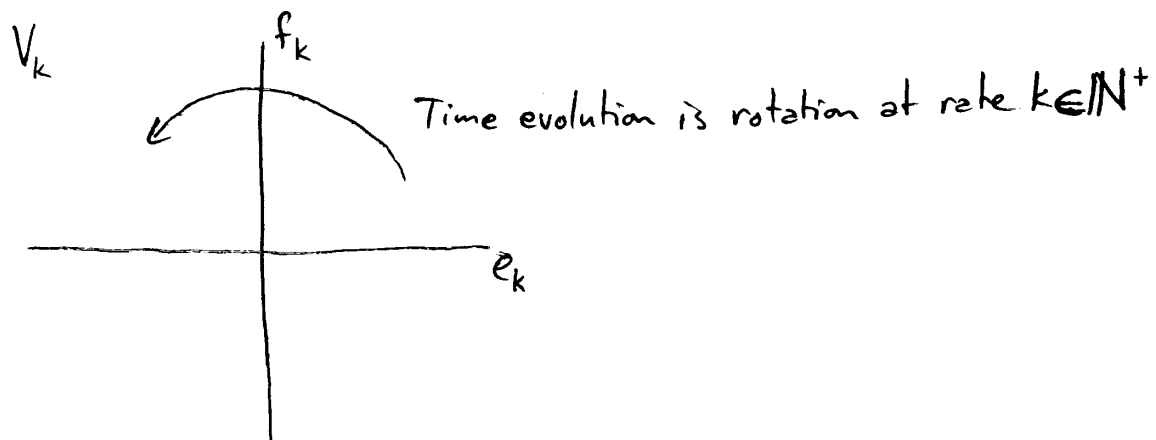
$$U(t) : V_k \longrightarrow V_k$$

$$\begin{aligned}
 U(s) e_k &= \cos k(t+s) \sin kx \\
 &= (\cos ks \cos kt - \sin ks \sin kt) \sin kx \\
 &= \cos ks e_k - \sin ks f_k
 \end{aligned}$$

$$\begin{aligned}
 U(s) f_k &= \sin k(t+s) \sin kx \\
 &= (\sin ks \cos kt + \cos ks \sin kt) \sin kx \\
 &= \sin ks e_k + \cos ks f_k
 \end{aligned}$$

So

$$U(s) \Big|_{V_k} = \begin{pmatrix} \cos ks & -\sin ks \\ \sin ks & \cos ks \end{pmatrix}$$



Just as we did with the harmonic oscillator, let's make V_k into a 1-dim complex vector space with

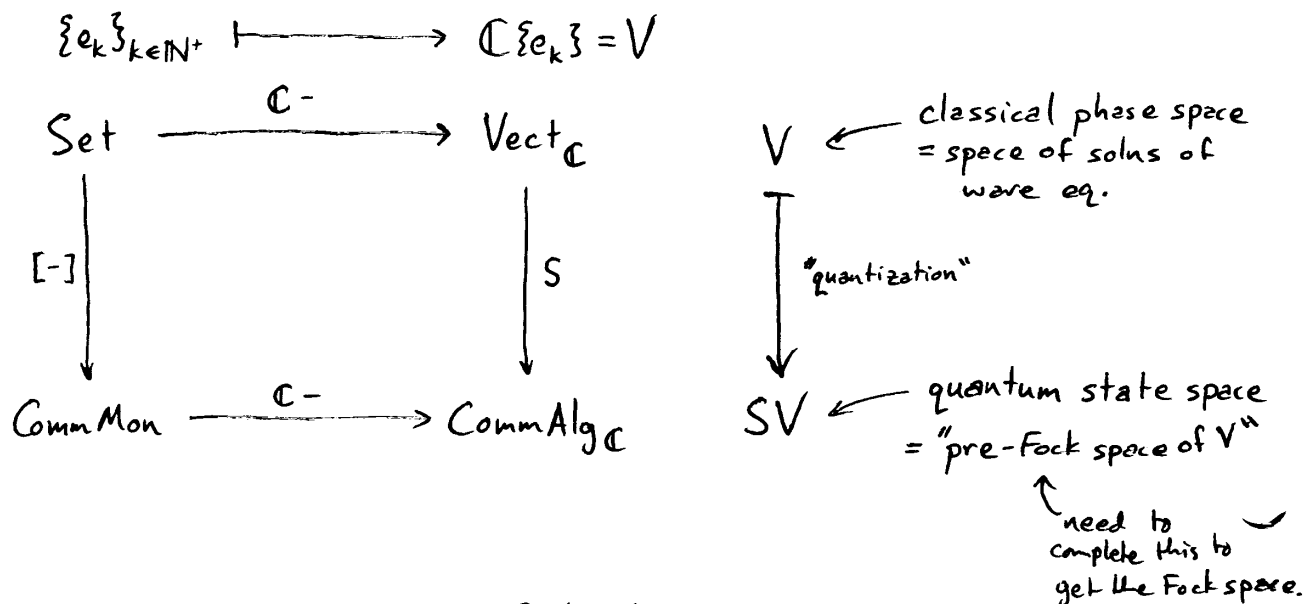
$$\begin{aligned}
 i e_k &= f_k \\
 i f_k &= -e_k
 \end{aligned}$$

Then we get

$$U(s) \Big|_{V_k} = e^{iks}$$

Now, let's QUANTIZE the string!

Recall:



V has basis of e_k 's so SV has a basis of monomials:

$$e_1^{p_1} e_2^{p_2} \dots e_k^{p_k}$$

where $k \in \mathbb{N}^+$ is arbitrary & $p_i \in \mathbb{N}$. Physically this is the state with " p_i particles of type i ". The monomial 1 is also called the vacuum.

To quantize time evolution for the string, we use the fact that $S: \text{Vect}_{\mathbb{C}} \rightarrow \text{Comm Alg}_{\mathbb{C}}$ is a functor: it not only sends objects (vector spaces)

to objects (comm. algs.), but morphisms (linear maps) to morphisms (algebra homomorphisms)! So, time evolution for the classical string:

$$U(t): V \longrightarrow V$$

gets sent to

$$SU(t): SV \longrightarrow SV$$

which is time evolution for the quantum string!

Here's how it goes: Given a linear operator

$$f: L \longrightarrow M$$

we have

$$Sf: SL \longrightarrow SM$$

$$l_1 \dots l_n \longmapsto f(l_1) \dots f(l_n) \quad l_i \in L$$

& Sf is linear, hence determined on all of SL - in fact a homomorphism.

In particular, $SU(t): SV \longrightarrow SV$ is given by:

$$\begin{aligned} SU(t)(e_1^{p_1} e_2^{p_2} \dots e_k^{p_k}) \\ &= (U(t)e_1)^{p_1} (U(t)e_2)^{p_2} \dots (U(t)e_k)^{p_k} \\ &= (e^{it} e_1)^{p_1} (e^{i2t} e_2)^{p_2} \dots (e^{ikt} e_k)^{p_k} \end{aligned}$$

$$= e^{it(1p_1+2p_2+\dots+kp_k)} e_1^{p_1} e_2^{p_2} \dots e_k^{p_k}$$

In fact

$$SU(t) = e^{itH}$$

where $H: SV \rightarrow SV$, the Hamiltonian for the quantized string, is:

$$H e_1^{p_1} \dots e_k^{p_k} = (1p_1 + 2p_2 + \dots + kp_k) e_1^{p_1} \dots e_k^{p_k}$$

Moral: the state $e_1^{p_1} \dots e_k^{p_k}$ is an energy eigenstate with energy $E = 1p_1 + \dots + kp_k$. So: the i th type of particle has energy i !!

Now: How many eigenstates are there with energy $n \in \mathbb{N}$?

We'll call this number $P(n)$, the n th partition number

$P(0) = 1$	(vacuum) $0 =$
$P(1) = 1$	$1 = 1$
$P(2) = 2$	$2 = 1 + 1, 2$
$P(3) = 3$	$3 = 1 + 1 + 1, 1 + 2, 3$
$P(4) = 4$	$4 = 1 + 1 + 1 + 1, 1 + 1 + 2, 1 + 3, 2 + 2, 4$
$P(5) = 7$	$5 = 1 + 1 + 1 + 1 + 1, 1 + 1 + 1 + 2, 1 + 1 + 3, 1 + 2 + 2, 1 + 4, 2 + 3, 5$
\vdots	

← Ramanujan studied these

The generating function for the $P(n)$ is, by the homework:

$$\sum_{n \in \mathbb{N}} P(n) z^n = \prod_{k \in \mathbb{N}^+} \frac{1}{1 - z^k},$$

the partition function of the string.

Ramanujan studied these & showed lots of bizarre things, e.g.:

$$P(5m+4) \equiv 0 \pmod{5}$$

How? He defined

$$\phi(z) = \prod_{k=1}^{\infty} (1 - z^k)$$

& he showed:

$$\sum_{m \geq 0} P(5m+4) z^m = 5 \underbrace{\frac{\phi(z^5)^5}{\phi(z)^6}}_{\text{where this} = \sum a_n z^n \text{ w. } a_n \in \mathbb{N}!}$$

These curiosities can be more systematically obtained via string theory!

