

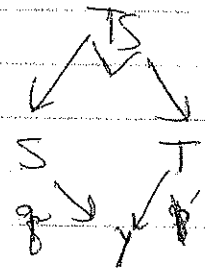
AN/2008

John Baez

a) Given a span of groupoids $\begin{matrix} & S & \\ p \swarrow & & \searrow q \\ X & & Y \end{matrix}$ (or S), we get a linear operator $\tilde{S} : H_0(X) \rightarrow H_0(Y)$
 $\tilde{S} = q_* \circ p^!$

b) Given composable spans $\begin{matrix} & S & & T & \\ p \swarrow & & \searrow q & & \searrow q' \\ X & & Y & & Z \end{matrix}$

we compose them via weak pullback



$$TS := \{ (s \in S, t \in T) \mid q(s) \cong p'(t) \}$$

Theorem

$$\tilde{TS} = \tilde{T} \circ \tilde{S}$$

c) We also have an identity span $1_X = \begin{matrix} & X & \\ 1_1 \swarrow & & \searrow 1_2 \\ X & & X \end{matrix}$. Then

Theorem

$$\tilde{1}_X = 1_{H_0(X)}$$

d) Adding spans? Given spans $\begin{matrix} & S & \\ & \swarrow & \searrow \\ X & & Y \end{matrix}$ & $\begin{matrix} & T & \\ & \swarrow & \searrow \\ X & & Y \end{matrix}$

there's an obvious way to construct $\begin{matrix} & S \sqcup T & \\ & \swarrow & \searrow \\ X & & Y \end{matrix}$

and this is denoted as $S+T$. So,

$S+T = S \amalg T$, with the same equation on the level of objects as well as morphisms!

- Theorem**
- ① $S+T \cong T+S$ (as groupoids/spans)
 - ② $\widetilde{S+T} = \widetilde{S} + \widetilde{T}$. ($= \widetilde{T} + \widetilde{S} = \widetilde{T+S}$).

② Remark: In adding spans, we're taking "OR" (logic) - which is what happens in quantum mechanics when one applies a linear operator to a state.

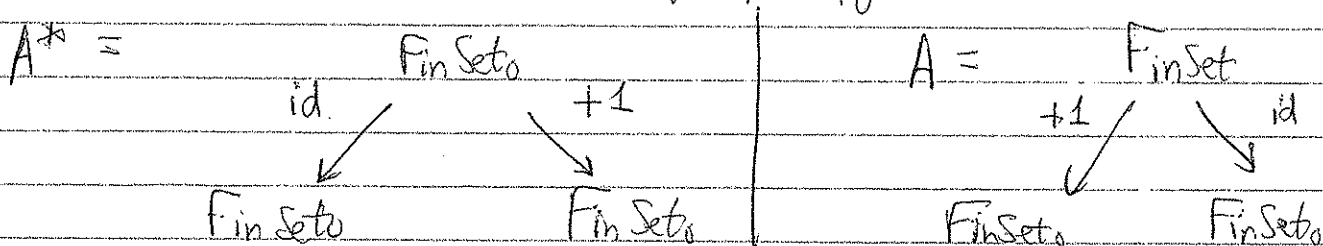
One gets superposition of states - and this is good.

"analogy" since we are trying to groupoidify quantum mechanics.

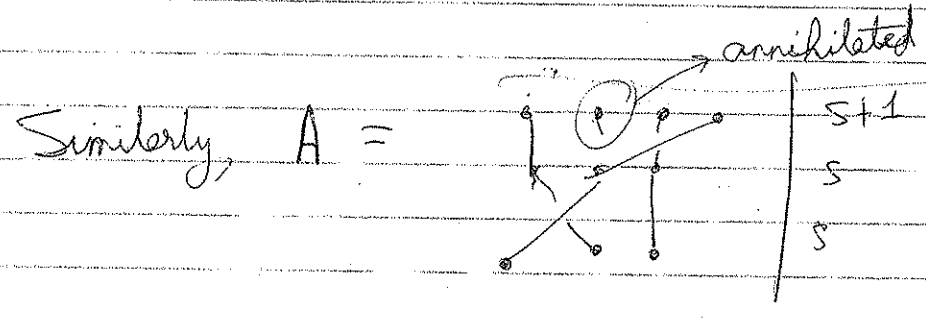
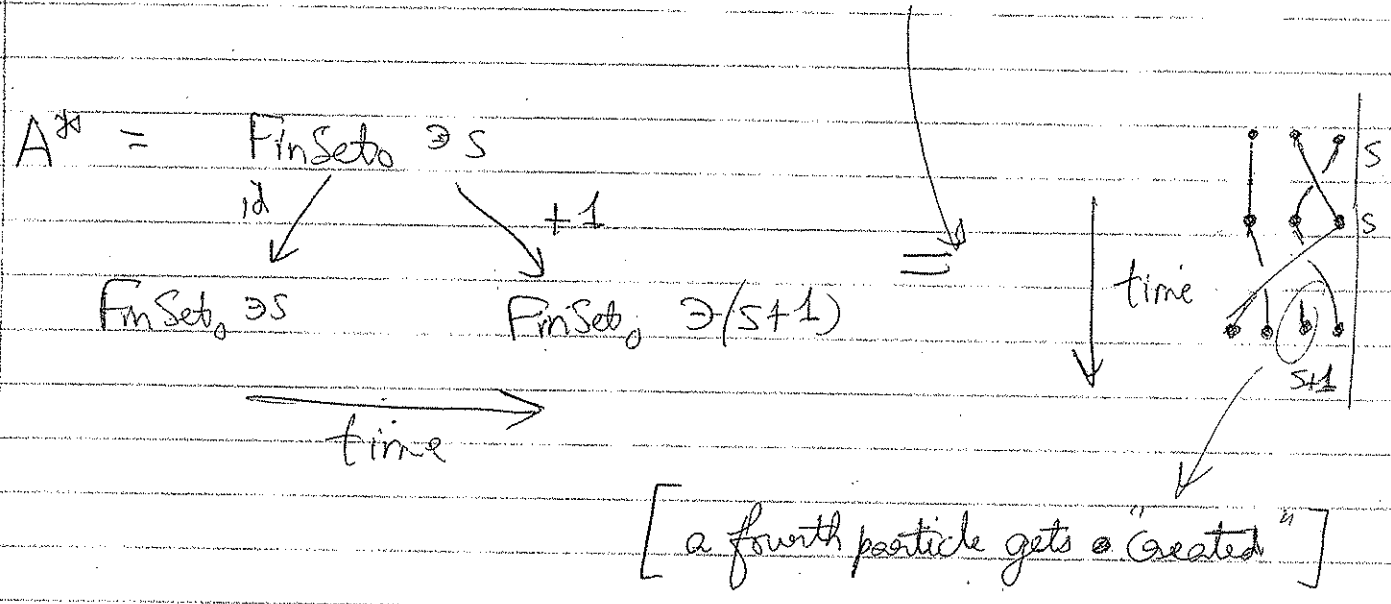
② @ Back to the Annihilation and Creation Operators

We have: $H_0(\text{FinSet}_0) \cong k[\mathbb{Z}]$ (k a ground field of char. 0)

and certain spans (that groupoidify a, a^*):



b) which we can draw as Feynman diagrams:



a) Now let's show : $AA^* \cong A^*A + Id$.

(and thus, $\widetilde{AA^*} = \widetilde{a^*a}$
 $\widetilde{AA^* + Id} = \widetilde{a^*a + Id}$)

Note We see using some more functorial properties, eg:

$S \underline{N} T \Rightarrow \widetilde{S} = \widetilde{T}$, and:

Moreover, we have the notion of the adjoint span:

Given $\left[\begin{array}{c} S \\ \swarrow \downarrow \searrow \\ X \end{array} \right] = S$ / the adjoint span is defined

to be $\left[\begin{array}{c} S^* (=S) \\ \swarrow \downarrow \searrow \\ Y \quad X \end{array} \right] = S^*$

Then Thm (1) All $H_0(X)$'s have inner products on them hence linear operators have adjoints!

(2) Thus, we have $H_0(X) \xrightarrow{\tilde{S}} H_0(Y)$

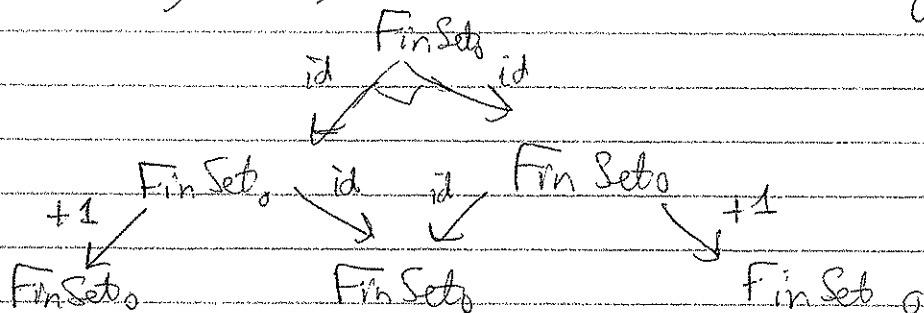
$\downarrow (\tilde{S})^*$
 $H_0(Y) \xrightarrow{(\tilde{S})^*} H_0(X)$

and we also have $H_0(Y) \xrightarrow{(\tilde{S}^*)} H_0(X)$.

Then, $\tilde{S}^* = \tilde{S}^*$

(3) $(S^*)^* = S$, $(\tilde{S}^*)^* = \tilde{S}$.

(c) We start, over, with A^*A (the "easy composite"):



The groupoid on top is the weak pullback:

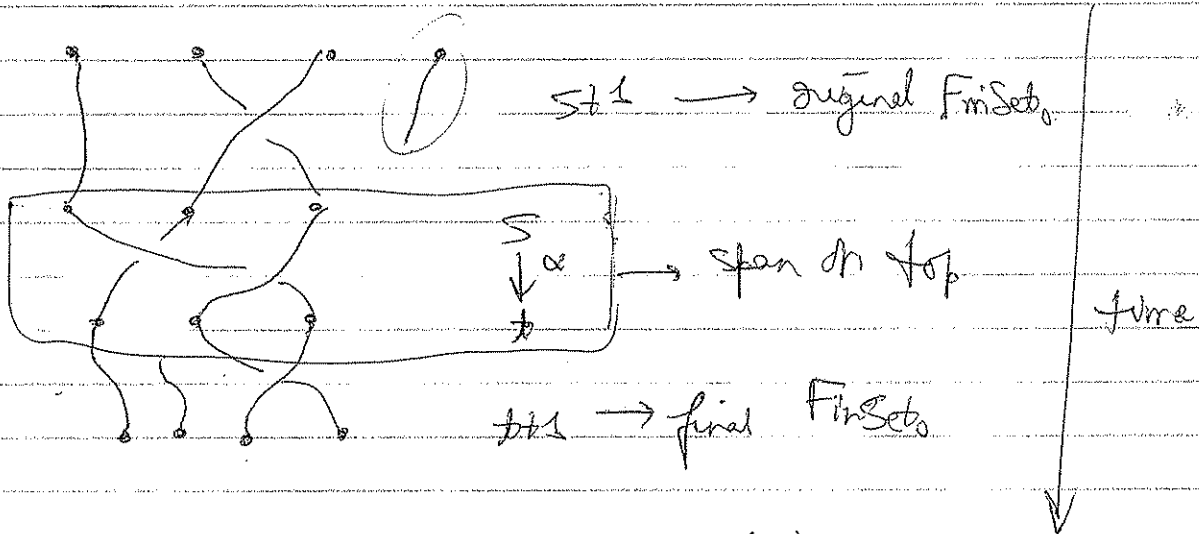
$$[s, t \in \text{FinSet}_0, \alpha: s \xrightarrow{\sim} t]$$

||?

FinSet₀

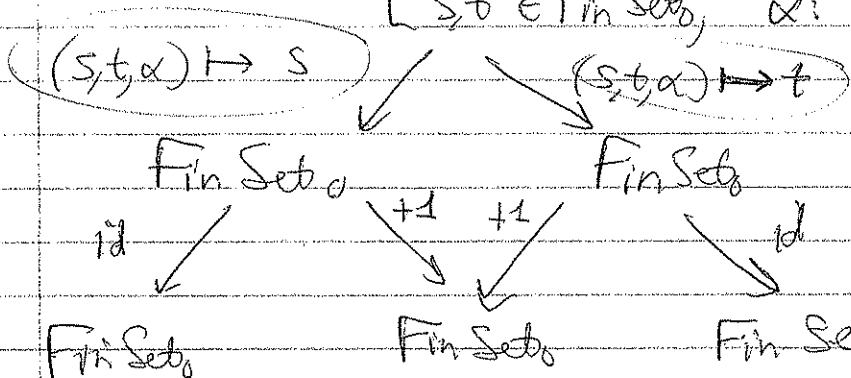
————— (*)

Let's now draw the Feynman diagram, but not using any simplification with, but using pairs of (isom.) finite sets.



This is how the composite span's elements "look".

$$[s, t \in \text{FinSet}_0, \alpha: s+1 \xrightarrow{\sim} t+1]$$



d) To show our desired commutation relation on the groupoid level, we ~~start to~~ first want to write AA^* as a sum of two spans:

Before we do this on the span level, let's do it at the groupoid level on top only:

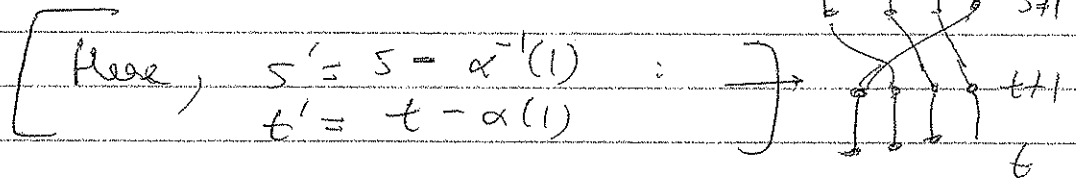
$$[s, t \in \text{FinSet}_0, \alpha: s+1 \xrightarrow{\sim} t+1]$$

$$\cong [s, t \in \text{FinSet}_0 : \alpha: s \xrightarrow{\sim} t]$$

$$+ [s, t \in \text{FinSet}_0, \alpha: s+1 \xrightarrow{\sim} t+1 \text{ s.t. } \alpha(1) \neq 1]$$

where $\checkmark \cong [s', t' \in \text{FinSet}_0 : \alpha: s' \xrightarrow{\sim} t']$

$$\cong \text{FinSet}_0$$



e) Now on the span level:

$$AA^* \text{ is } [s, t \in \text{FinSet}_0, \alpha: s+1 \xrightarrow{\sim} t+1]$$

$$(s, t, \alpha) \mapsto s \quad \swarrow \quad \searrow \quad (s, t, \alpha) \mapsto t$$

FinSet_0

FinSet_0

This span is equivalent to the sum:

$$[(s, t) \in \text{Fin Set}_0, \alpha: s \xrightarrow{\sim} t]$$

$$(s, t, \alpha) \mapsto s$$

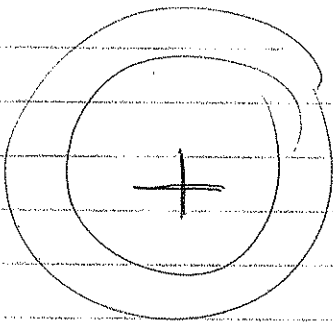
$$\downarrow$$

$$\text{Fin Set}_0$$

$$(s, t, \alpha) \mapsto t$$

$$\downarrow$$

$$\text{Fin Set}_0$$



(and this is $\cong \text{Fin Set}_0 = \text{id}$)

$$\begin{array}{ccc} & \text{id} & \\ & \swarrow & \searrow \\ \text{Fin Set}_0 & & \text{Fin Set}_0 \end{array}$$

$$[(s', t') \in \text{Fin Set}_0, \alpha: s' \xrightarrow{\sim} t']$$

$$(s', t', \alpha) \mapsto s' + 1$$

$$\downarrow$$

$$\text{Fin Set}_0$$

$$(s', t', \alpha) \mapsto t' + 1$$

$$\downarrow$$

$$\text{Fin Set}_0$$

(which is the span A^*A)

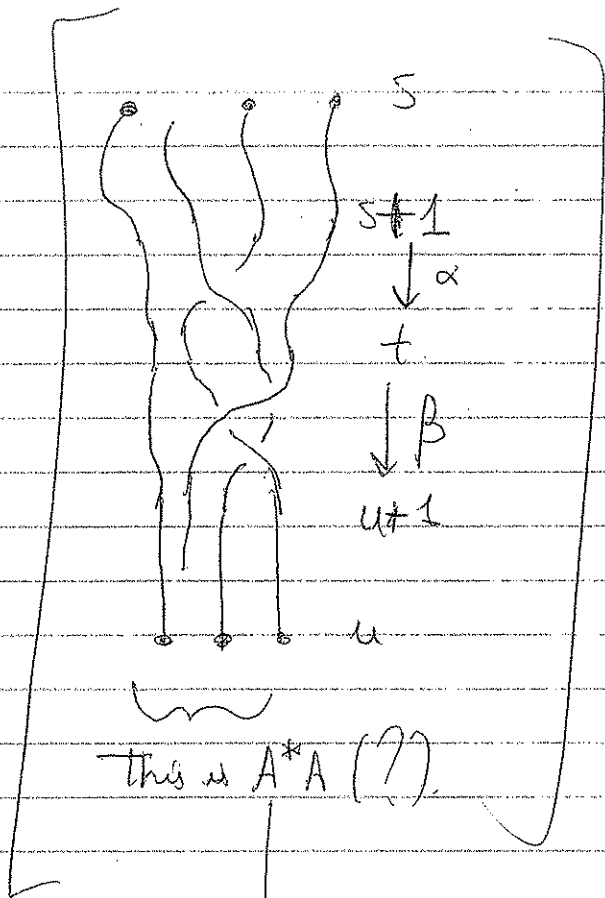
Hence the sum AA^* of these spans is equivalent to

$$1_{\text{Fin Set}_0} + A^*A$$

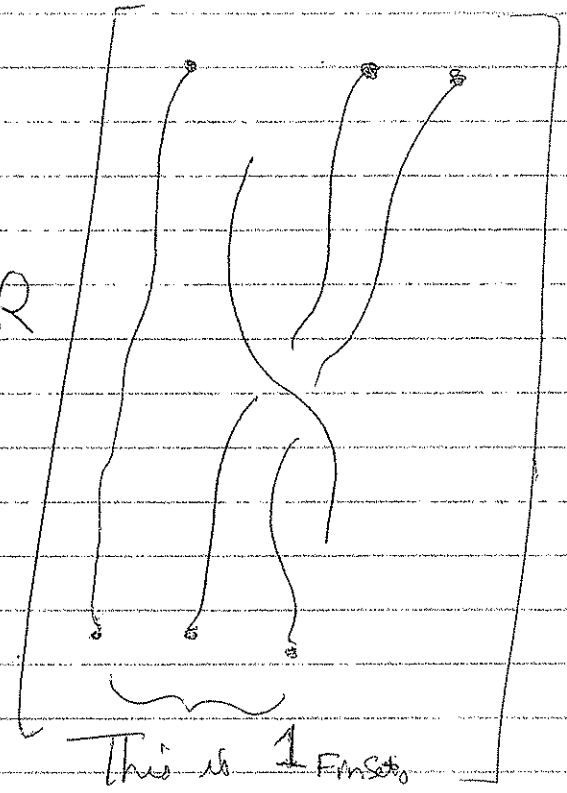
(Proved)

(f) OR VIA Feynman diagrams

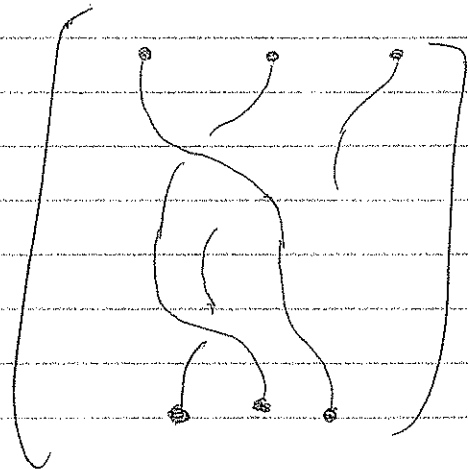
AA^* comes in two cases



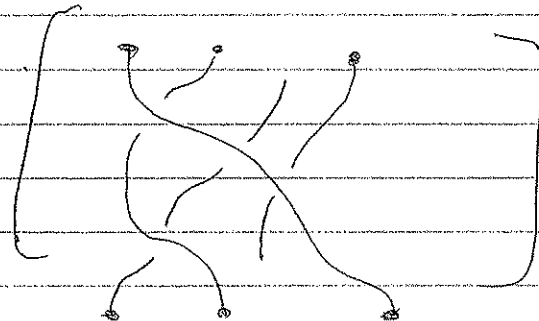
OR



Well, A^*A really is



But one can just make the arrows longer:



→ and this is what we essentially drew at the top of this page.