

3/2008

# John Baez - The Harmonic Oscillator (n deg. of freedom)

1) a) We'll groupoidify representations of  $\mathbb{R}$   $\mathfrak{sl}_n$ , which comes into play once we consider the harmonic oscillator with  $n$ -degrees of freedom.

b) On Thursdays, Jim is talking about the Hall Algebra arising from  $\text{Rep}_{\mathbb{F}_q}(A_2)$ , which generalizes to all  $A_n$  as well.

c) As we'll see/say later,  $q$ -deforming b)  $\Rightarrow$  c) (this has to do with the fact that the Hall Alg. is a "part" =  $U_q(\mathfrak{sl}^+)$  of the  $q$ -deformed  $UEA = U(\mathfrak{gl}_n)$ ).

X

2) So let us first (easily) generalize the theory of the Harmonic Oscillator to  $n$ -variables. So, we have

$$\text{a) } H_0(\text{FinSet}_0^n) \cong H_0(\text{FinSet}_0) \otimes^n \cong k[z] \otimes^n \cong k[z_1, \dots, z_n]$$

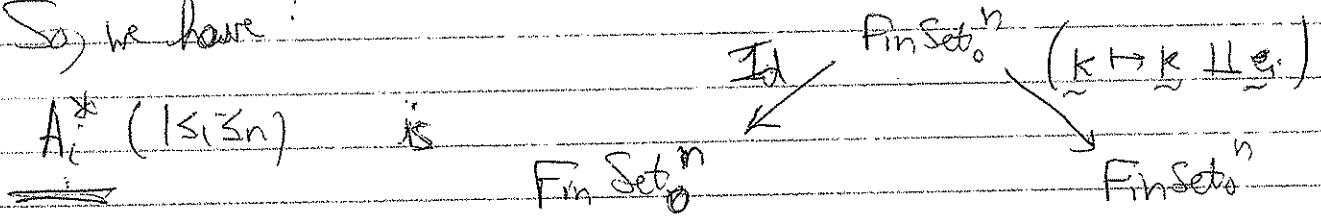
where the isoclass of  $(k_1, \dots, k_n) \in \text{FinSet}_0^n$  corresponds to  $z_1^{k_1} \dots z_n^{k_n} \in k[z_1, \dots, z_n]$ .

We have  $n$  creation and annihilation operators, whose groupoidified versions are spans.

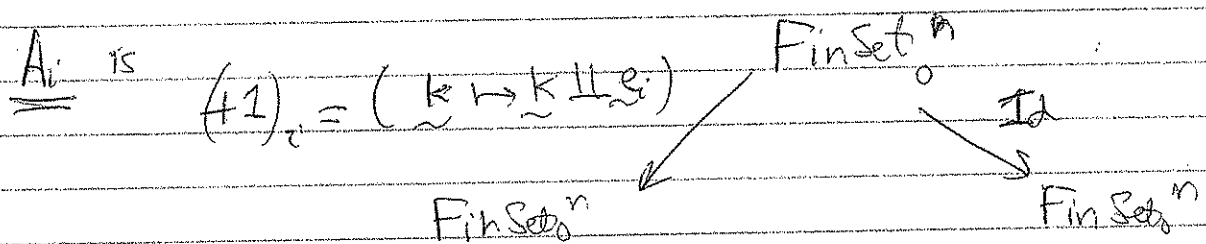
(b) Given  $i$ , denote  $\underline{e}_i = (\emptyset, \emptyset, \rightarrow \{i\} = 1, \emptyset, \dots, \emptyset) \in \text{FinSet}_0^n$

Then the creation operator corresponds to  $[k \mapsto k \cup \underline{e}_i]$   
 i.e.  $(k_1, \dots, k_n) \mapsto (k_1, \dots, k_i + 1, \dots, k_n)$

So, we have:



This "creates a particle of type  $i$ ".



This annihilates a particle of type  $i$ .

(c) These give operator  $\widehat{A}_i = a_i$ ,  $\widehat{A}_i^* = a_i^*$  from the

Fock space  $= k[z_1, \dots, z_n]$

to itself, and explicitly,  $\left[ \begin{array}{l} a_i = \partial / \partial z_i \\ a_i^* = z_i \end{array} \right]$

These satisfy:

$$a_i a_j = a_j a_i$$

$$a_i^* a_j^* = a_j^* a_i^*$$

$$a_i a_j^* = a_j^* a_i + \delta_{ij} \cdot 1$$

These are called the "CCR" = "canonical commutation relns"

So, we hope that it also holds on the groupoid level

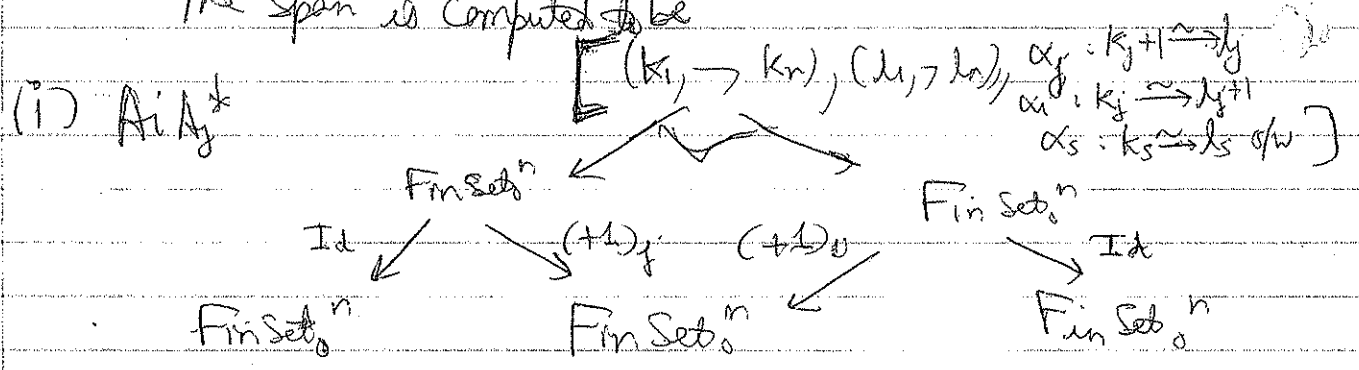
$$A_i A_j = A_j A_i$$

$$A_i^* A_j^* = A_j^* A_i^* \quad (\text{this follows from the previous line})$$

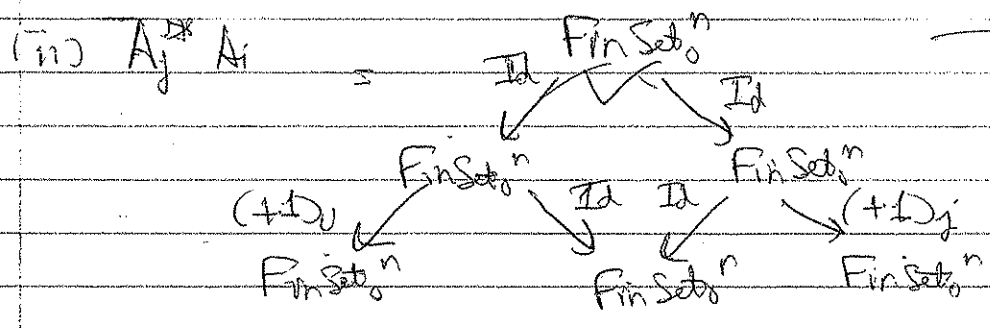
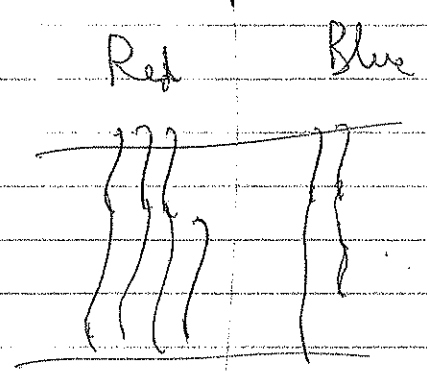
$$A_i A_j^* = A_j^* A_i + \delta_{ij} \cdot 1 \quad (\text{this is already done for } i=j)$$

(d) PE  $\downarrow$  Let's first do the last line only, and for  $i \neq j$

The span is computed to be



In Feynman diagrams, we have



Why are these the same? It's because the ~~fact~~ fact -  
 composite span is  $\cong$

$$\left[ \binom{\min(k_i, l_j)}{i+j} \right]_{1 \leq i, j \leq n}$$

$(+1)_i \swarrow \quad \searrow (+1)_j$

$$= \text{FunSet}_0^n$$

$(+1)_i \swarrow \quad \searrow (+1)_j$

$$\text{FunSet}_0^n \quad \text{FunSet}_0^n$$

and this makes it really  
 easy to see that the desired  
 c.c.r. holds.

② We'll make more remarks on this after the notes for this class.

③ So, we've groupoidified all these things:

① The assoc. alg.  $W_n$  generated by the  $a_i, a_i^* (i=1, n)$   
 This is called the ~~Weyl~~  $n^{\text{th}}$  Weyl algebra.

This has a good-given rep. on  $k[z_1, \dots, z_n]$  with

$$a_i \mapsto \frac{\partial}{\partial z_i}, \quad a_i^* \mapsto m_{z_i}$$

A typical element in this rep. is like  $z_1^3 \frac{\partial}{\partial z_2} + z_2 z_3 \frac{\partial^2}{\partial z_2 \partial z_1}$

So,  $W_n = \{ \text{polynomial coeff. diff. operators} \}$

This becomes a Lie algebra w/  $[a, b] = ab - ba$   
 and then it has lots of Lie subalgebras, including!

b)  $\{ \text{Polynomial coeff vector fields} \} = \text{Diff}^1$   
 $= \mathbb{Q} \cdot k[z_1, \dots, z_n] \cdot \bigoplus k \left( \frac{\partial}{\partial z_i} \right)$

c)  $\{ \text{Homogeneous linear coeff vector fields} \} \cong \mathfrak{gl}(n)$ , i.e.

e.g.  $\begin{pmatrix} 4 & 2 \\ 0 & 3 \end{pmatrix} \mapsto 4z_1 \frac{\partial}{\partial z_1} + 2z_1 \frac{\partial}{\partial z_2} + 3z_2 \frac{\partial}{\partial z_2}$

d) Note that  $[f \cdot \frac{\partial}{\partial z_i}, g \cdot \frac{\partial}{\partial z_j}] = f \cdot \frac{\partial}{\partial z_i} (g \cdot \frac{\partial}{\partial z_j}) - g \cdot \frac{\partial}{\partial z_j} (f \cdot \frac{\partial}{\partial z_i})$   
 $\forall f, g, i, j$ . So this explains why c) works, ~~is~~ and can be generalized.

e) The Heisenberg Lie Algebra  $\mathfrak{h}_n = \text{Span} \{ a_i, a_i^*, 1 \}$   
 This is a  $(2n+1)$ -dim Lie alg.

f)  $\{ \text{Homogeneous quadratic expressions in } a_i, a_i^*, \text{ plus const} \}$   
~~is spanned~~ is spanned by  $a_i^* a_j$   $i < j$   
 $a_i a_j$   $i < j$   
 $a_i^* a_i^*$   $i < j$   
 1

g)  $\boxed{\mathbb{Q}_n}$  Find commutation relations.  
 Is there a codim 1 Lie alg  $\cong \mathfrak{sp}(2n)$ ?

ANS: One can check that (i)  $[a_i^* a_j, a_k^* a_l] = \delta_{jk} a_i^* a_l - \delta_{il} a_k^* a_j$

(ii)  $[a_i a_j, a_k^* a_l] = \delta_{jk} a_i a_l + \delta_{il} a_j a_k$

(iii)  $[a_i^* a_j^*, a_k^* a_l] = -\delta_{il} a_j^* a_k^* - \delta_{jl} a_i^* a_k^*$

And finally,  $[a_i a_j, a_k^* a_l^*] = \delta_{jk} a_i^* a_l + \delta_{il} a_i^* a_j^* + \delta_{jl} a_k^* a_l + \delta_{il} a_k^* a_j + (\delta_{jk} \delta_{il} + \delta_{ik} \delta_{jl})$

This is where scalars may occur!

(h) How to amend? Well, change the basis/vectors from  $a_i^* a_j$  to  $B_{ij} = a_i^* a_j + \frac{1}{2} \delta_{ij} \mathbb{1}$  (so check  $k \neq l$ )

Then (i) above becomes  $[B_{ij}, B_{kl}] = [a_i^* a_j, a_k^* a_l]$  ( $\mathbb{1}$  is central)  
 $= \delta_{jk} a_i^* a_l + \frac{1}{2} \delta_{jk} \delta_{il} - \delta_{il} a_k^* a_j - \frac{1}{2} \delta_{jk} \delta_{il}$

(i)':  $[B_{ij}, B_{kl}] = B_{il} - B_{kj}$

(ii) stays the same, and similar to  $\downarrow$ , (iii) changes to:

(iii)':  $[a_i a_j, a_k^* a_l^*] = \delta_{jk} B_{il} + \delta_{il} B_{lj} + \delta_{jl} B_{kl} + \delta_{il} B_{kj}$   
 and these span the symplectic Lie algebra. (Proved)

④ (Continuing ②③ above)

① Firstly, an ERRATUM: The weak pullback over  $\mathcal{D}^1 X, Y, Z$  is the terminal object, not the initial one, in the category of objects over  $X, Y, Z$ .

② Now, here is a slightly general criterion that can be used to prove ~~the~~ some of the CER's (i.e. all the CER's with commuting spans).

**Defn** A functor is faithful and reflects isomorphisms if  $(F: \mathcal{C} \rightarrow \mathcal{D})$

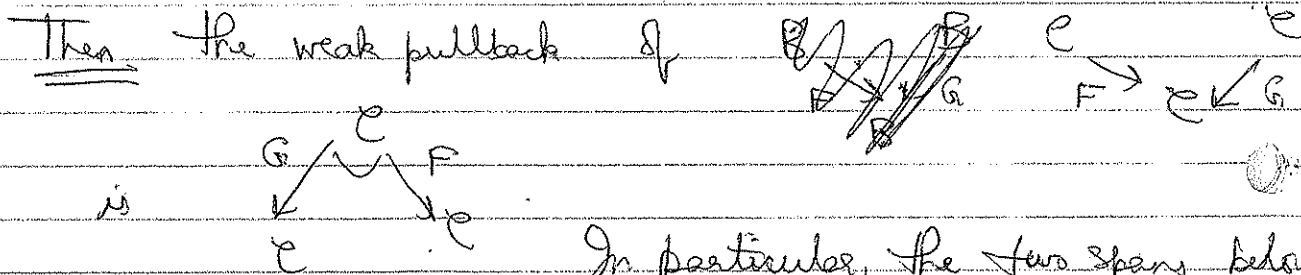
if ①  $F(\text{Hom}_{\mathcal{C}}(X, Y)) \rightarrow \text{Hom}_{\mathcal{D}}(F(X), F(Y))$  is injective for all objects  $X, Y$  in  $\mathcal{C}$

and ② whenever  $F(X) \cong F(Y)$  in  $\mathcal{D}$ , we must have  $X \cong Y$  in  $\mathcal{C}$

③ **Lemna** Say  $\mathcal{C}$  is a category with functors  $F, G: \mathcal{C} \rightarrow \mathcal{C}$  that ~~are both isomorphisms~~

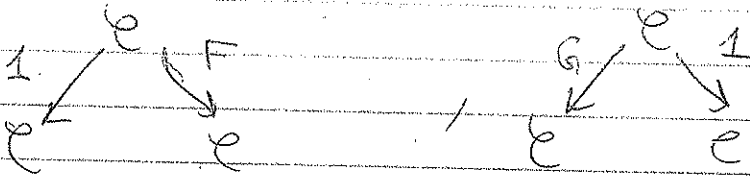
(i) are faithful, (ii) reflect isomorphisms, and (iii) commute.

Also say that whenever  $F(x) \cong G(y)$  in  $\mathcal{C}$ , we must have  $x \cong G(z)$ ,  $y \cong F(z')$  for some  $z, z'$  in  $\mathcal{C}$ .



In particular, the two spans below commute

(and this implies most of the CER's above).



(d) **Proof** We first show that the weak pullback is, indeed,  $\mathcal{C}$ .

The weak pullback is  $\{(x, y) \in \text{Ob}(\mathcal{C}) \times \text{Ob}(\mathcal{C}) ; \alpha: F(x) \xrightarrow{\sim} G(y)\}$

By the given condn, this gives  $(z, z') \in \mathcal{C}$  with

$$G(z) \cong x, F(z') \cong y.$$

But then  $F(x) \cong F(G(z)) = FG(z)$

$$\alpha \downarrow$$

$$G(y) \cong G(F(z')) = GF(z') \cong FG(z')$$

↓  
we can weaken this!

ie  $FG(z) \cong FG(z')$ . Since  $F, G$  both reflect isms,

$$z \cong z' \text{ in } \mathcal{C}, \text{ ie } (x, y) \cong \text{~~(x, y)~~$$

**CHECK** This can be used to construct an equivalence between  $\mathcal{C}$  and the weak pullback. (By universality, the functor in the other direction exists:  $z \mapsto (G, F)(z)$ )

**CHECK 2** The above spans commute (this should be similar to what was done in class - and hence fairly straightforward!).