

Solutions to suggested homework problems from
An Introduction to Partial Differential Equations by Yehuda Pinchover and Jacob Rubinstein

Suggested problems: Exercises 8.5, 8.6, 8.7, 8.8, 8.11

8.5. (a) Show that the function

$$\begin{aligned} G(x, y; \xi, \eta) &:= \Gamma(x - \xi, y - \eta) - \Gamma(x - \xi, y - \tilde{\eta}) \\ &= -\frac{1}{2\pi} \ln \left(\frac{\sqrt{(x - \xi)^2 + (y - \eta)^2}}{\sqrt{(x - \xi)^2 + (y + \eta)^2}} \right) \end{aligned} \quad (8.23)$$

is indeed the Green function in \mathbb{R}_+^2 , and that its derivative in the y direction for $y = 0$ is the Poisson kernel which is given by

$$K(x, 0; \xi, \eta) := \frac{\eta}{\pi((x - \xi)^2 + \eta^2)} \quad (8.24)$$

for all $(x, 0) \in \partial\mathbb{R}_+^2$ and $(\xi, \eta) \in \mathbb{R}_+^2$.

Solution. Recall that the fundamental solution of the Laplace equation with a pole at (ξ, η) is given by

$$\begin{aligned} \Gamma(x - \xi, y - \eta) &:= -\frac{1}{2\pi} \ln(\sqrt{(x - \xi)^2 + (y - \eta)^2}) \\ &= -\frac{1}{4\pi} \ln((x - \xi)^2 + (y - \eta)^2) \end{aligned}$$

and satisfies

$$\Delta \Gamma = -\delta(x - \xi, y - \eta)$$

for all $(x, y) \in \mathbb{R}_+^2$ with $(x, y) \neq (\xi, \eta)$, where

$$\delta(x - \xi, y - \eta) := \begin{cases} \infty & \text{if } (x, y) = (\xi, \eta), \\ 0 & \text{if } (x, y) \neq (\xi, \eta). \end{cases}$$

Given the upper half plane \mathbb{R}_+^2 , the inverse point of (x, y) with respect to the real line is

$$(\tilde{x}, \tilde{y}) := (x, -y).$$

By combining (8.13) with the boundary condition of (8.11) from the textbook, the possible expressions of the given function are

$$\begin{aligned} G(x, y; \xi, \eta) &:= \Gamma(x - \xi, y - \eta) - \Gamma(x - \xi, y - \tilde{\eta}) \\ &= \Gamma(x - \xi, y - \eta) - \Gamma(x - \xi, y + \eta) \\ &= -\frac{1}{2\pi} \ln(\sqrt{(x - \xi)^2 + (y - \eta)^2}) + \frac{1}{2\pi} \ln(\sqrt{(x - \xi)^2 + (y + \eta)^2}) \\ &= -\frac{1}{2\pi} \ln \left(\frac{\sqrt{(x - \xi)^2 + (y - \eta)^2}}{\sqrt{(x - \xi)^2 + (y + \eta)^2}} \right). \end{aligned} \quad (8.23)$$

So we have

$$\begin{aligned} \Delta G(x, y; \xi, \eta) &= \Delta(\Gamma(x - \xi, y - \eta) - \Gamma(x - \xi, y + \eta)) \\ &= \Delta \Gamma(x - \xi, y - \eta) - \Delta \Gamma(x - \xi, y + \eta) \\ &= -\delta(x - \xi, y - \eta) - 0 \\ &= -\delta(x - \xi, y - \eta) \end{aligned}$$

and

$$\begin{aligned} G(x, 0; \xi, \eta) &= -\frac{1}{2\pi} \ln \left(\frac{\sqrt{(x - \xi)^2 + (0 - \eta)^2}}{\sqrt{(x - \xi)^2 + (0 + \eta)^2}} \right) \\ &= -\frac{1}{2\pi} \ln \left(\frac{\sqrt{(x - \xi)^2 + \eta^2}}{\sqrt{(x - \xi)^2 + \eta^2}} \right) \\ &= -\frac{1}{2\pi} \ln(1) \\ &= -\frac{1}{2\pi} 0 \\ &= 0. \end{aligned}$$

In summary, $G(x, y; \xi, \eta)$ solves

$$\begin{aligned}\Delta G &= -\delta(x - \xi, y - \eta) & (x, y) \in \mathbb{R}_+^2, \\ G(x, 0; \xi, \eta) &= 0 & (x, y) \in \partial\mathbb{R}_+^2,\end{aligned}\tag{8.14}$$

meaning that G is indeed the Green function. And the derivative of G in the y direction is

$$\begin{aligned}G_y(x, y; \xi, \eta) &= \frac{\partial}{\partial y} \left(-\frac{1}{2\pi} (\ln(\sqrt{(x - \xi)^2 + (y - \eta)^2}) - \ln(\sqrt{(x - \xi)^2 + (y + \eta)^2})) \right) \\ &= -\frac{1}{2\pi} \frac{\partial}{\partial y} (\ln(\sqrt{(x - \xi)^2 + (y - \eta)^2})) + \frac{1}{2\pi} \frac{\partial}{\partial y} (\ln(\sqrt{(x - \xi)^2 + (y + \eta)^2})) \\ &= -\frac{1}{2\pi} \frac{y - \eta}{(x - \xi)^2 + (y - \eta)^2} + \frac{1}{2\pi} \frac{y + \eta}{(x - \xi)^2 + (y + \eta)^2}.\end{aligned}$$

At $y = 0$, we obtain

$$\begin{aligned}G_y(x, 0; \xi, \eta) &= -\frac{1}{2\pi} \frac{0 - \eta}{(x - \xi)^2 + (0 - \eta)^2} + \frac{1}{2\pi} \frac{0 + \eta}{(x - \xi)^2 + (0 + \eta)^2} \\ &= \frac{1}{2\pi} \frac{\eta}{((x - \xi)^2 + \eta^2)} + \frac{1}{2\pi} \frac{\eta}{((x - \xi)^2 + \eta^2)} \\ &= \frac{\eta}{\pi((x - \xi)^2 + \eta^2)} \\ &= K(x, 0; \xi, \eta)\end{aligned}$$

for all $(x, 0) \in \partial\mathbb{R}_+^2$ and $(\xi, \eta) \in \mathbb{R}_+^2$, as desired. \square

- (b) Using a reflection principle and part (a), find the Green function of the positive quarter plane $\{(x, y) \in \mathbb{R}^2 \mid x > 0, y > 0\}$.

Solution. Given the positive quarter plane $D := \{(x, y) \in \mathbb{R}^2 \mid x > 0, y > 0\}$, the inverse points of (x, y) with respect to ∂D are

$$\begin{aligned}(x_1, y_1) &:= (x, -y), \\ (x_2, y_2) &:= (-x, y), \\ (x_3, y_3) &:= (-x, -y).\end{aligned}$$

Using the reflection principle, we have the function

$$\begin{aligned}G(x, y; \xi, \eta) &= \Gamma(x - \xi, y - \eta) - \Gamma(x - \xi_1, y - \eta_1) - \Gamma(x - \xi_2, y - \eta_2) + \Gamma(x - \xi_3, y - \eta_3) \\ &= \Gamma(x - \xi, y - \eta) - \Gamma(x - \xi, y + \eta) - \Gamma(x + \xi, y - \eta) + \Gamma(x + \xi, y + \eta) \\ &= -\frac{1}{2\pi} \ln(\sqrt{(x - \xi)^2 + (y - \eta)^2}) + \frac{1}{2\pi} \ln(\sqrt{(x - \xi)^2 + (y + \eta)^2}) \\ &\quad + \frac{1}{2\pi} \ln(\sqrt{(x + \xi)^2 + (y - \eta)^2}) - \frac{1}{2\pi} \ln(\sqrt{(x + \xi)^2 + (y + \eta)^2}) \\ &= -\frac{1}{4\pi} \ln((x - \xi)^2 + (y - \eta)^2) + \frac{1}{4\pi} \ln((x - \xi)^2 + (y + \eta)^2) \\ &\quad + \frac{1}{4\pi} \ln((x + \xi)^2 + (y - \eta)^2) - \frac{1}{4\pi} \ln((x + \xi)^2 + (y + \eta)^2) \\ &= \boxed{-\frac{1}{4\pi} \ln \left(\frac{((x - \xi)^2 + (y - \eta)^2)((x + \xi)^2 + (y + \eta)^2)}{((x - \xi)^2 + (y + \eta)^2)((x + \xi)^2 + (y - \eta)^2)} \right)}.\end{aligned}$$

Now, it remains to show that this function is indeed the Green function. We have

$$\begin{aligned}\Delta G(x, y; \xi, \eta) &= \Delta \Gamma(x - \xi, y - \eta) - \Delta \Gamma(x - \xi_1, y - \eta_1) - \Delta \Gamma(x - \xi_2, y - \eta_2) + \Delta \Gamma(x - \xi_3, y - \eta_3) \\ &= -\delta(x - \xi, y - \eta) - 0 - 0 + 0 \\ &= -\delta(x - \xi, y - \eta).\end{aligned}$$

We also have, for all $x \geq 0$,

$$\begin{aligned}G(x, 0; \xi, \eta) &= -\frac{1}{4\pi} \ln \left(\frac{((x - \xi)^2 + (0 - \eta)^2)((x + \xi)^2 + (0 + \eta)^2)}{((x - \xi)^2 + (0 + \eta)^2)((x + \xi)^2 + (0 - \eta)^2)} \right) \\ &= -\frac{1}{4\pi} \ln \left(\frac{((x - \xi)^2 + \eta^2)((x + \xi)^2 + \eta^2)}{((x - \xi)^2 + \eta^2)((x + \xi)^2 + \eta^2)} \right) \\ &= -\frac{1}{4\pi} \ln(1) \\ &= -\frac{1}{4\pi} 0 \\ &= 0\end{aligned}$$

and, for all $y \geq 0$,

$$\begin{aligned}
G(0, y; \xi, \eta) &= -\frac{1}{4\pi} \ln \left(\frac{((0 - \xi)^2 + (y - \eta)^2)((0 + \xi)^2 + (y + \eta)^2)}{((0 - \xi)^2 + (y + \eta)^2)((0 + \xi)^2 + (y - \eta)^2)} \right) \\
&= -\frac{1}{4\pi} \ln \left(\frac{(\xi^2 + (y - \eta)^2)(\xi^2 + (y + \eta)^2)}{(\xi^2 + (y + \eta)^2)(\xi^2 + (y - \eta)^2)} \right) \\
&= -\frac{1}{4\pi} \ln(1) \\
&= -\frac{1}{4\pi} 0 \\
&= 0.
\end{aligned}$$

In summary, $G(x, y; \xi, \eta)$ solves

$$\begin{aligned}
\Delta G(x, y; \xi, \eta) &= -\delta(x - \xi, y - \eta) & (x, y) \in D, \\
G(x, y; \xi, \eta) &= 0 & (x, y) \in \partial D,
\end{aligned} \tag{8.14}$$

meaning that G is indeed the Green function. □

8.6. Let \mathbb{R}_+^2 be the upper half-plane. Find the Neumann function of \mathbb{R}_+^2 .

Solution. Given the upper half-plane \mathbb{R}_+^2 , the inverse point of (x, y) with respect to the real line is

$$(\tilde{x}, \tilde{y}) := (x, -y).$$

The possible expressions of the given function are

$$\begin{aligned}
N(x, y; \xi, \eta) &:= \Gamma(x - \xi, y - \eta) + \Gamma(x - \tilde{\xi}, y - \tilde{\eta}) + C \\
&= \Gamma(x - \xi, y - \eta) + \Gamma(x - \xi, y + \eta) + C \\
&= -\frac{1}{2\pi} \ln(\sqrt{(x - \xi)^2 + (y - \eta)^2}) - \frac{1}{2\pi} \ln(\sqrt{(x - \xi)^2 + (y + \eta)^2}) + C \\
&= -\frac{1}{2\pi} \ln(\sqrt{(x - \xi)^2 + (y - \eta)^2} \sqrt{(x - \xi)^2 + (y + \eta)^2}) + C,
\end{aligned}$$

where C is a constant. This expression is based on the fact that one must place another positive charge (instead of negative charge for Dirichlet case) on the image point $(\tilde{\xi}, \tilde{\eta}) = (\xi, -\eta)$, in order to satisfy the Neumann condition $\partial_n N(x, y; \xi, \eta) = 0 = -\frac{1}{L}$ with $L = \infty$. We have

$$\begin{aligned}
\Delta N(x, y; \xi, \eta) &= \Delta(\Gamma(x - \xi, y - \eta) - \Gamma(x - \xi, y + \eta) + C) \\
&= \Delta\Gamma(x - \xi, y - \eta) - \Delta\Gamma(x - \xi, y + \eta) - \Delta C \\
&= -\delta(x - \xi, y - \eta) - 0 + 0 \\
&= -\delta(x - \xi, y - \eta)
\end{aligned}$$

and, for all $(x, 0; \xi, \eta) \in \partial\mathbb{R}_+^2$,

$$\begin{aligned}
\partial_n(x, 0; \xi, \eta) &= \partial_n \left(-\frac{1}{2\pi} \ln(\sqrt{(x - \xi)^2 + (0 - \eta)^2} \sqrt{(x - \xi)^2 + (0 + \eta)^2}) + C \right) \\
&= -\frac{1}{2\pi} \partial_n \ln(\sqrt{(x - \xi)^2 + \eta^2} \sqrt{(x - \xi)^2 + \eta^2}) - \partial_n C \\
&= -\frac{1}{2\pi} \partial_n \ln((x - \xi)^2 + \eta^2) \\
&= -\frac{1}{2\pi} \nabla(\ln((x - \xi)^2 + \eta^2)) \cdot \hat{n} \\
&= -\frac{1}{2\pi} \left(\frac{\partial}{\partial x}(\ln((x - \xi)^2 + \eta^2)), \frac{\partial}{\partial y}(\ln((x - \xi)^2 + \eta^2)) \right) \cdot (0, -1) \\
&= \frac{1}{2\pi} \frac{\partial}{\partial y} \ln((x - \xi)^2 + \eta^2) \\
&= \frac{1}{2\pi} 0 \\
&= 0 \\
&= -\frac{1}{L}
\end{aligned}$$

with $L = \infty$. In summary, $N(x, y; \xi, \eta)$ solves

$$\begin{aligned}\Delta N(x, y; \xi, \eta) &= -\delta(x - \xi, y - \eta) \quad (x, y) \in \mathbb{R}_+^2, \\ \partial_n N(x, 0; \xi, \eta) &= -\frac{1}{L} \quad (x, y) \in \partial\mathbb{R}_+^2,\end{aligned}\tag{8.29}$$

meaning that G is indeed the Neumann function. □

8.7. (a) Let u be a smooth function with a compact support in \mathbb{R}^2 . Prove

$$\lim_{\epsilon \rightarrow 0^+} \int_{\mathbb{R}^2} \rho_\epsilon(\vec{x}) u(\vec{x}) d\vec{x} = u(\vec{y}) = \int_{\mathbb{R}^2} \delta(\vec{x} - \vec{y}) u(\vec{x}) d\vec{x}.\tag{8.9}$$

Proof. Recall from page 212 of the textbook that $\rho_\epsilon(\vec{x})$ has compact support in $B(0, \epsilon)$ and satisfies

$$\int_{\mathbb{R}^2} \rho_\epsilon(\vec{x}) d\vec{x} = \int_{B(0, \epsilon)} \rho_\epsilon(\vec{x}) d\vec{x} = 1,$$

which implies

$$\begin{aligned}\int_{\mathbb{R}^2} \rho_\epsilon(\vec{x}) u(\vec{x}) d\vec{x} - u(\vec{y}) &= \int_{\mathbb{R}^2} \rho_\epsilon(\vec{x}) u(\vec{x}) d\vec{x} - u(\vec{y}) \int_{\mathbb{R}^2} \rho_\epsilon(\vec{x}) d\vec{x} \\ &= \int_{\mathbb{R}^2} \rho_\epsilon(\vec{x}) u(\vec{x}) d\vec{x} - \int_{\mathbb{R}^2} \rho_\epsilon(\vec{x}) u(\vec{y}) d\vec{x} \\ &= \int_{\mathbb{R}^2} \rho_\epsilon(\vec{x}) u(\vec{x}) - \rho_\epsilon(\vec{x}) u(\vec{y}) d\vec{x} \\ &= \int_{\mathbb{R}^2} \rho_\epsilon(\vec{x}) (u(\vec{x}) - u(\vec{y})) d\vec{x}.\end{aligned}$$

Since u is smooth in \mathbb{R}^2 , it is continuous at $y \in \mathbb{R}^2$. So we can invoke the ϵ - δ definition of continuity, which states: For all $\epsilon > 0$, there exists $\delta > 0$ that satisfies $|u(x) - u(y)| < \epsilon$ for all $y \in B_\delta(x)$. Applying this definition and using the triangle inequality for integrals, we have

$$\begin{aligned}\left| \int_{\mathbb{R}^2} \rho_\epsilon(\vec{x}) u(\vec{x}) d\vec{x} - u(\vec{y}) \right| &= \left| \int_{\mathbb{R}^2} \rho_\epsilon(\vec{x}) (u(\vec{x}) - u(\vec{y})) d\vec{x} \right| \\ &\leq \int_{\mathbb{R}^2} |\rho_\epsilon(\vec{x}) u(\vec{x}) - \rho_\epsilon(\vec{x}) u(\vec{y})| d\vec{x} \\ &= \int_{\mathbb{R}^2} \rho_\epsilon(\vec{x}) |u(\vec{x}) - u(\vec{y})| d\vec{x} \\ &< \int_{\mathbb{R}^2} \rho_\epsilon(\vec{x}) \epsilon d\vec{x} \\ &< \epsilon \int_{\mathbb{R}^2} \rho_\epsilon(\vec{x}) d\vec{x} \\ &= \epsilon \cdot 1 \\ &= \epsilon \\ &\rightarrow 0\end{aligned}$$

as $\epsilon \rightarrow 0^+$. This implies

$$\int_{\mathbb{R}^2} \rho_\epsilon(\vec{x}) u(\vec{x}) d\vec{x} \rightarrow u(y)$$

as $\epsilon \rightarrow 0^+$, as desired. □

(b) Find the constant c in

$$\rho(\vec{x}) := \begin{cases} c \exp(\frac{1}{|\vec{x}|^2 - 1}) & \text{if } |\vec{x}| \leq 1, \\ 0 & \text{if } |\vec{x}| > 1, \end{cases}\tag{8.10}$$

and verify directly that ρ_ϵ is an approximation of the Dirac delta function.

Solution. We need to find the constant c that satisfies

$$\int_{\mathbb{R}^2} \rho(\vec{x}) d\vec{x} = 1.$$

We have

$$\begin{aligned}
 1 &= \int_{\mathbb{R}^2} \rho(\vec{x}) d\vec{x} = c \int_{B(0,1)} \exp\left(\frac{1}{|\vec{x}|^2 - 1}\right) d\vec{x} \\
 &= c \int_0^{2\pi} \int_0^1 \exp\left(\frac{1}{r^2 - 1}\right) r dr d\theta \\
 &= 2\pi c \int_0^1 \exp\left(\frac{1}{r^2 - 1}\right) r dr \\
 &\approx 0.466512c,
 \end{aligned}$$

where the approximate value is taken from [this computation on WolframAlpha](#). We find $c \approx \frac{1}{0.466512} = \boxed{2.143566}$. \square

Remark. Although the integral

$$\int_0^1 \exp\left(\frac{1}{r^2 - 1}\right) r dr$$

is convergent, we are unable to evaluate it analytically. We can try to invoke the Taylor series of the exponential function to obtain

$$\int_0^1 \exp\left(\frac{1}{r^2 - 1}\right) r dr = \sum_{n=0}^{\infty} \frac{1}{n!} \int_0^1 \frac{r}{(r^2 - 1)^n} dr,$$

but the Taylor series entails an antiderivative that contains a divergent term. Our only recourse is to resort to numerical methods or advanced calculators in order to approximate the integral, which means we can only approximate the value of c at best.

8.8. Let $k \neq 0$. Show that the function $G_k(x; \xi) = \frac{1}{2k} e^{-k|x-\xi|}$ is a fundamental solution of the equation

$$-u'' + k^2 u = 0$$

for all $-\infty < x < \infty$.

Hint: Use one of Green's identities.

Proof. Given the function

$$\begin{aligned}
 G_k(x; \xi) &:= \frac{1}{2k} e^{-k|x-\xi|} \\
 &= \begin{cases} \frac{1}{2k} e^{-k(x-\xi)} & \text{if } x \geq \xi, \\ \frac{1}{2k} e^{k(x-\xi)} & \text{if } x < \xi, \end{cases}
 \end{aligned}$$

we obtain its first and second derivatives

$$\begin{aligned}
 G'_k(x; \xi) &= \begin{cases} -\frac{1}{2} e^{-k(x-\xi)} & \text{if } x > \xi, \\ \frac{1}{2} e^{k(x-\xi)} & \text{if } x < \xi, \end{cases} \\
 G''_k(x; \xi) &= \begin{cases} \frac{k}{2} e^{-k(x-\xi)} & \text{if } x > \xi, \\ \frac{k}{2} e^{k(x-\xi)} & \text{if } x < \xi. \end{cases}
 \end{aligned}$$

For all $x \neq \xi$, we have

$$\begin{aligned}
 -G''_k(x; \xi) + k^2 G_k(x; \xi) &= -\begin{cases} \frac{k}{2} e^{-k(x-\xi)} & \text{if } x > \xi, \\ \frac{k}{2} e^{k(x-\xi)} & \text{if } x < \xi. \end{cases} + k^2 \begin{cases} \frac{1}{2k} e^{-k(x-\xi)} & \text{if } x > \xi, \\ \frac{1}{2k} e^{k(x-\xi)} & \text{if } x < \xi, \end{cases} \\
 &= \begin{cases} -\frac{k}{2} e^{-k(x-\xi)} + \frac{k}{2} e^{-k(x-\xi)} & \text{if } x > \xi, \\ -\frac{k}{2} e^{k(x-\xi)} + \frac{k}{2} e^{k(x-\xi)} & \text{if } x < \xi, \end{cases} \\
 &= \begin{cases} 0 & \text{if } x > \xi, \\ 0 & \text{if } x < \xi, \end{cases} \\
 &= 0,
 \end{aligned}$$

which implies in particular $-G''_k(x; \xi) + k^2 G_k(x; \xi) = 0$ for all $x \in (-\infty, x - \epsilon) \cup (x + \epsilon, \infty)$, where $\epsilon > 0$ is arbitrary. By

using Green's third identity (integration by parts), we have

$$\begin{aligned}
\int_{-\infty}^{\infty} u(x)(-G_k''(x; \xi) + k^2 G_k(x; \xi)) dx &= \int_{\xi-\epsilon}^{\xi+\epsilon} u(x)(-G_k''(x; \xi) + k^2 G_k(x; \xi)) dx \\
&\quad + \int_{(-\infty, \xi-\epsilon) \cup (\xi+\epsilon, \infty)} u(x)(-G_k''(x; \xi) + k^2 G_k(x; \xi)) dx \\
&= \int_{\xi-\epsilon}^{\xi+\epsilon} u(x)(-G_k''(x; \xi) + k^2 G_k(x; \xi)) dx + \int_{(-\infty, \xi-\epsilon) \cup (\xi+\epsilon, \infty)} u(x) \cdot 0 dx \\
&= \int_{\xi-\epsilon}^{\xi+\epsilon} u(x)(-G_k''(x; \xi) + k^2 G_k(x; \xi)) dx \\
&= - \int_{\xi-\epsilon}^{\xi+\epsilon} u(x) G_k''(x; \xi) dx + \int_{\xi-\epsilon}^{\xi+\epsilon} k^2 u(x) G_k(x; \xi) dx \\
&= - \left(u(x) G_k'(x; \xi) \Big|_{\xi-\epsilon}^{\xi+\epsilon} + \int_{\xi-\epsilon}^{\xi+\epsilon} u'(x) G_k'(x; \xi) dx \right) \\
&\quad + \int_{\xi-\epsilon}^{\xi+\epsilon} k^2 u(x) G_k(x; \xi) dx \\
&= -(u(\xi + \epsilon) G_k'(\xi + \epsilon; \xi) - u(\xi - \epsilon) G_k'(\xi - \epsilon; \xi)) \\
&\quad + \int_{\xi-\epsilon}^{\xi+\epsilon} -u'(x) G_k'(x; \xi) + k^2 u(x) G_k(x; \xi) dx \\
&=: A_\epsilon + B_\epsilon.
\end{aligned}$$

We have

$$\begin{aligned}
A_\epsilon &= -(u(\xi + \epsilon) G_k'(\xi + \epsilon; \xi) - u(\xi - \epsilon) G_k'(\xi - \epsilon; \xi)) \\
&= -u(\xi + \epsilon) G_k'(\xi + \epsilon; \xi) + u(\xi - \epsilon) G_k'(\xi - \epsilon; \xi) \\
&= -u(\xi + \epsilon) \left(-\frac{1}{2} e^{-k((\xi+\epsilon)-\xi)} \right) + u(\xi - \epsilon) \left(\frac{1}{2} e^{k((\xi-\epsilon)-\xi)} \right) \\
&= \frac{1}{2} u(\xi + \epsilon) e^{-k\epsilon} + \frac{1}{2} u(\xi - \epsilon) e^{-k\epsilon} \\
&\rightarrow \frac{1}{2} u(\xi) + \frac{1}{2} u(\xi) \\
&= u(\xi)
\end{aligned}$$

as $\epsilon \rightarrow 0^+$. As $u(x)$ and $G_k(x; \xi)$ and their first derivatives are bounded, we have

$$\begin{aligned}
|u(x)| &\leq C_1, \\
|u'(x)| &\leq C_2, \\
|G_k(x; \xi)| &\leq D_1, \\
|G_k'(x; \xi)| &\leq D_2,
\end{aligned}$$

where C_1, C_2, D_1, D_2 are constants, and so, by the triangle inequality and the triangle inequality for integrals, we have

$$\begin{aligned}
|B_\epsilon| &= \left| \int_{\xi-\epsilon}^{\xi+\epsilon} -u'(x) G_k'(x; \xi) + k^2 u(x) G_k(x; \xi) dx \right| \\
&\leq \int_{\xi-\epsilon}^{\xi+\epsilon} |-u'(x) G_k'(x; \xi) + k^2 u(x) G_k(x; \xi)| dx \\
&\leq \int_{\xi-\epsilon}^{\xi+\epsilon} |-u'(x) G_k'(x; \xi)| + |k^2 u(x) G_k(x; \xi)| dx \\
&= \int_{\xi-\epsilon}^{\xi+\epsilon} |u'(x)| |G_k'(x; \xi)| + k^2 |u(x) G_k(x; \xi)| dx \\
&\leq \int_{\xi-\epsilon}^{\xi+\epsilon} C_2 D_2 + k^2 C_1 D_1 dx \\
&= 2(C_2 D_2 + k^2 C_1 D_1) \epsilon \\
&\rightarrow 0,
\end{aligned}$$

which implies $B_\epsilon \rightarrow 0$, as $\epsilon \rightarrow 0^+$. Therefore, we conclude

$$\begin{aligned} \int_{-\infty}^{\infty} u(x)(-G_k''(x; \xi) + k^2 G_k(x; \xi)) dx &= \lim_{\epsilon \rightarrow 0^+} \int_{-\infty}^{\infty} u(x)(-G_k''(x; \xi) + k^2 G_k(x; \xi)) dx \\ &= \lim_{\epsilon \rightarrow 0^+} (A_\epsilon + B_\epsilon) \\ &= \lim_{\epsilon \rightarrow 0^+} A_\epsilon + \lim_{\epsilon \rightarrow 0^+} B_\epsilon \\ &= u(\xi) + 0 \\ &= u(\xi), \end{aligned}$$

and so, according to Definition 8.3(b) in pages 212-213 of the textbook, $G_k(x; \xi)$ satisfies

$$-G_k''(x; \xi) + k^2 G_k(x; \xi) = \delta(x - \xi)$$

for all $-\infty < x < \infty$, which means that $G_k(x; \xi)$ is a fundamental solution of the equation $-u'' + k^2 u = 0$. \square

8.11. Let $D_R := \mathbb{R}^2 \setminus B_R$ be the exterior of the disk with radius R centered at the origin. Find the (Dirichlet) Green function of D_R .

Remark. For my solution below, I am following Example 8.14 of the textbook with necessary modifications for D_R .

Solution. Let $(x, y) \in D_R$. The point

$$(\tilde{x}, \tilde{y}) := \frac{R^2}{x^2 + y^2} (x, y) = \left(\frac{R^2}{x^2 + y^2} x, \frac{R^2}{x^2 + y^2} y \right)$$

is the inverse point of (x, y) with respect to the circle ∂D_R . Define the function

$$G_R(x, y; \xi, \eta) := \Gamma(x - \xi, y - \eta) - \Gamma\left(\frac{\sqrt{\xi^2 + \eta^2}}{R} (x - \tilde{\xi}, y - \tilde{\eta})\right)$$

and set

$$\begin{aligned} r &= \sqrt{(x - \xi)^2 + (y - \eta)^2}, \\ r^* &= \sqrt{\left(x - \frac{R^2}{\rho^2} \xi\right)^2 + \left(y - \frac{R^2}{\rho^2} \eta\right)^2}, \\ \rho &= \sqrt{\xi^2 + \eta^2}. \end{aligned}$$

Then we have

$$\begin{aligned} G_R(x, y; \xi, \eta) &= \Gamma(x - \xi, y - \eta) - \Gamma\left(\frac{\sqrt{\xi^2 + \eta^2}}{R} (x - \tilde{\xi}, y - \tilde{\eta})\right) \\ &= \Gamma(x - \xi, y - \eta) - \Gamma\left(\frac{\sqrt{\xi^2 + \eta^2}}{R} \left(x - \frac{R^2}{\xi^2 + \eta^2} \xi\right), \frac{\sqrt{\xi^2 + \eta^2}}{R} \left(y - \frac{R^2}{\xi^2 + \eta^2} \eta\right)\right) \\ &= \Gamma(x - \xi, y - \eta) - \Gamma\left(\frac{\rho}{R} \left(x - \frac{R^2}{\rho^2} \xi\right), \frac{\rho}{R} \left(y - \frac{R^2}{\rho^2} \eta\right)\right) \\ &= -\frac{1}{2\pi} \ln(\sqrt{(x - \xi)^2 + (y - \eta)^2}) + \frac{1}{2\pi} \ln\left(\sqrt{\frac{\rho^2}{R^2} \left(x - \frac{R^2}{\rho^2} \xi\right)^2 + \frac{\rho^2}{R^2} \left(y - \frac{R^2}{\rho^2} \eta\right)^2}\right) \\ &= -\frac{1}{2\pi} \ln(\sqrt{(x - \xi)^2 + (y - \eta)^2}) + \frac{1}{2\pi} \ln\left(\frac{\rho}{R} \sqrt{\left(x - \frac{R^2}{\rho^2} \xi\right)^2 + \left(y - \frac{R^2}{\rho^2} \eta\right)^2}\right) \\ &= -\frac{1}{2\pi} \ln(r) + \frac{1}{2\pi} \ln\left(\frac{\rho r^*}{R}\right) \\ &= -\frac{1}{2\pi} \left(\ln(r) - \ln\left(\frac{\rho r^*}{R}\right)\right) \\ &= -\frac{1}{2\pi} \ln\left(\frac{Rr}{\rho r^*}\right) \end{aligned}$$

for all $(x, y) \neq (\xi, \eta)$. Now, it remains to show that this function is indeed the Green function. We have

$$\begin{aligned} \Delta G_R(x, y; \xi, \eta) &= \Delta \left(\Gamma(x - \xi, y - \eta) - \Gamma\left(\frac{\rho}{R} \left(x - \frac{R^2}{\rho^2} \xi\right), \frac{\rho}{R} \left(y - \frac{R^2}{\rho^2} \eta\right)\right) \right) \\ &= \Delta \Gamma(x - \xi, y - \eta) - \Delta \Gamma\left(\frac{\rho}{R} \left(x - \frac{R^2}{\rho^2} \xi\right), \frac{\rho}{R} \left(y - \frac{R^2}{\rho^2} \eta\right)\right) \\ &= -\delta(x - \xi, y - \eta) - 0 \\ &= -\delta(x - \xi, y - \eta). \end{aligned}$$

For all $(x, y) \in \partial D_R$ (that is, $x^2 + y^2 = R^2$), we have

$$\begin{aligned}(\tilde{x}, \tilde{y}) &:= \left(\frac{R^2}{x^2 + y^2} x, \frac{R^2}{x^2 + y^2} y \right) \\&= \left(\frac{R^2}{R^2} x, \frac{R^2}{R^2} y \right) \\&= (x, y),\end{aligned}$$

and so we have

$$\begin{aligned}G_R(x, y; \xi, \eta) &= \Gamma(x - \xi, y - \eta) - \Gamma\left(\frac{\sqrt{\xi^2 + \eta^2}}{R}(x - \tilde{\xi}, y - \tilde{\eta})\right) \\&= \Gamma(x - \xi, y - \eta) - \Gamma\left(\frac{R}{R}(x - \xi, y - \eta)\right) \\&= \Gamma(x - \xi, y - \eta) - \Gamma(x - \xi, y - \eta) \\&= 0.\end{aligned}$$

In summary, $G(x, y; \xi, \eta)$ solves

$$\begin{aligned}\Delta G &= -\delta(x - \xi, y - \eta) & (x, y) \in D, \\G(x, y; \xi, \eta) &= 0 & (x, y) \in \partial D,\end{aligned}\tag{8.14}$$

meaning that G is indeed the Green function. □